

A long time ago in a galaxy far,
far away....

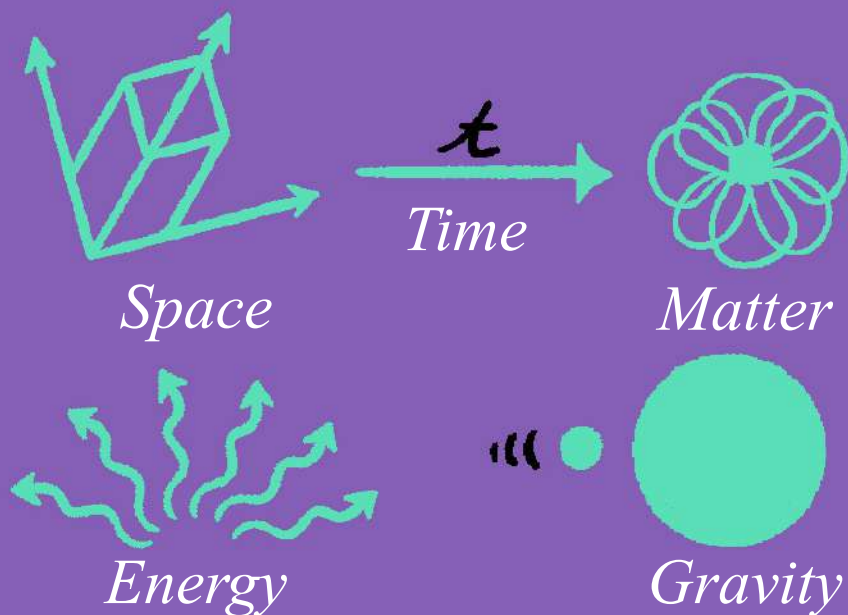


There are particular objects in the universe,
including the universe itself, that question
astrophysicists...

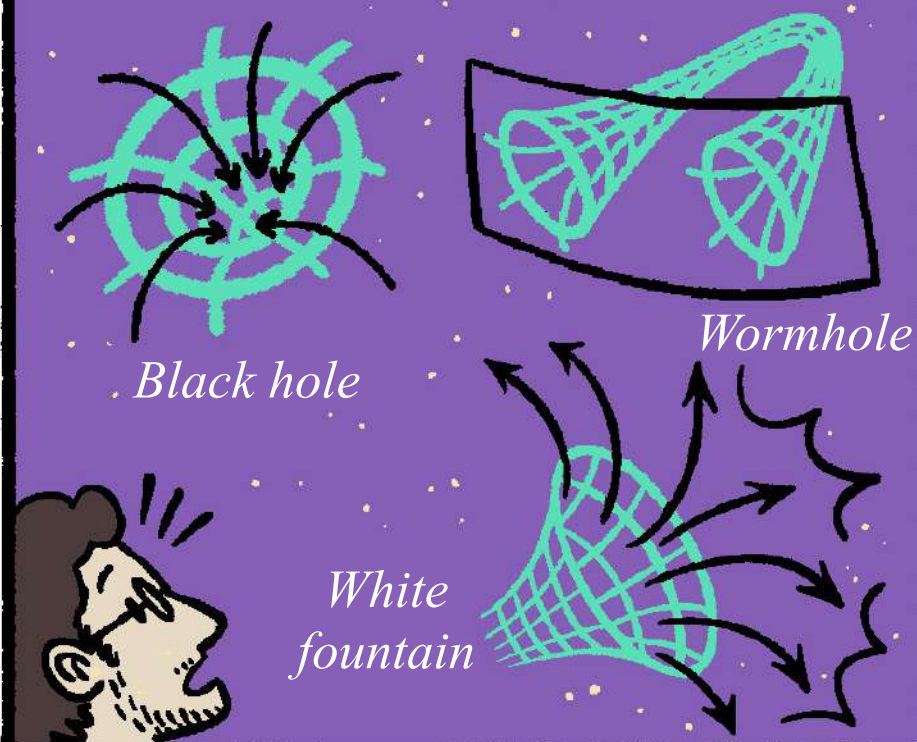
"Remember to look up at the stars and not down at your feet",
like Hawking said...

Futuroscope (France)

These objects mingle...



Like a...



And all these objects
can rotate, making
their understanding
even more complex...



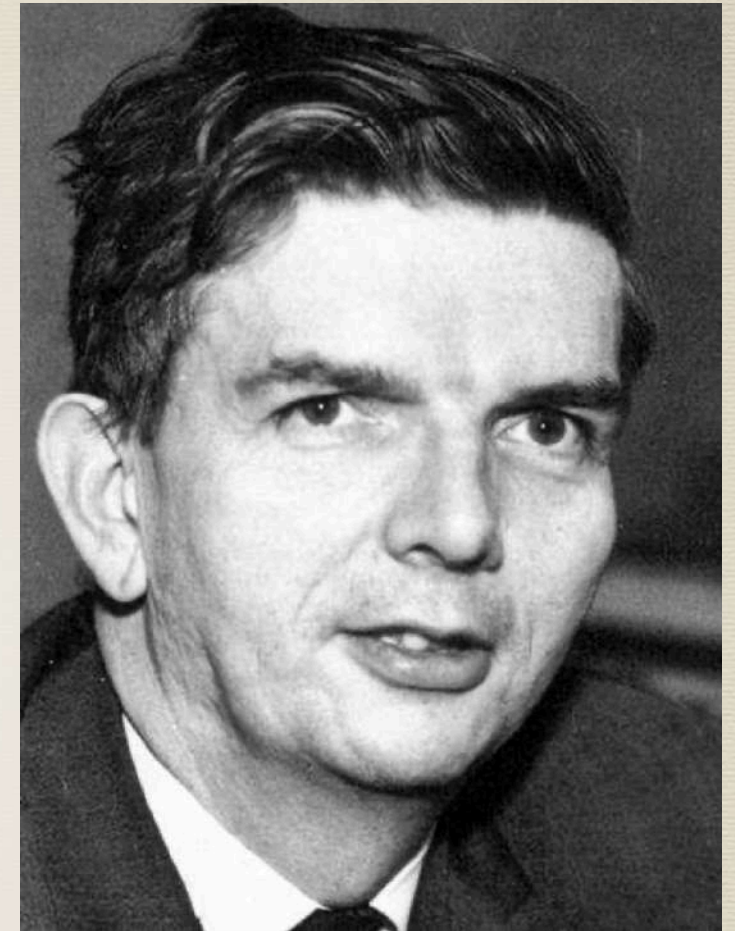
The Black Hole Prison of Fort William in Calcutta

Physicist Robert Dicke in 1961 spoke of peculiar celestial bodies (gravitationally completely collapsed stars) similar to "the black hole of Calcutta".

The Black Hole of Calcutta.

A scientist, writing of the black hole of Calcutta and its atmosphere, says:

"On the 20th of June, 1756, about 8 o'clock in the evening, 146 men were forced at the point of the bayonet into a dungeon 18 feet square. They had been but a few minutes confined in this infernal prison before every one fell into a perspiration so profuse that no idea can be formed of it. This brought on a raging thirst, the most difficult respiration and an outrageous delirium. Such was the horror of their situation that every insult that could be devised against the guard without and all the opprobrious names the viceroy and his officers could be loaded with were repeated to provoke the guard to fire upon them and terminate their sufferings. Before 11 o'clock the same evening one-third of the men were dead, and before 6 next morning only 23 came out alive, but most of them in a high putrid fever. All these dreadful effects were occasioned by the want of atmospheric air and by their breathing a superabundant quantity of nitrogen emitted from their lungs."

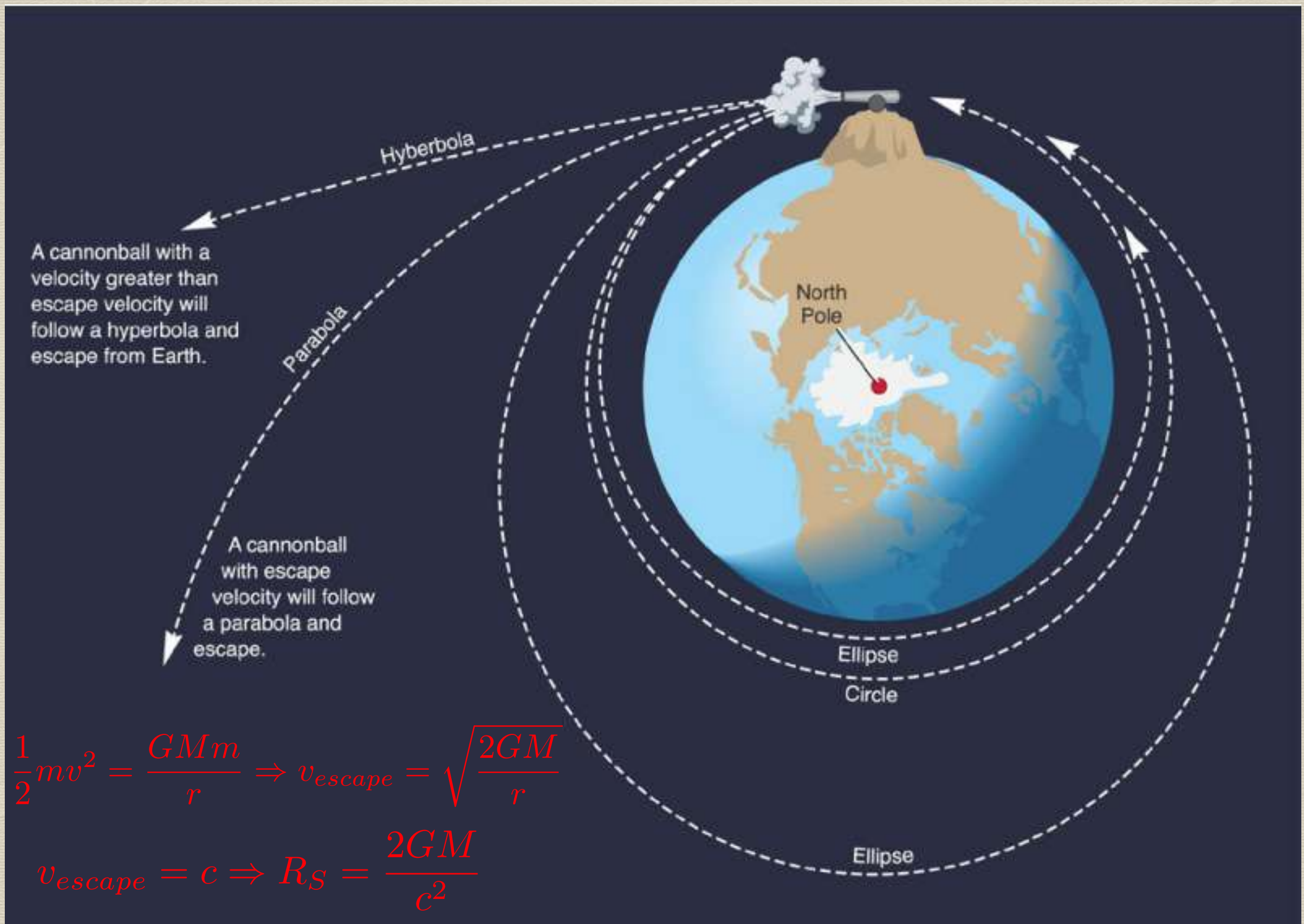


Robert Dicke (1916-1997).

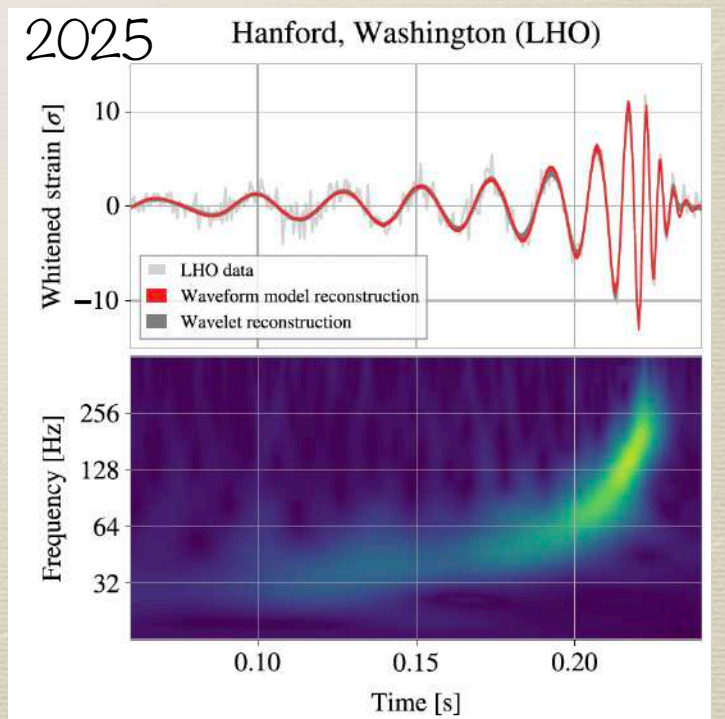
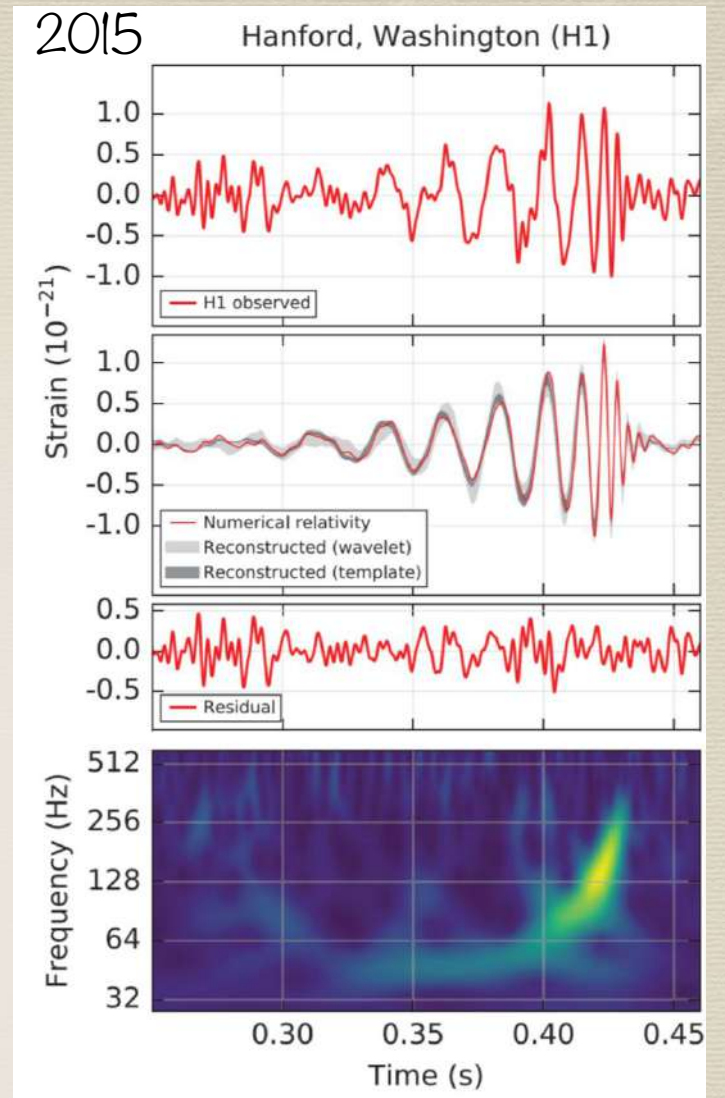
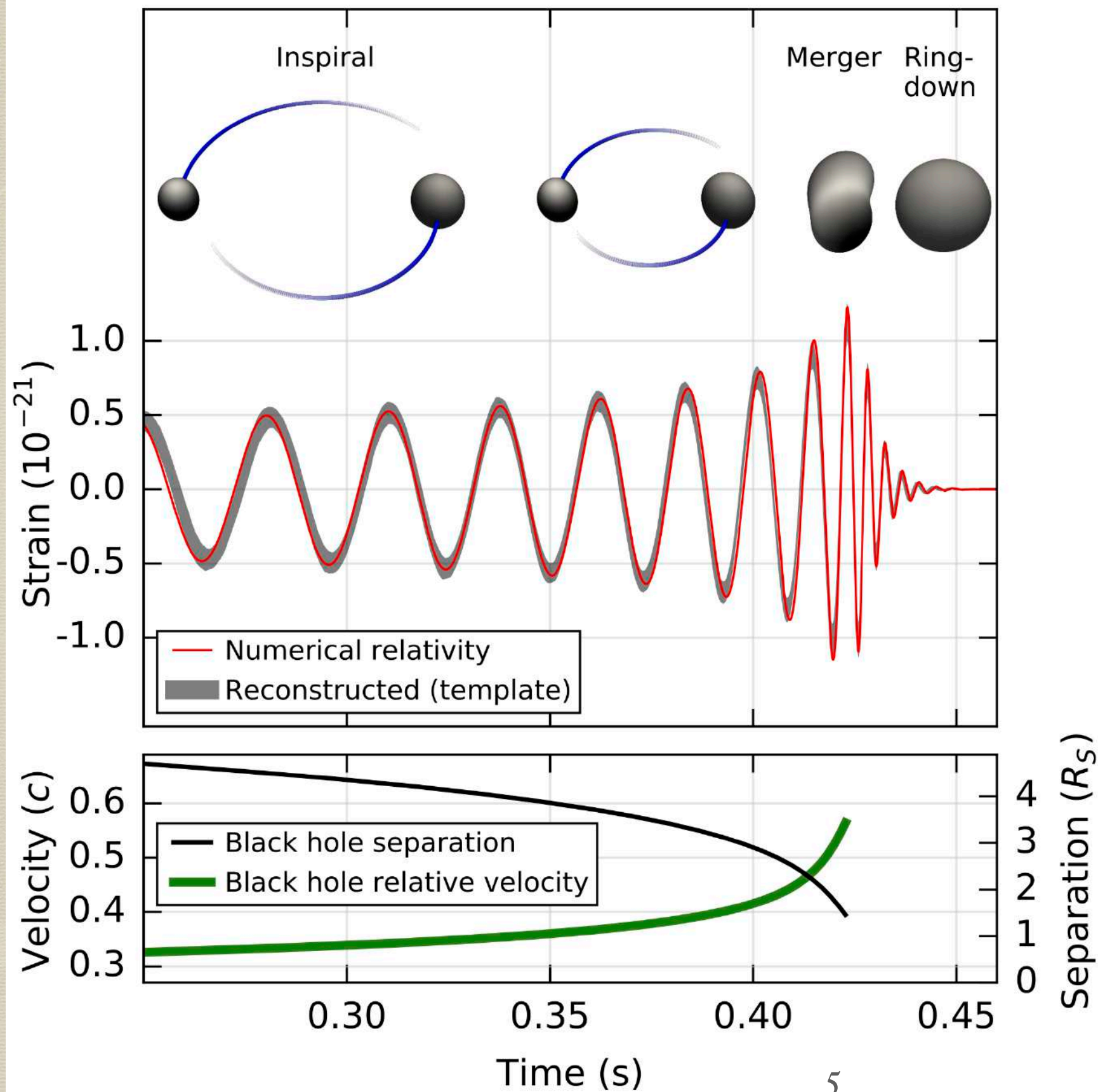
*To the memory of the 123 persons
Who perished in the Black Hole prison
Of old Fort William*

"The Black Hole was a small dungeon in the old Fort William in Calcutta where troops of the Nawab of Bengal, Siraj ud-Daulah, held British prisoners of war after the capture of the fort on 20 June 1756. British and Anglo-Indian soldiers and civilians were held overnight in conditions so cramped that many died from suffocation, heat exhaustion and crushing."

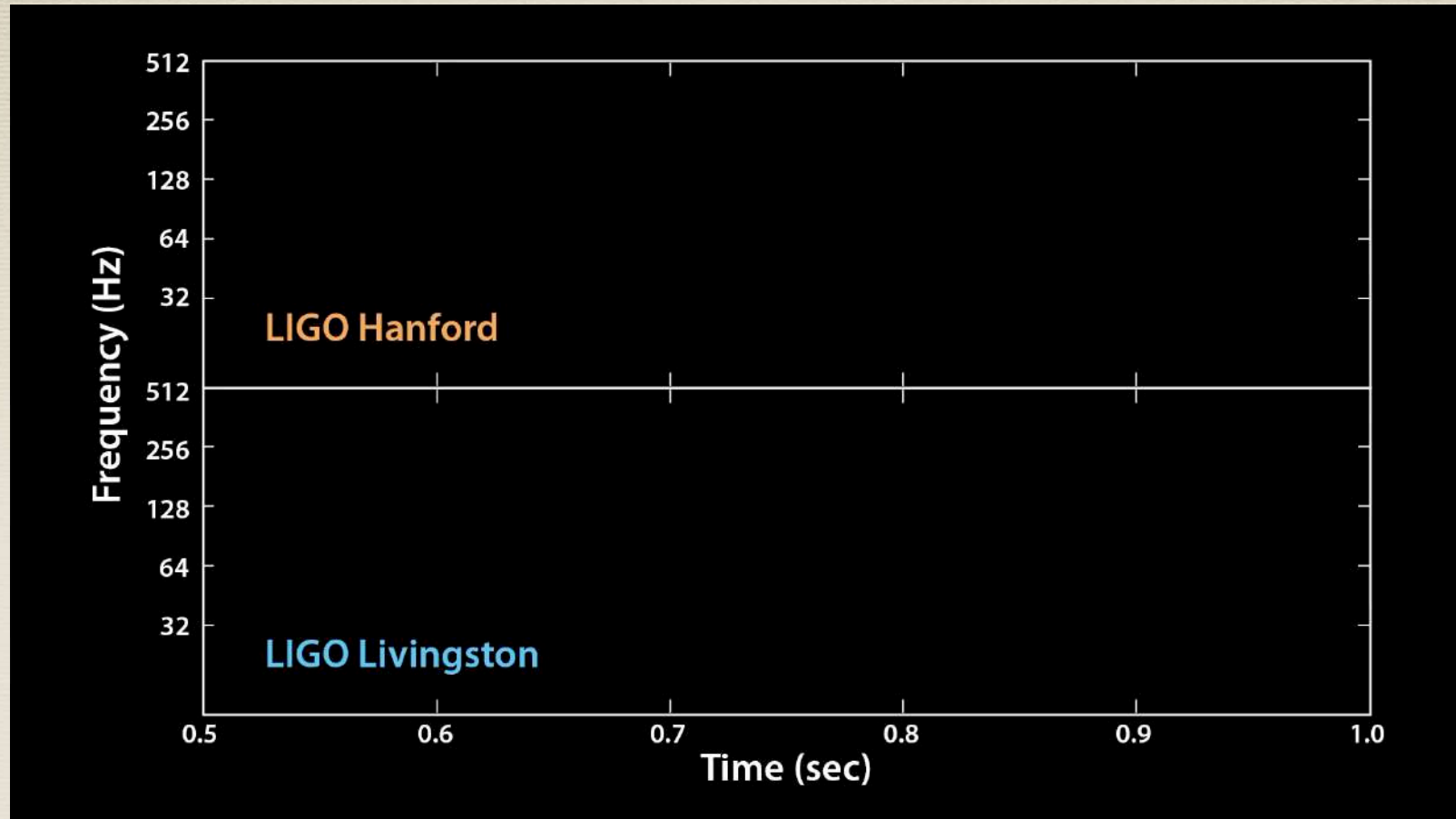
Escape Velocity and Schwarzschild Horizon



A Black Holes Merger as seen by the LIGO Interferometer in 2015/2025



The Sound of Two Black Holes Colliding



Audio Credit: Caltech/MIT/LIGO Lab

Gravitational waves sent out from a pair of colliding black holes have been converted to sound waves, as heard in this animation. On September 14, 2015, LIGO observed gravitational waves from the merger of two black holes, each about 30 times the mass of our sun. The incredibly powerful event, which released 50 times more energy than all the stars in the observable universe, lasted only fractions of a second.

In the first two runs of the animation, the sound-wave frequencies exactly match the frequencies of the gravitational waves. The second two runs of the animation play the sounds again at higher frequencies that better fit the human hearing range.

The animation ends by playing the original frequencies again twice.

As the black holes spiral closer and closer in together, the frequency of the gravitational waves increases. Scientists call these sounds "chirps," because some events that generate gravitation waves would sound like a bird's chirp.

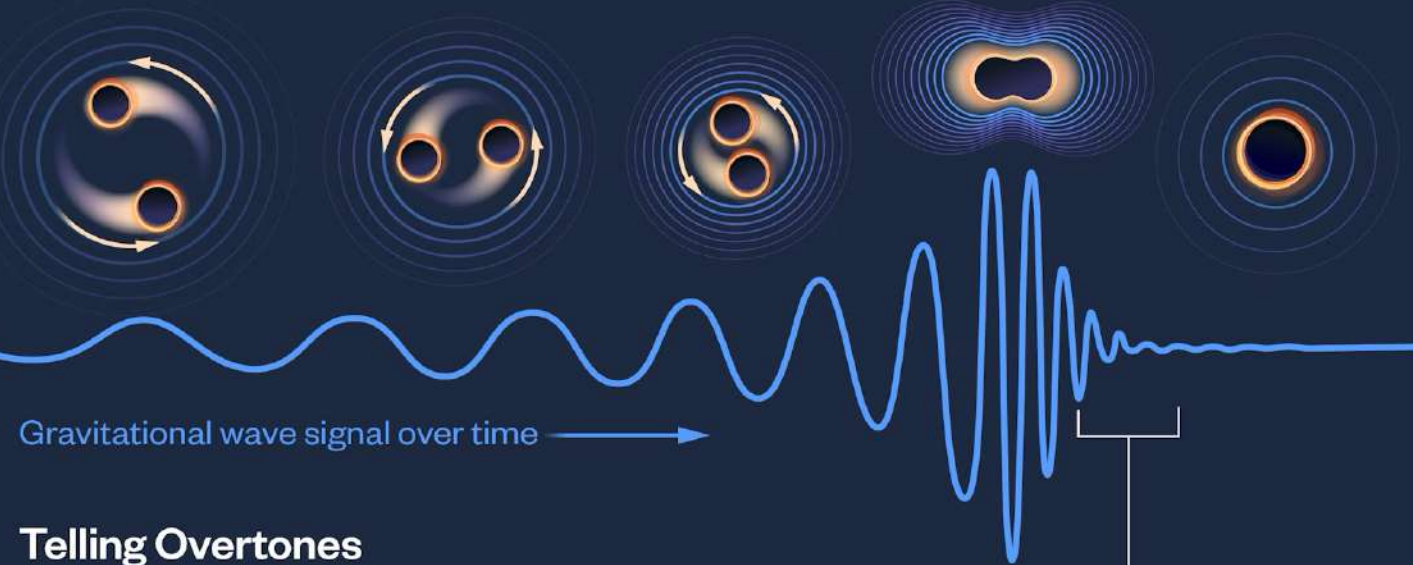
Clear Signal Sheds Light on Black Holes

When two black holes (BHs) collide and merge, they release gravitational waves. These waves can be detected by sensitive instruments on Earth, allowing scientists to determine the mass and spin of the BHs. The clearest BH merger signal yet, named GW250114, recorded by LIGO in January 2025, offers new insights into these mysterious cosmic giants.

Two BHs release gravitational waves as they spiral inward

Merger

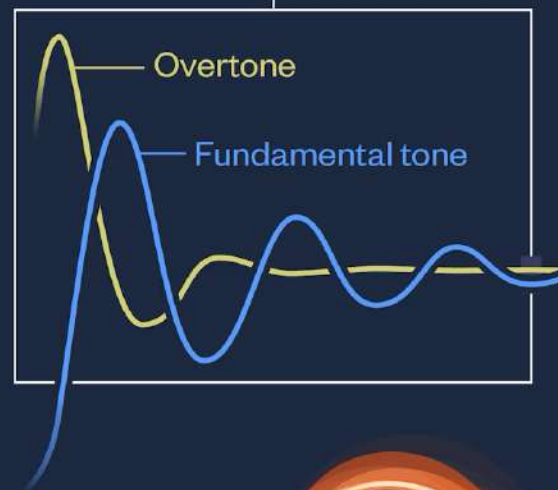
BH stabilizes



Gravitational wave signal over time

Telling Overtones

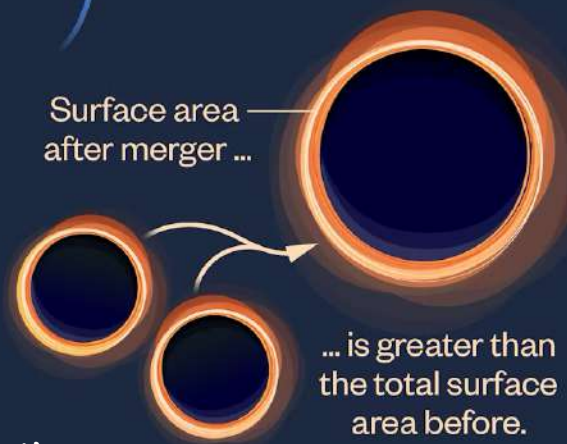
A fleeting secondary tone was detected in the signal, offering a rare chance to test the “**Kerr solution**,” which describes a rotating BH using only mass and spin. Excitingly, the mass and spin values from this **overtone** matched those from the **fundamental tone**. If they had differed, it would imply that additional properties are necessary to describe a BH, but a match confirms that — at least for this BH — no other details are needed.



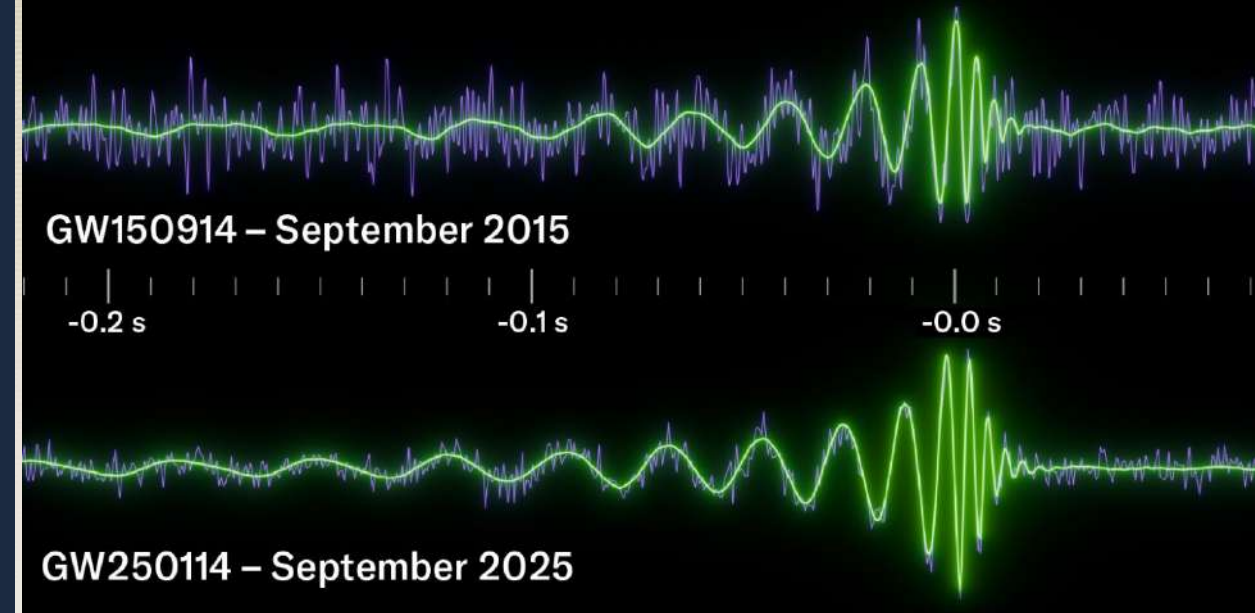
Forever Growing

The signal also tested **Hawking’s area theorem**, which states that a BH’s surface area can never decrease — it can only grow. Surface area of a BH is determined by the area of its **event horizon** and is proportional to the square of the BH’s mass. Comparing the BHs before and after the merger confirmed that the surface area had increased, supporting the theorem.

Courtesy L. Reading-Ikkanda/Simons Foundatio

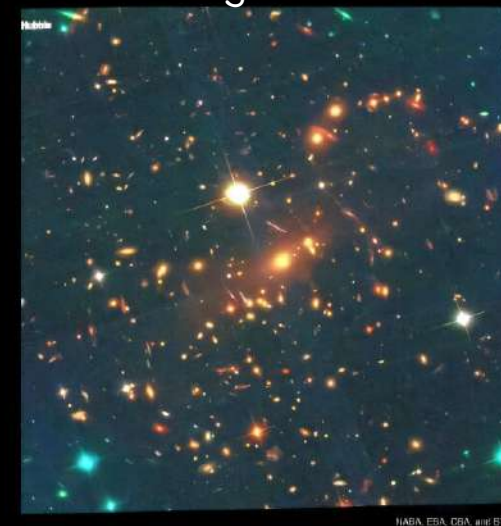


Courtesy LIGO/J. Tissino (GSSI)/R. Hurt (Caltech-IPAC)

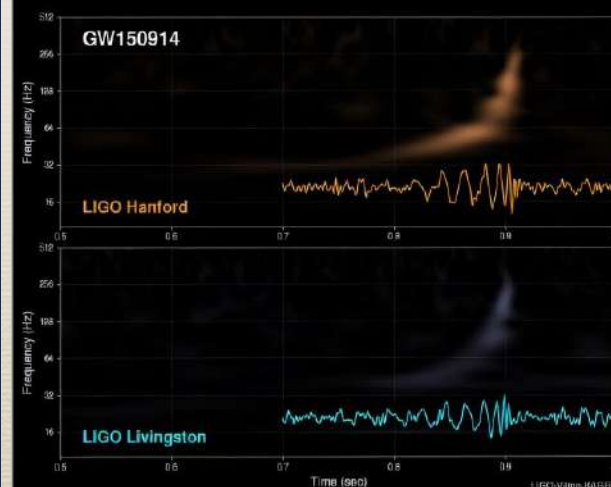


A New Perspective on the Universe

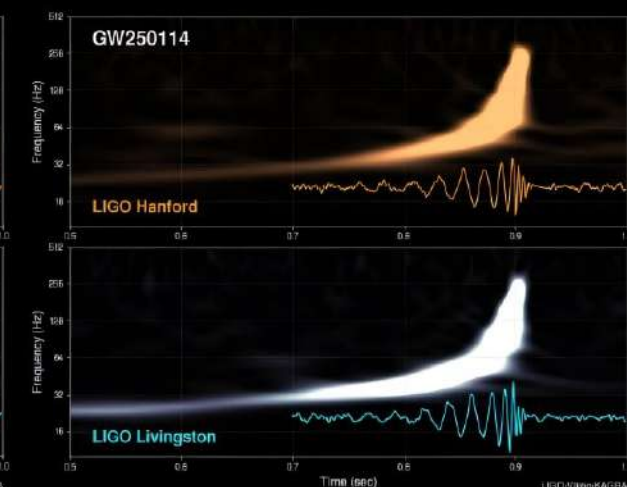
Courtesy Dr. Derek Davis (Caltech, LIGO Laboratory)



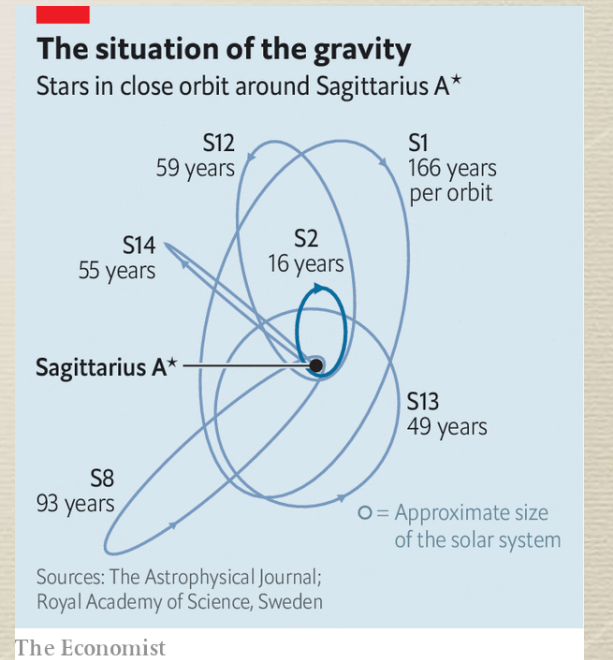
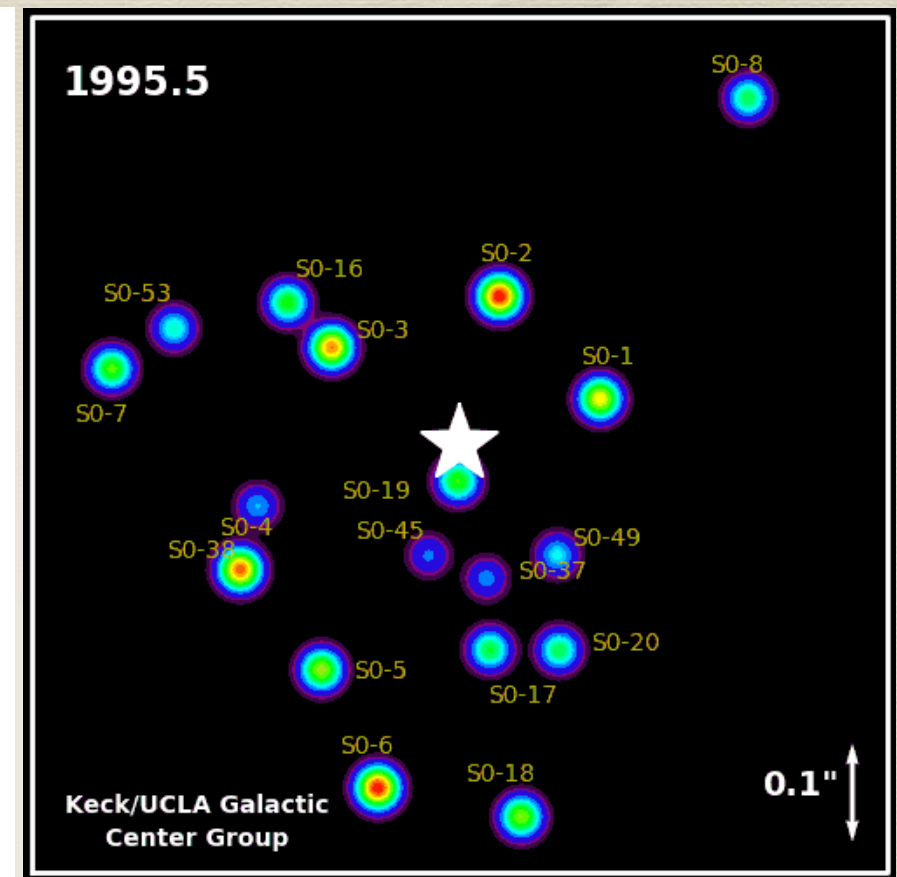
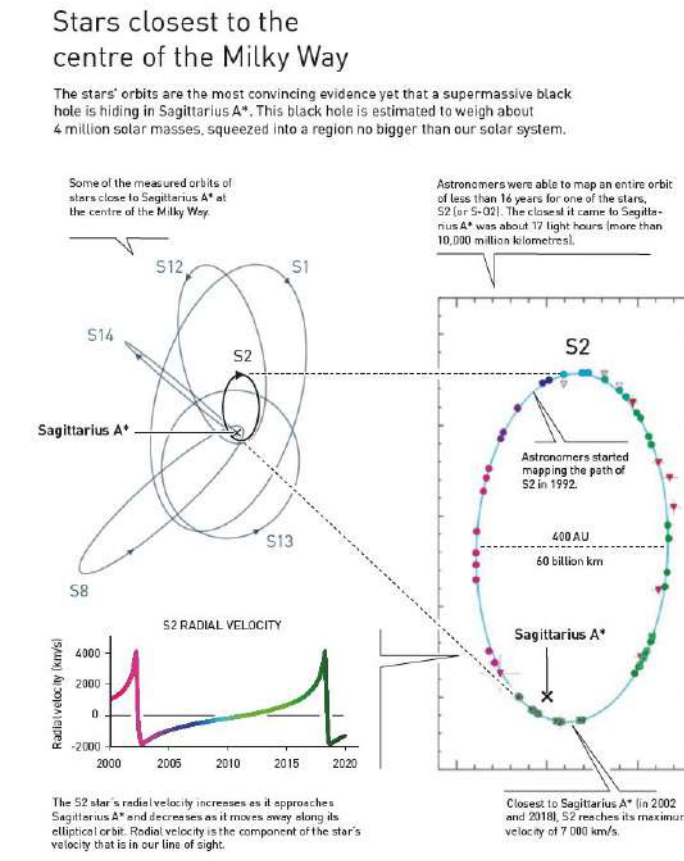
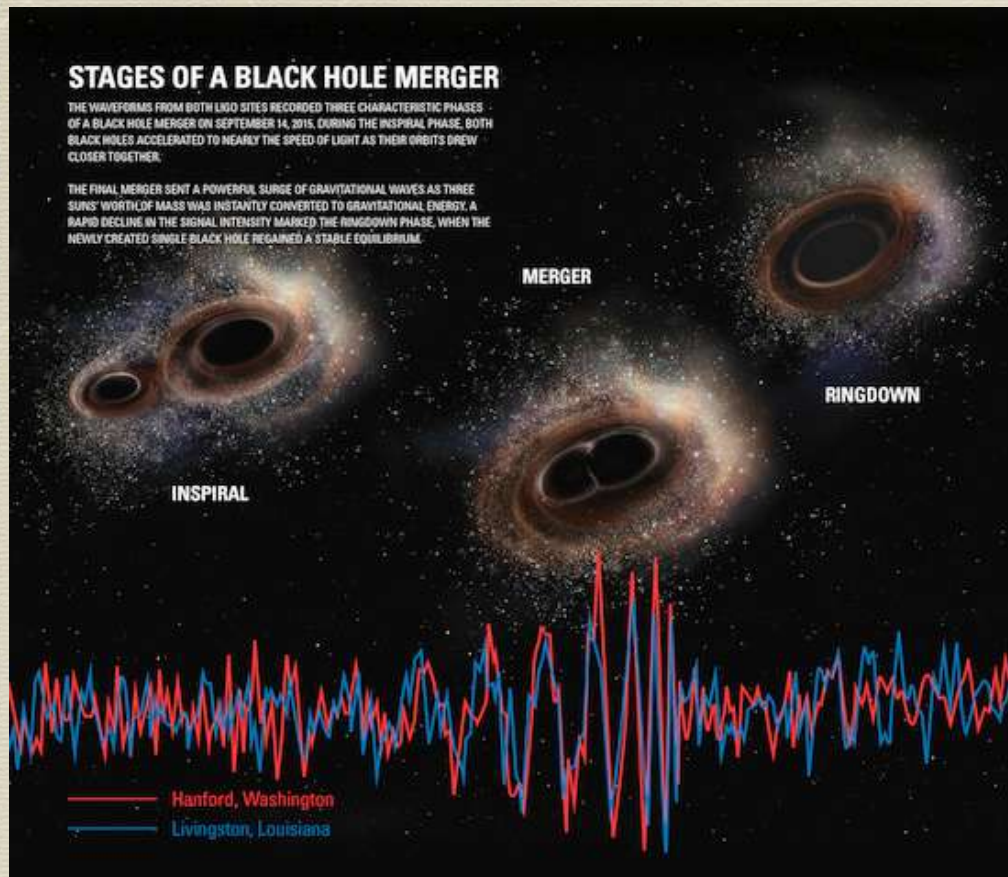
LIGO - First Observing Run (2015)



LIGO - Fourth Observing Run (2025)

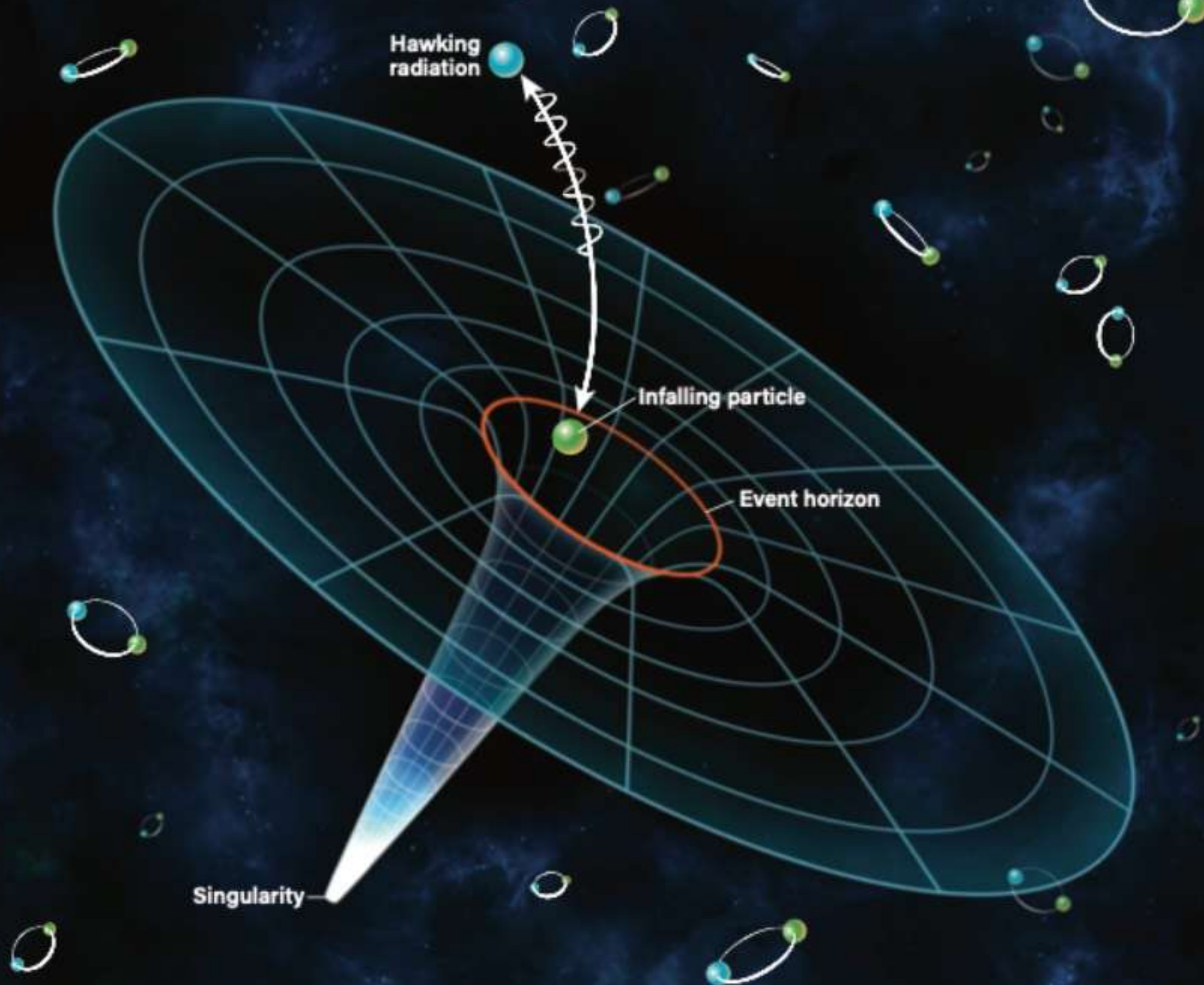


A New Golden Age for Black Hole Physics



"As you probe ever closer to the horizons of the black holes, Nature might have new surprises in store... » Ulf Danielsson, during the announcement of the Nobel Prize in Physics (October 2020).

Black Holes Should Radiate a Quantum Glow

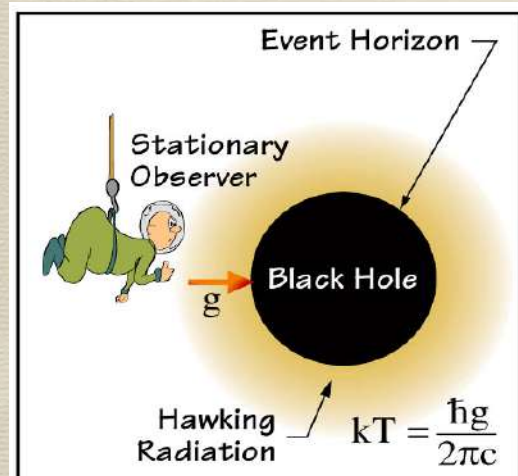


Hawking/Unruh Temperature

$$E = h\nu = \frac{hc}{\lambda} \quad \frac{1}{2}mv^2 = \frac{GMm}{r} \Rightarrow v_{escape} = \sqrt{\frac{2GM}{r}} \quad \lambda \simeq R_S = \frac{2GM}{c^2}$$

$$E \simeq \frac{hc^3}{2GM}$$

1974



A stationary observer outside the black hole would see the thermal Hawking radiation.

$$v_{escape} = c \Rightarrow R_S = \frac{2GM}{c^2}$$

$$T \simeq \frac{hc^3}{2k_B GM}$$

$$T_H = \frac{hc^3}{16\pi^2 k_B GM}$$

$$E \simeq k_B T$$

1976



Surface gravity (peeling rate of the rays bundle) = acceleration of a plunging observer at the horizon seen by an asymptotic observer

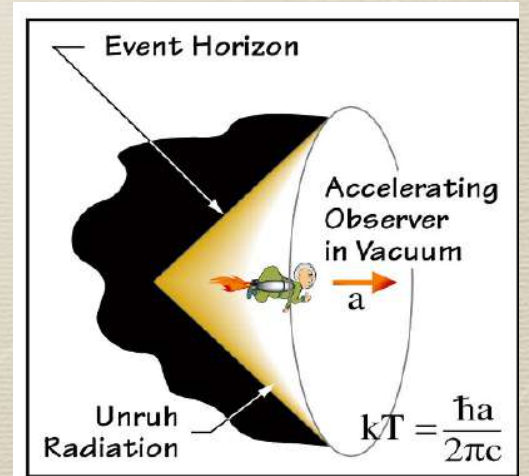
$$\kappa = \frac{v_{escape}^2(R_S)}{R_S} = \frac{2GM}{R_S} = \frac{c^4}{2GM}$$

$$T_H = \frac{\hbar\kappa}{2\pi k_B c}$$

$$T_U = \frac{\hbar a}{2\pi k_B c}$$

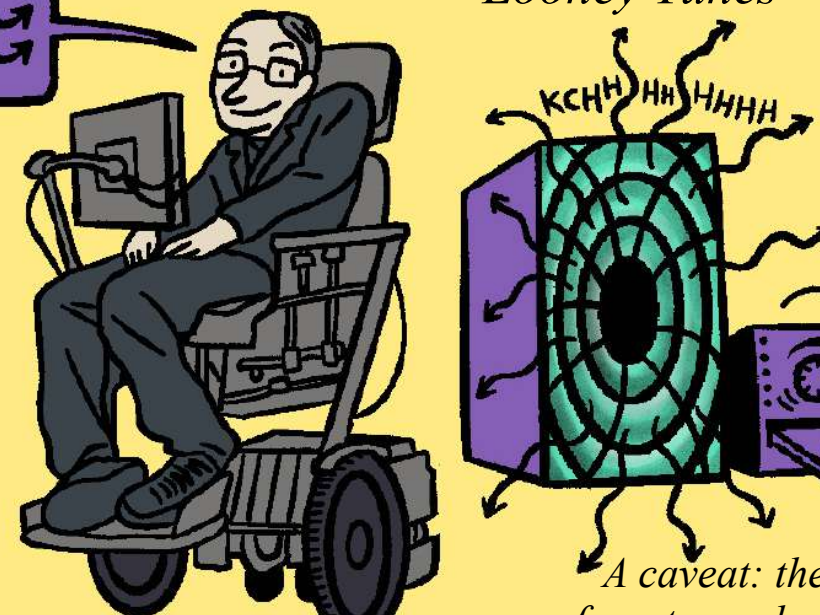
“God not only plays dice, he also sometimes throws the dice where they cannot be seen.”

Stephen Hawking



An accelerating observer in vacuum would see a similar Hawking-like radiation called Unruh radiation.

The late Stephen Hawking studied them and made incredible predictions.



Looney Tunes

Black holes radiate, in the sense that they amplify the fluctuations of the quantum vacuum, like a HIFI system boosting the electric noise without a CD. The vacuum fluctuations are the universe music.

A caveat: the quantum glow originates from transplanckian scales at the horizon.

These rays, these radiations, are unfortunately unobservable, completely masked by the primordial light of the Big Bang. It's like trying to see a firefly in a spotlight or hearing a whisper in a rock concert...



LIGHT OF THE BIG-BANG

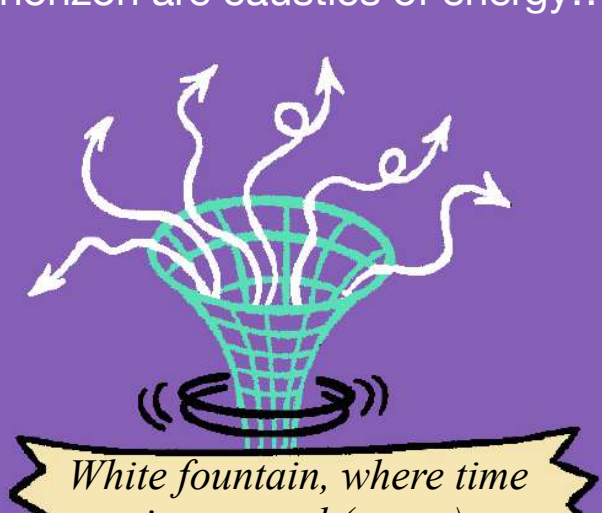
Black hole radiation

AAAAAH! I CAN'T SEE A THING!!

EVERYTHING THAT IS YET

UNKNOWN

We don't know if white fountains exist in the universe since they are considered as unstable as their horizon are caustics of energy...



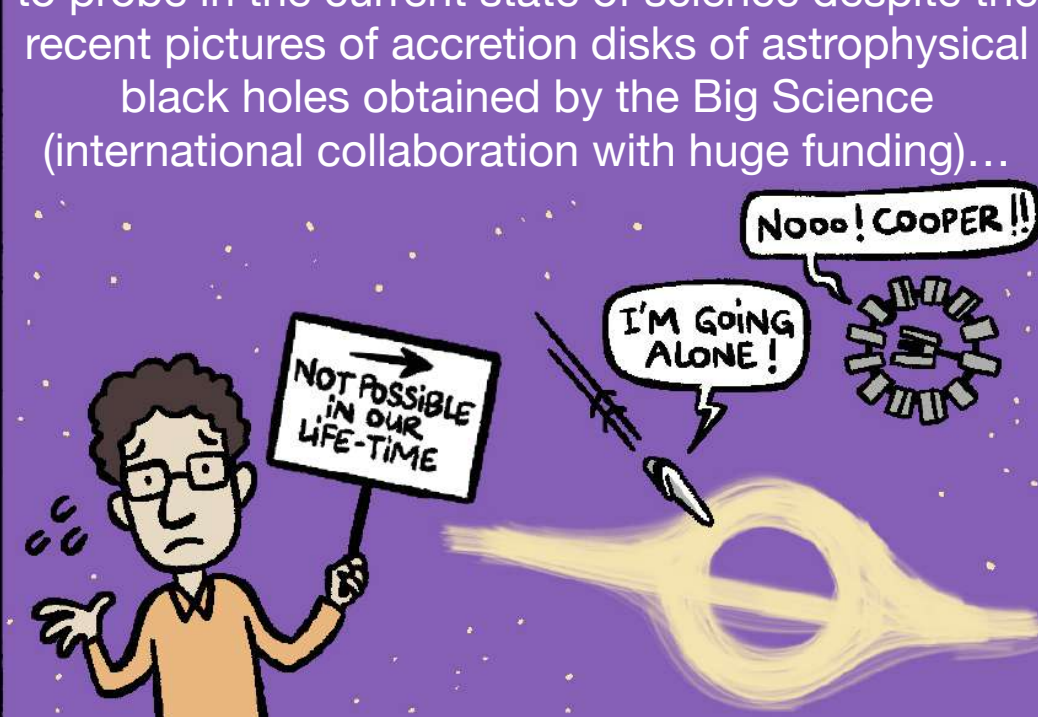
*White fountain, where time is reversed ($t \rightarrow -t$)
Myth or reality??*

We don't know what is inside the «heart» of a black hole...



«The undiscovered country from whose bourn No traveller returns, puzzles the will», in Hamlet.

Its horizon and its central singularity are impossible to probe in the current state of science despite the recent pictures of accretion disks of astrophysical black holes obtained by the Big Science (international collaboration with huge funding)...



NOT POSSIBLE IN OUR LIFE-TIME

I'M GOING ALONE!

Nooo! COOPER!!

New Scientist

WEEKLY 27 May 2023

CAN NEW ALZHEIMER'S
TREATMENTS FINALLY
CURB DEMENTIA?

WHY THINKING ABOUT
YOUR FUTURE SELF
CAN CHANGE YOUR
LIFE TODAY

HOW TO BUILD A BLACK HOLE

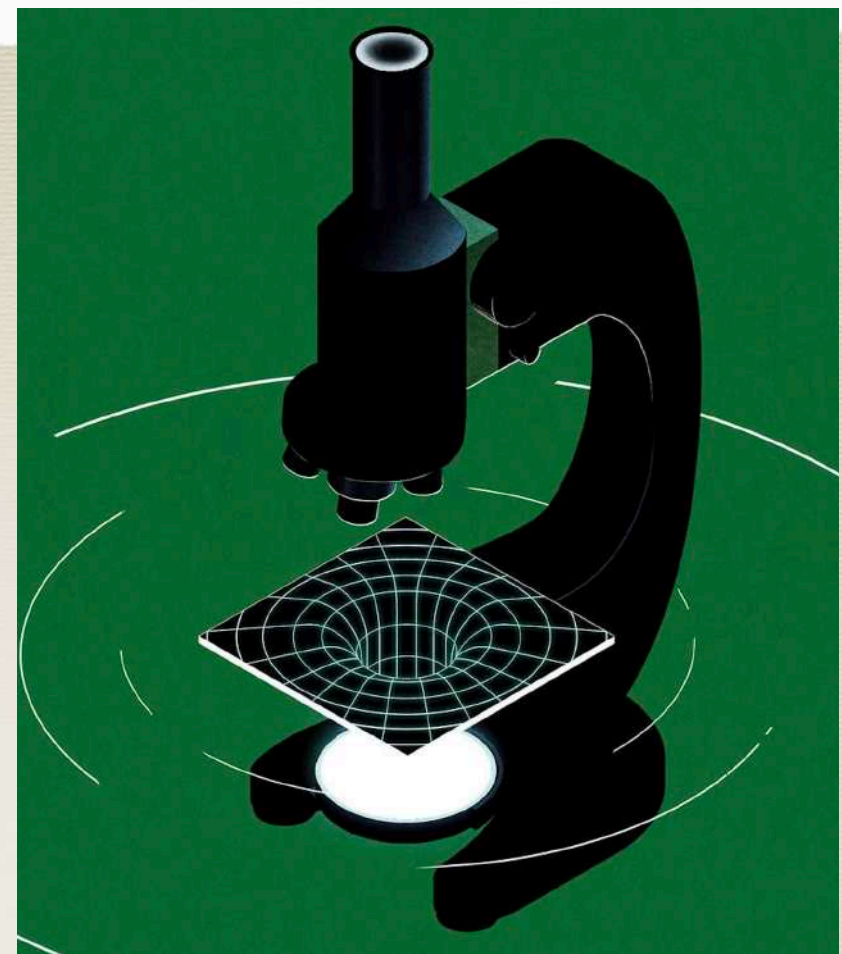
Inside the audacious attempts to recreate
extreme gravity in the lab



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LOST BEES / DO OCTOPUSES HAVE NIGHTMARES? / **BAT CAVES**



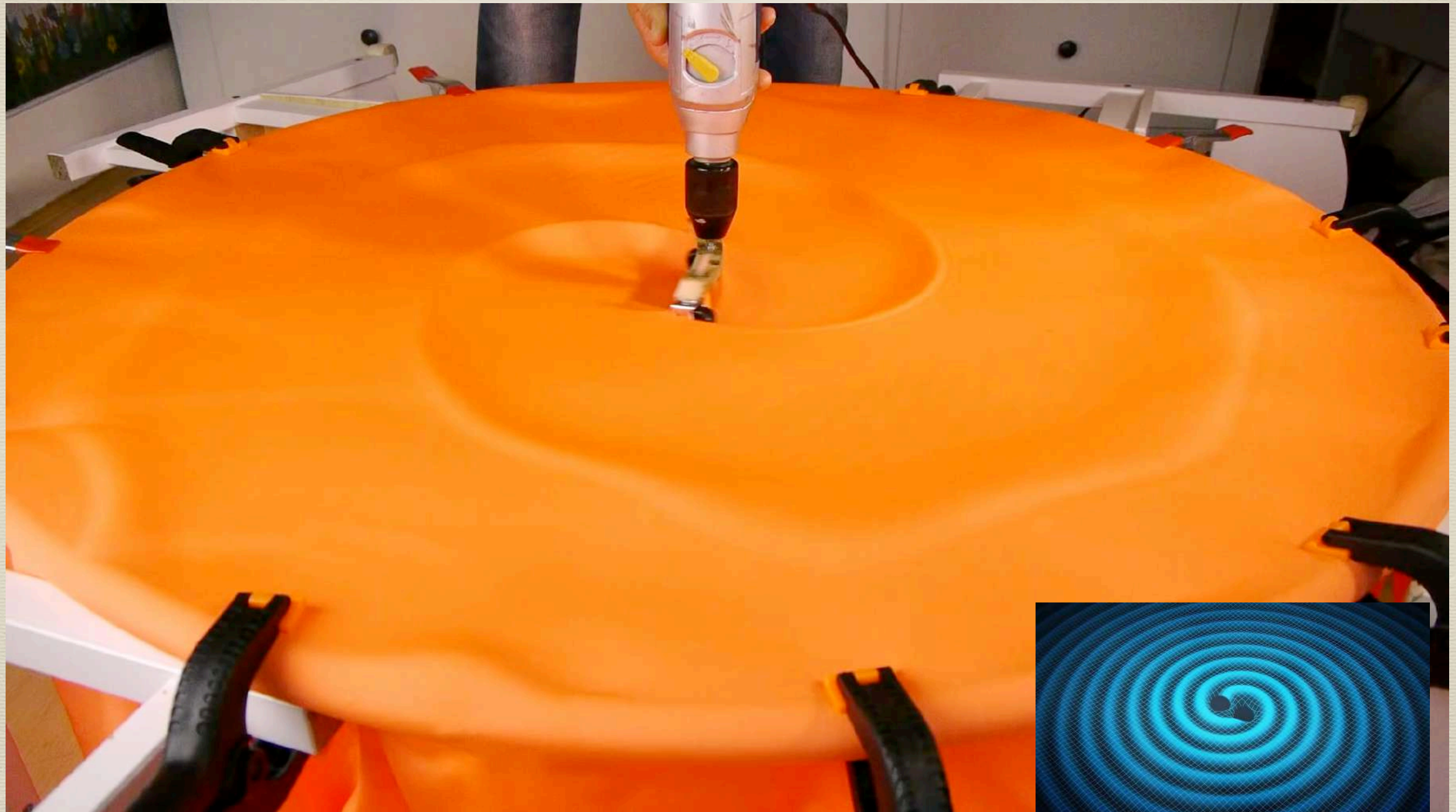
An Economic Black Hole for Tourists



Analogue Simulations at Home

« It seems to me that the test of 'Do we or do we not understand a particular subject in physics?' is
'Can we make a mechanical model of it?' »

Lord Kelvin

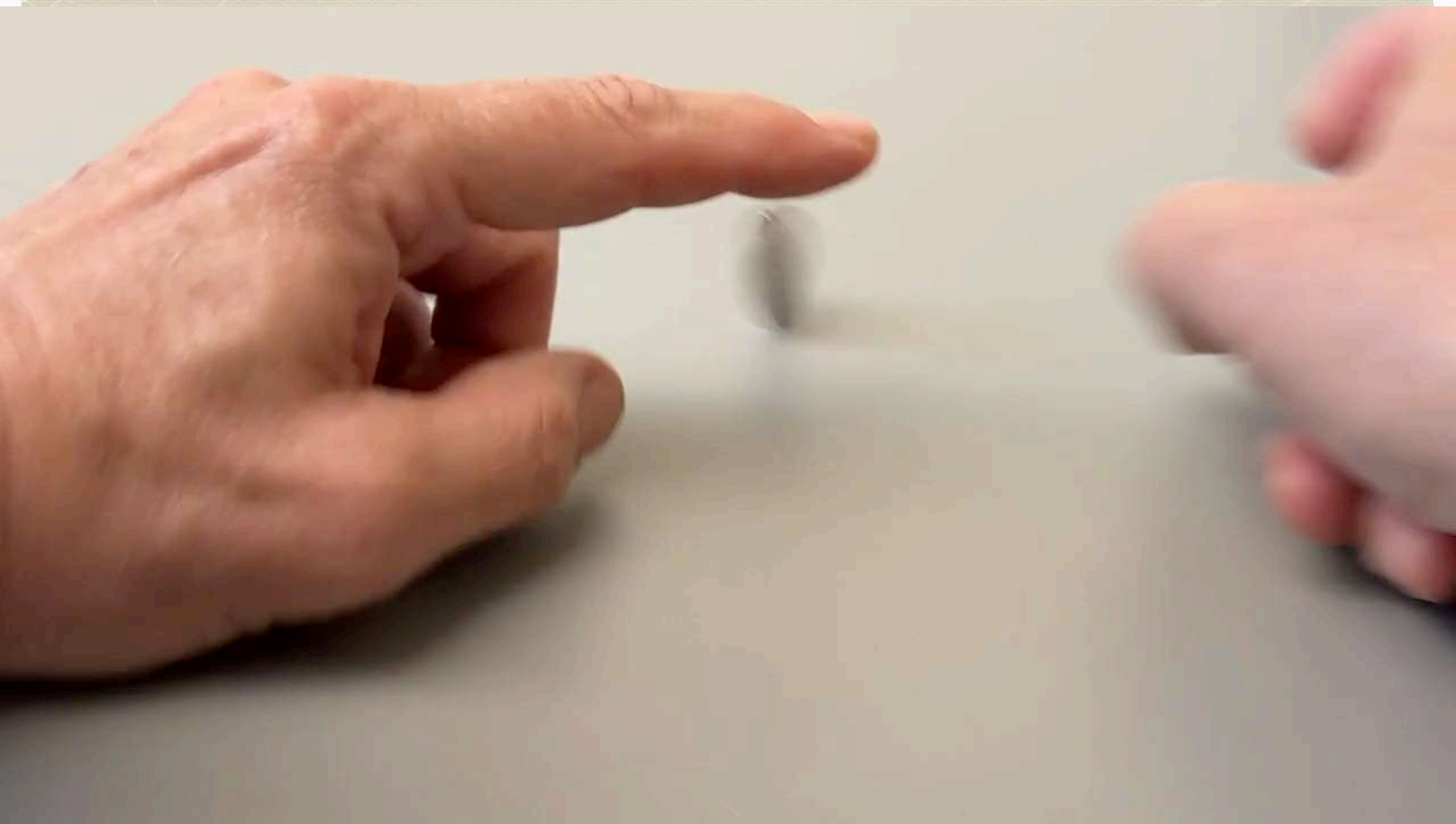


Ingredients : a tablecloth made of stretched spandex lycra + a drill + two roulettes

<https://www.youtube.com/watch?v=dw7U3BYMs4U>

<http://thekidshouldseethis.com/post/69096529285>

The Inspiring Power of Analogies



D.D. Meringolo, F. Conidi, A. Mercuri, M. Sposato, R.C. Barberi and G. Pucci, On the analogy between spinning disks coming to rest and merging black holes, *American Journal of Physics*, 93(7), 551-556 (2025).

Finite Time Singularity of the Euler's Disk

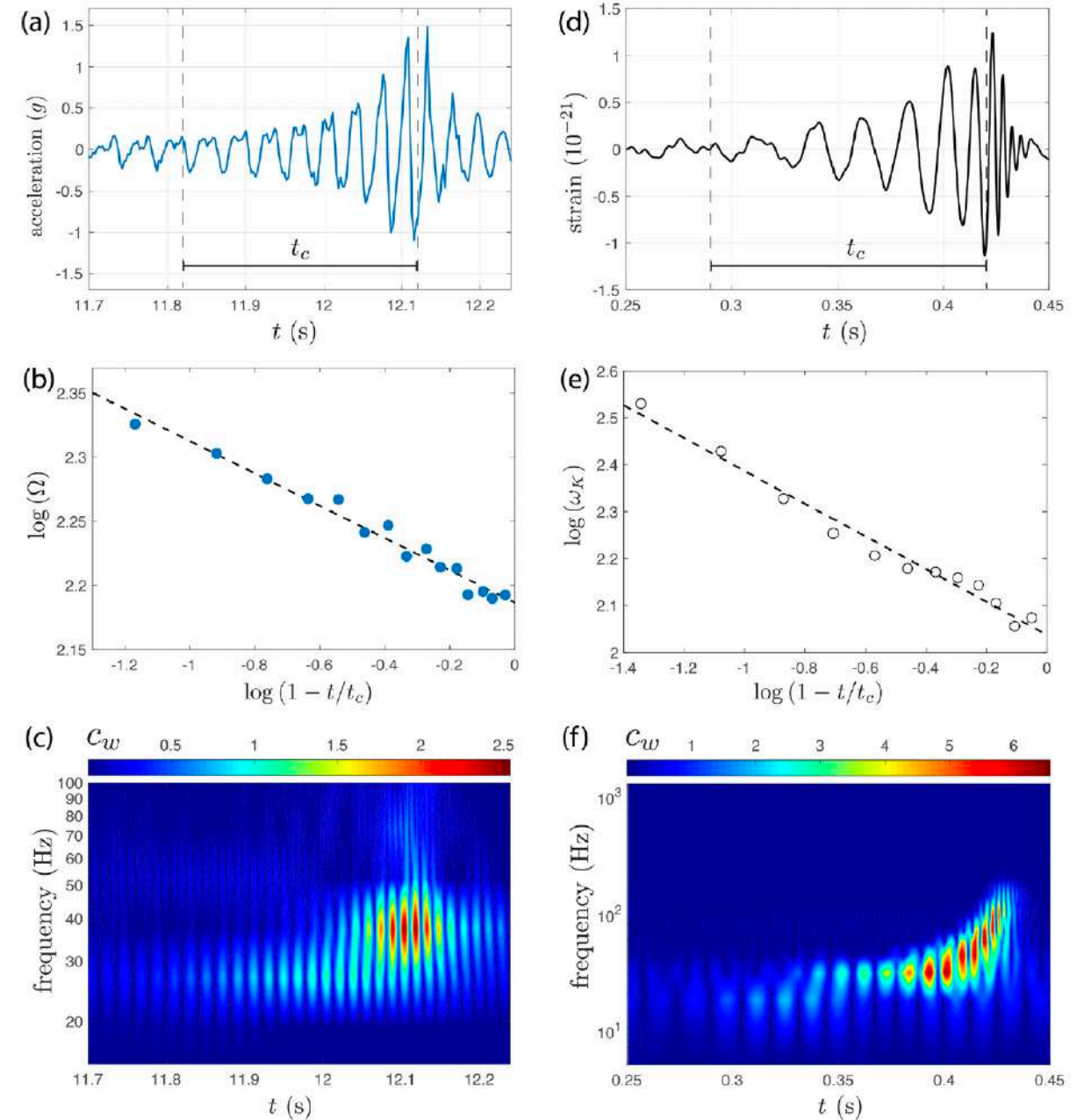


Spinning Disk or Black Holes Merging

(Small Science versus Big Science)

Table I. Comparison of the relevant quantities of a disk coming to rest and two black holes merging.

Coming to rest of a spinning disk	Merger of two orbiting black holes
Energy conservation, $dE_{\text{disk}}/dt + P_{\text{diss}} = 0$	Energy conservation, $dE_{\text{orb}}/dt + P_{\text{rad}} = 0$
Disk energy, $E_{\text{disk}} = \frac{3}{2}mgR\alpha$	Orbital energy, $E_{\text{orb}} = -\frac{1}{2}G\mu M/l$
Dissipated power, $P_{\text{diss}} = k\alpha^{-\gamma}$	Gravitational radiation, $P_{\text{rad}} = \frac{32}{5}G^4\mu^2M^3/c^5l^5$
Coming to rest time, $t_c = 3\alpha_0^{\gamma+1}mgR/2k(\gamma+1)$	Coalescence time, $t_c = \frac{5}{256}c^5l_0^4/G^3\mu M^2$
Contact angle, $\alpha = \alpha_0(1-t/t_c)^{1/\gamma+1}$	Orbital separation, $l = l_0(1-t/t_c)^{1/4}$
Precession angular velocity, $\Omega = \Omega_0(1-t/t_c)^{-1/2(\gamma+1)}$	Orbital angular velocity, $\omega_K = \omega_{K0}(1-t/t_c)^{-3/8}$

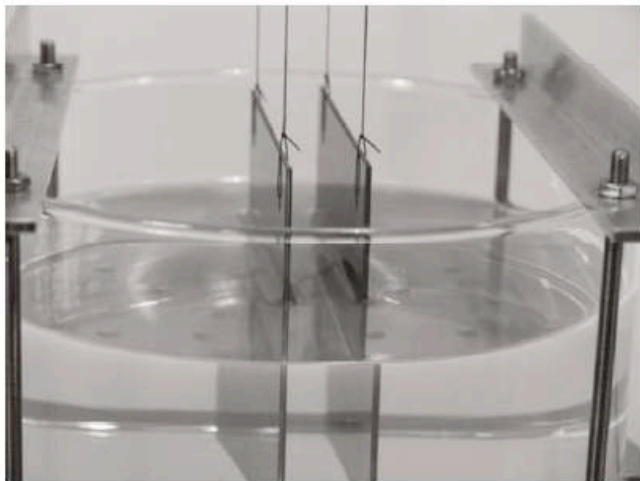


D.D. Meringolo, F. Conidi, A. Mercuri, M. Sposato, R.C. Barberi and G. Pucci, On the analogy between spinning disks coming to rest and merging black holes, American Journal of Physics, 93(7), 551-556 (2025).

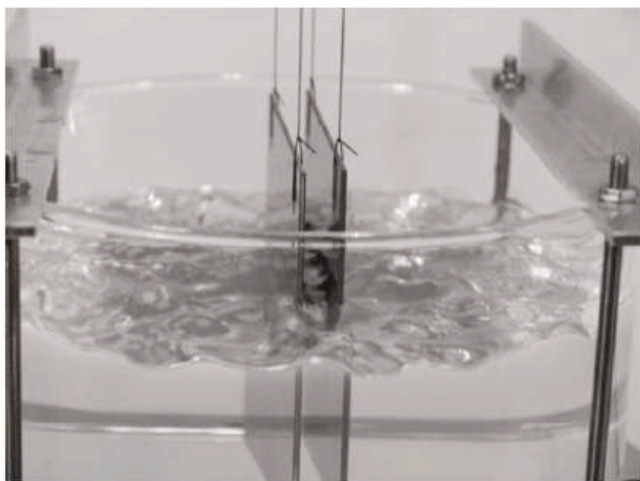
Casimir Effect with Water Waves : Another Quantum Fluctuations Process.



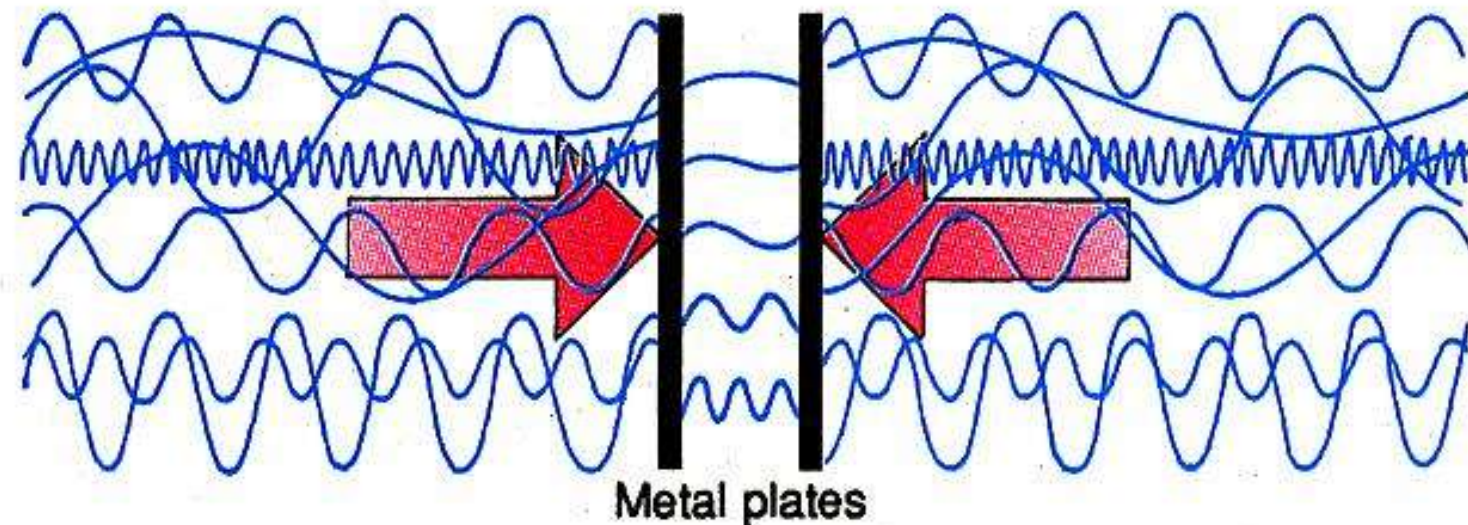
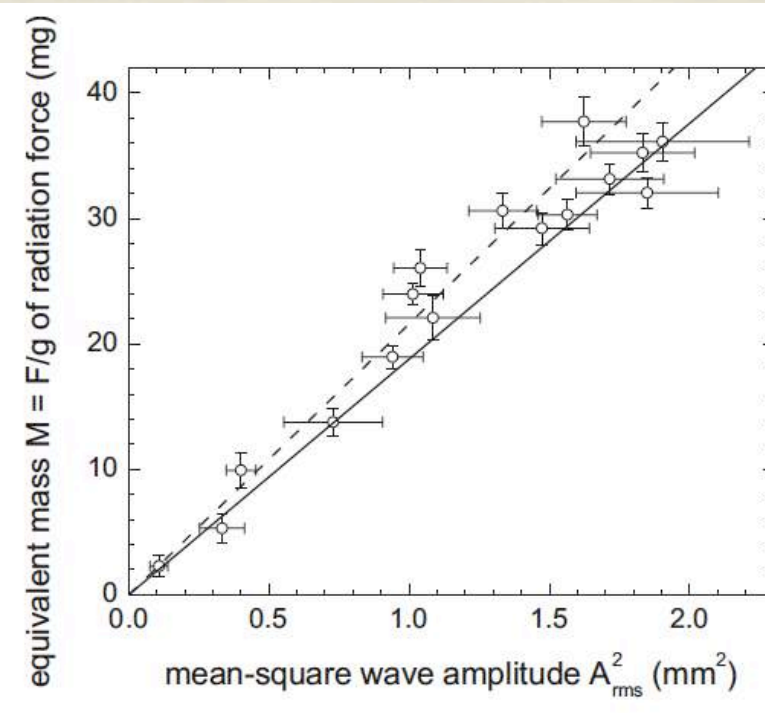
(a)



(b)



(c)



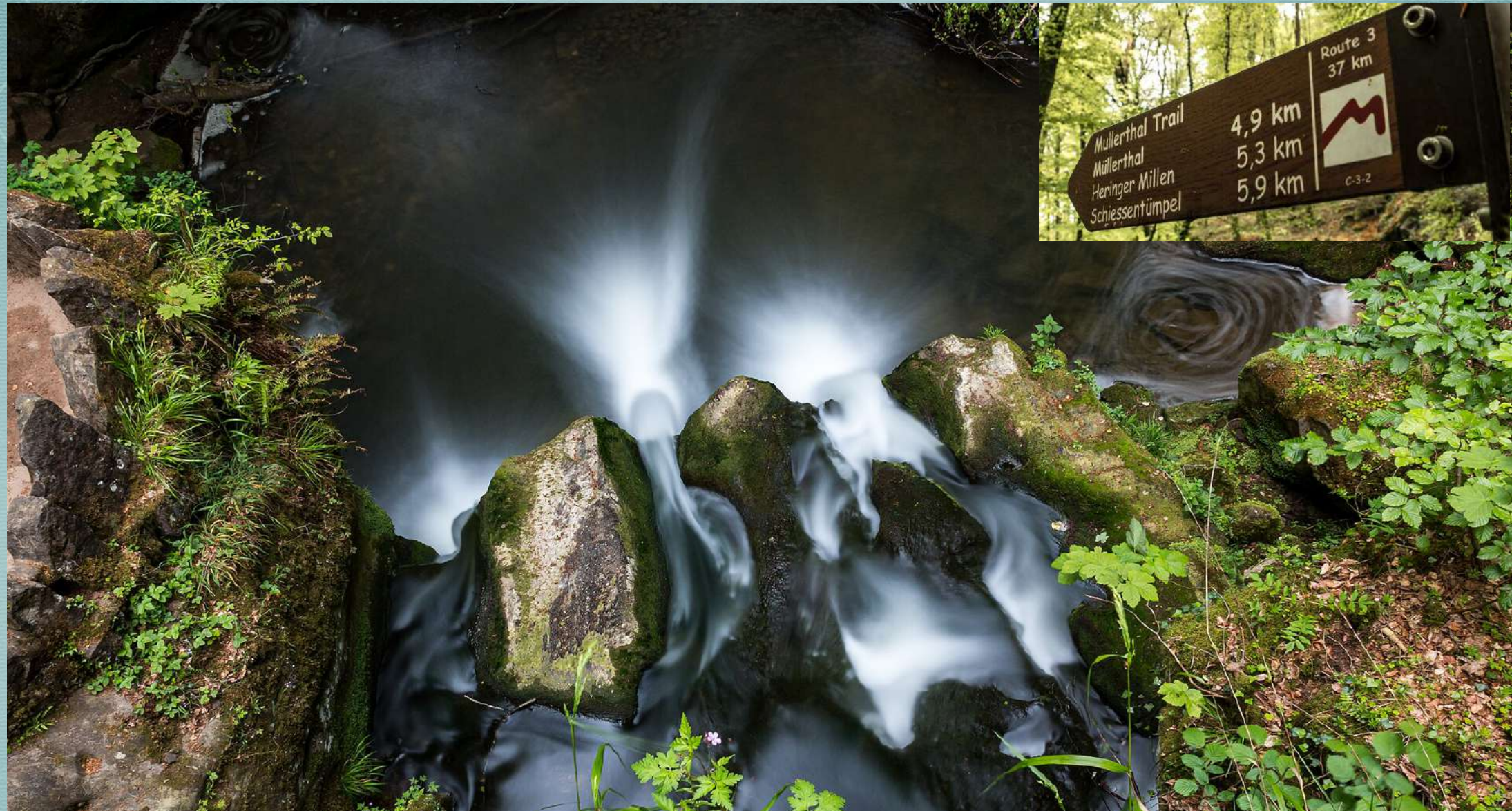
The Casimir effect: an imbalance in the quantum fluctuations of empty space can push two metal plates together

A water wave analog of the Casimir effect

Bruce C. Denardo,^{a)} Joshua J. Puda, and Andrés Larraza

Am. J. Phys., Vol. 77, No. 12, December 2009

ANALOGUE AQUATIC SPACETIMES



Germain Rousseaux (Research Director CNRS) and Co.,

Team Leader of Curiosity,

Institut Pprime, Futuroscope, Poitiers, (France)

IAS Luxembourg University, Quantum Analogue Workshop,

15th-16th September 2025 Luxembourg



A Warning

G. BRUMFIEL/D. ALLISON

Cosmos in a Bottle,
Nature, Vol. 451,
17, 236, January
2008.

“Frankly,” says Wolfgang Ketterle, a Nobel-Prize-winning condensed-matter physicist at the Massachusetts Institute of Technology in Cambridge, “I don’t think a table-top experiment will answer fundamental questions about the cosmos any time soon.” “It is really pushing experimental technique to its limits,” says Ketterle. Basing a career on such analogies would be “scientific suicide”, he says, especially given their tentative link to actual cosmology.



ROSETTA STONE

LP disc	black hole
label	interior of the black hole
playing groove	exterior of the black hole
dead wax	black hole horizon
play hole	central singularity
RPM	Angular momentum
sound speed	light speed
Mister Fahrenheit	Hawking Radiation
supersonic (wo)man	analogue experimentalist

(Don't stop me now)
(Don't stop me)
'Cause I'm having a good time,
having a good time

I'm a shooting star, leaping through the sky
Like a tiger defying the laws of gravity

I'm burnin' through the sky, yeah
Two hundred degrees
That's why they call me Mister Fahrenheit
I'm traveling at the speed of light

I wanna make a supersonic man out of you

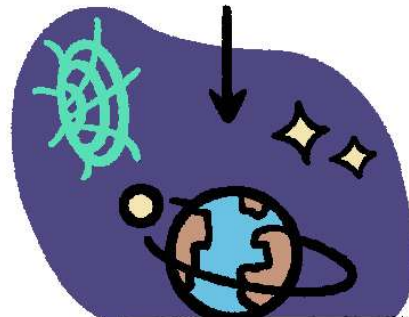


**SO,
WE PROPOSE...**

Analogue
Gravity &
Cosmology

Where...

General
Relativity



Physical
Hydrodynamics



Meet

It seems to me that the test of 'Do we or do we not understand a particular subject in Physics?' is 'Can we make a mechanical model of it?'



*Lord
Kelvin*



Alexis Bossard

Germain
Rousseaux

Nicolas James

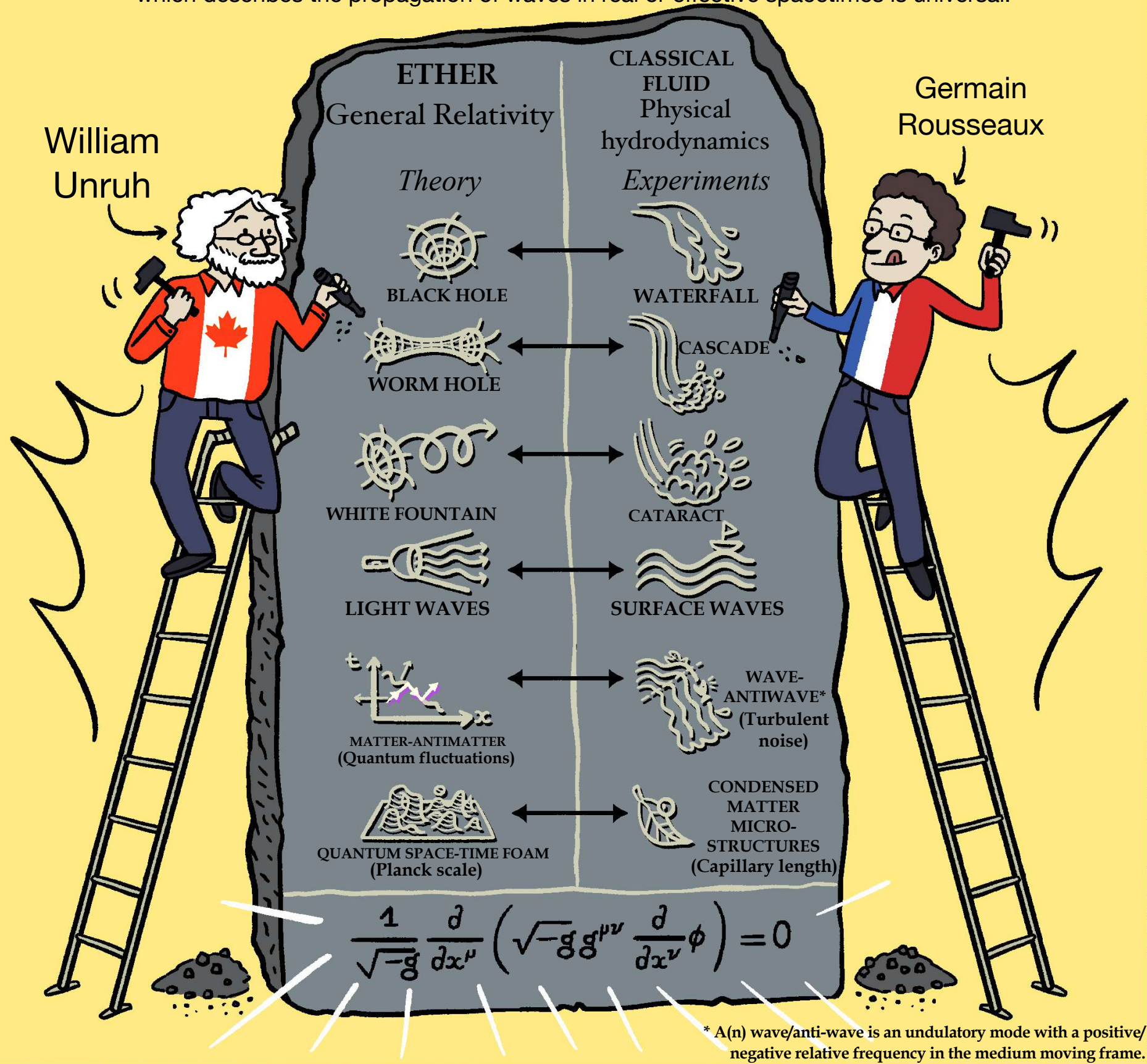
Scott Robertson

Ludivine Goncalves

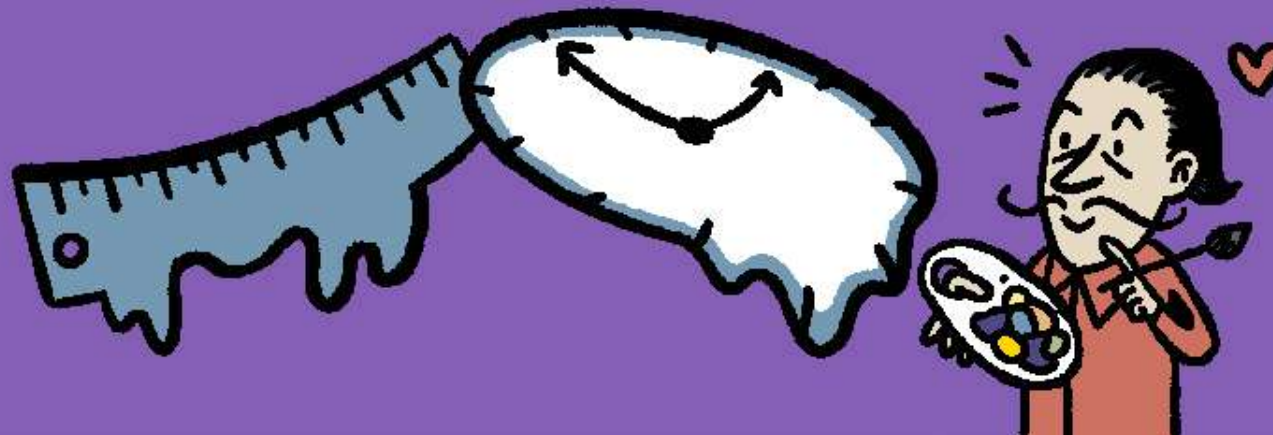
Alessia Biondi

Satadal Datta

We have been developing a Rosetta Stone which allows us to transpose questions and potential solutions to other areas that are more controllable, but not necessarily better known... The analogy with hydrodynamics allow us to approach these objects mathematically and experimentally. Our goal is to provide a maximum of correspondences, by “filling the dictionary” on the both sides of the analogy... Because the acoustic metric which describes the propagation of waves in real or effective spacetimes is universal!



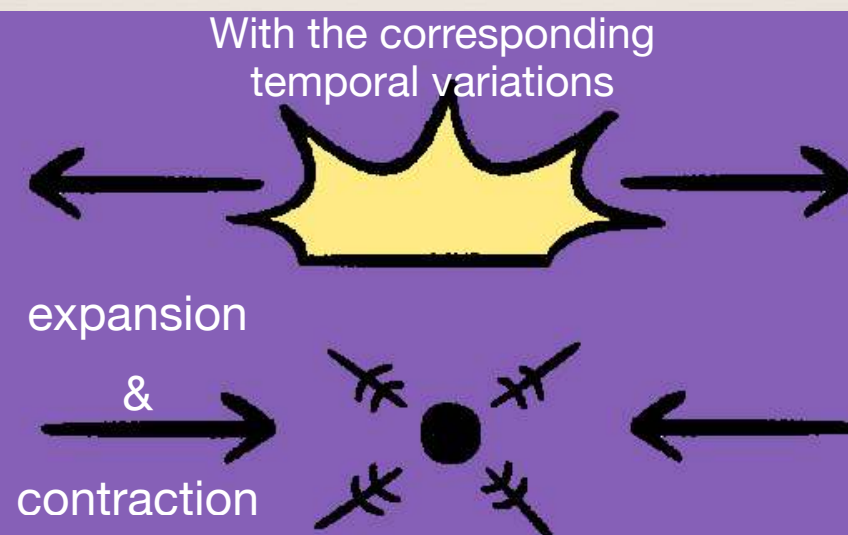
The spatio-temporal dependencies are synthesized in the concept of metric which underlies our measurement standards (rulers and clocks).



Acoustic Metric whose determinant is g

$$g_{\mu\nu} = \Omega \left(\begin{array}{c|c} -(c^2 - v^2) & -v^T \\ \hline -v & \mathbf{I} \end{array} \right)$$

c is the wave speed / v is the flow speed



Acoustic Rays as Geodesics in an Effective Space-Time (R. White 1973)

dx/dt : speed of the
wavefront

c : wave speed
(acoustic, water
waves, phonons in
a BEC)

$\mathbf{v}$: flow/medium
speed

$\hat{\mathbf{n}}$: normal vector of
norm equal to one

ds : space-time
interval

$\mathbf{u}$: column four-
vector

$\mathbf{u}^T$: line four vector

$g_{\mu\nu}$: acoustic metric

$$\frac{d\mathbf{x}}{dt} = c\hat{\mathbf{n}} + \mathbf{v}, \quad \hat{\mathbf{n}} = \frac{1}{c} \left[\frac{d\mathbf{x}}{dt} - \mathbf{v} \right]$$

$$1 = \frac{1}{c^2} \left[\frac{d\mathbf{x}}{dt} - \mathbf{v} \right] \cdot \left[\frac{d\mathbf{x}}{dt} - \mathbf{v} \right] = \frac{1}{c^2} \left[\left(\frac{d\mathbf{x}}{dt} \right)^2 - 2\mathbf{v} \cdot \frac{d\mathbf{x}}{dt} + \mathbf{v}^2 \right]$$

$$c^2 dt^2 = d\mathbf{x}^2 - 2(\mathbf{v} \cdot d\mathbf{x}) dt + \mathbf{v}^2 dt^2$$

$$-(c^2 - \mathbf{v}^2) dt^2 - 2(\mathbf{v} \cdot d\mathbf{x}) dt + d\mathbf{x} \cdot d\mathbf{x} = 0$$

$$ds^2 = g_{\mu\nu} dx^\mu dx^\nu \quad \mathbf{u}^T \equiv (dx^\mu)^T = (dt, d\mathbf{x})$$

$$(dt, d\mathbf{x}) \begin{pmatrix} -(c^2 - \mathbf{v}^2)dt - (\mathbf{v} \cdot d\mathbf{x}) \\ -\mathbf{v}dt + d\mathbf{x} \end{pmatrix} = 0$$

$$(dt, d\mathbf{x}) \begin{pmatrix} -(c^2 - \mathbf{v}^2) & -\mathbf{v}^T \\ -\mathbf{v} & \mathbb{I}_{3 \times 3} \end{pmatrix} \begin{pmatrix} dt \\ d\mathbf{x} \end{pmatrix} = 0$$

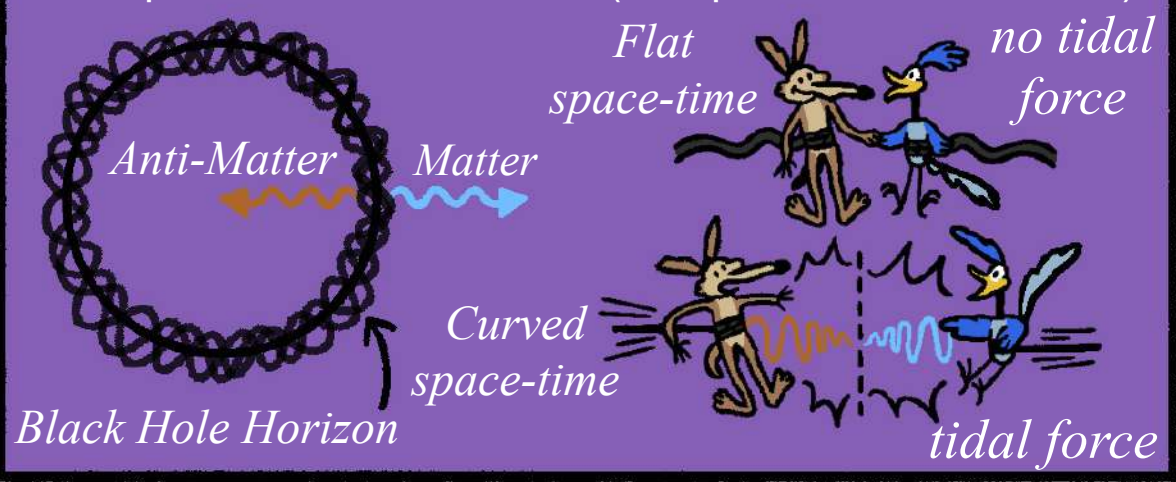
$$g_{\mu\nu} = \Omega^2 \begin{pmatrix} -(c^2 - \mathbf{v}^2) & -\mathbf{v}^T \\ -\mathbf{v} & \mathbb{I}_{3 \times 3} \end{pmatrix}$$

QUESTIONS WE
ARE WORKING ON

1

WHAT IS HAPPENING INSIDE A BLACK HOLE?

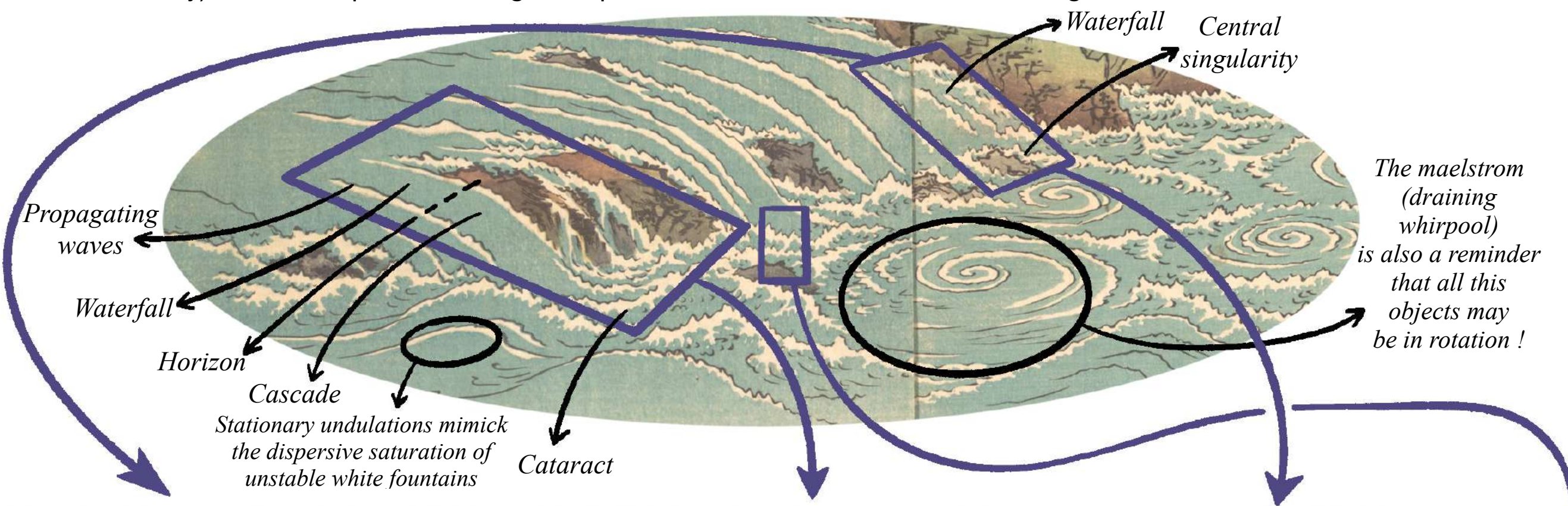
Vacuum fluctuates: virtual pairs of particles and anti-particles become real when polarized by an external field (electric/gravity field according to Schwinger/Hawking, an analogy due to Gribov). A horizon separates the pairs and the escaping mode (particle or wave) is amplified by the absorption of the anti-mode (anti-particle or anti-wave).



If we personify these particles, in a duo of linked characters, like Road Runner (BipBip in French) and Wile E. Coyote, we can develop 3 **classical** and 1 **quantum-like** scenarii regarding the fate of these particles in the vicinity of a black hole. BipBip always escapes (when the antiparticles appear, the particles extract energy of the black hole!), but what will happen to Wile E. Coyote?



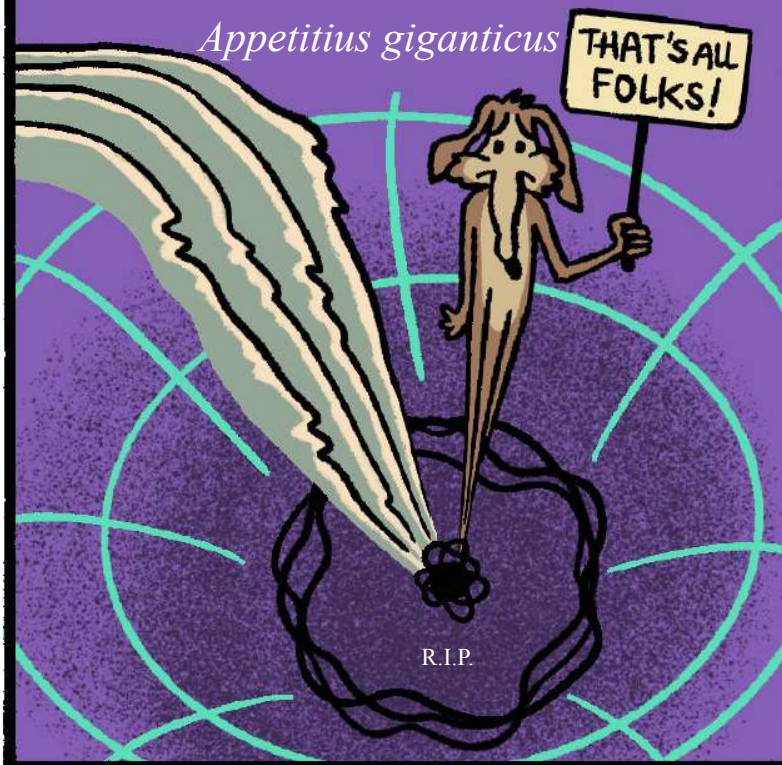
Interestingly, in this extract from this Hiroshige woodblock print (Whirlpools of Naruto Straits in Awa Province, in the 19th century), we have a perfect analogous representation of the 4 scenarii describing one's fate within a black hole...



Scenario 1:

The waterfall leads to a turbulent heart... Coyote is spaghetti-fied by the black hole during his fall towards the black hole heart.

Appetitus giganticus



Scenario 2:

The waterfall (black hole) leads to a cascade (a wormhole) and ends in a cataract (white fountain) from which Coyote comes out!



Scenario 3:

The coyote always has a stick of dynamite to blow up every thing and he may hit the heart of the black hole.



Coyote manages to “spit out” the black hole, and can come out by surfing the tidal wave created... Like a space-time tsunami!

We will try to reproduce this creation of nonlinear waves (separating Gust from Wind!) by simply closing a guillotine in the downstream side of a waterfall created by an obstacle in an open water channel.



Back-reaction (feedback mechanism of fluctuations on some background)

Trans-planckian scenario:

We can even imagine a 4th scenario happening at the microscopic scale where the capillary length plays the role of the Planck scale in order to emulate quantum gravity scenarii for which capillary waves are superluminal and may seed Hawking radiation escaping from an analogue black hole as in our reverse interstellar travel experiment without amplification by anti-matter...



A Hydraulic Black Hole in the Morbihan Gulf (France)



Gois of Berder Island, rising tide : « Pascalian » black hole by water depths mismatch.

2D Hydraulic Black Hole

big tidal range
big upstream
water depth
flood tide

Transcritical accelerating
non-dispersive waterfall



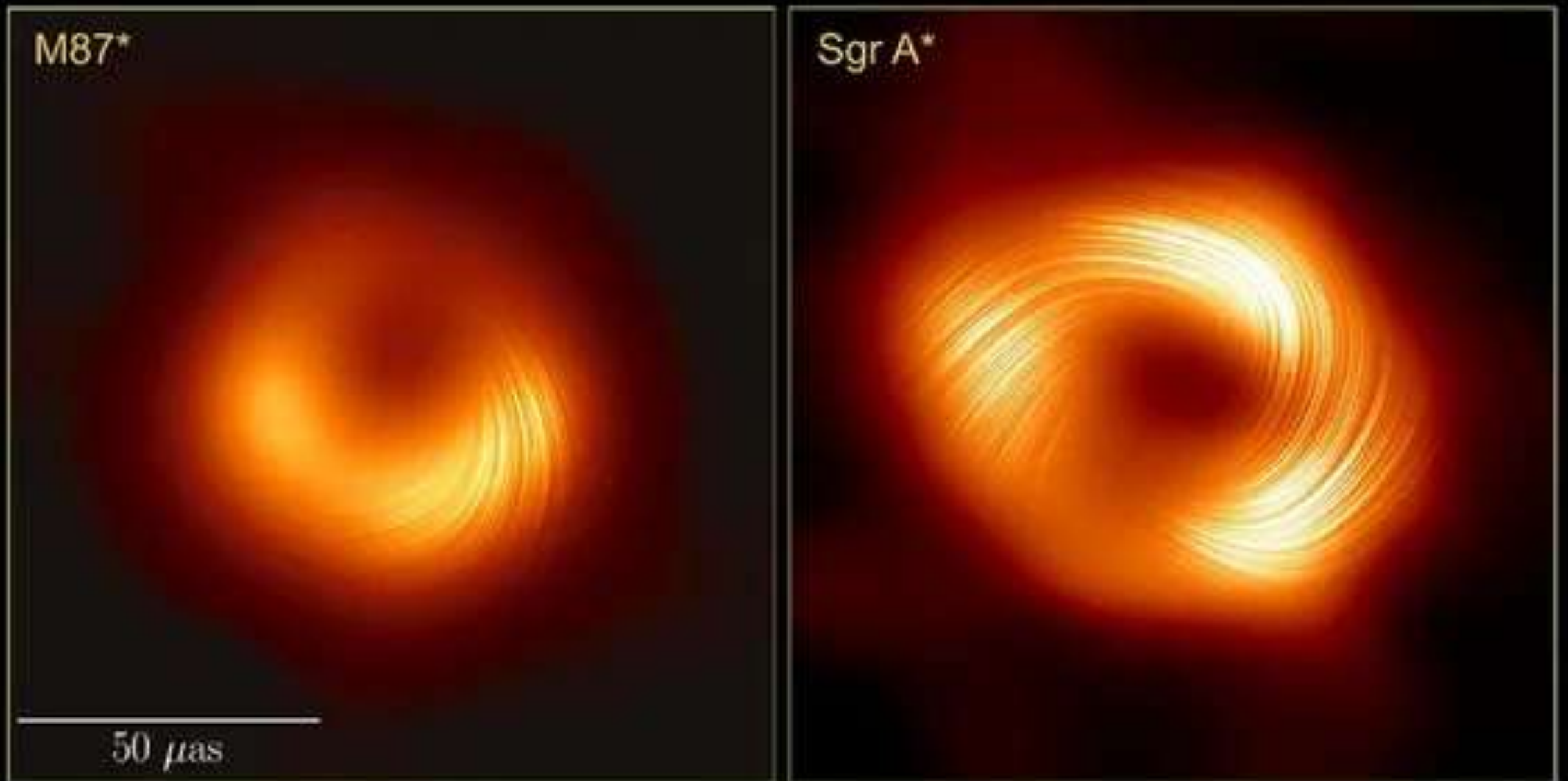
The Spaghetification in a Natural Hydraulic Black Hole



The Spaghettification in an Astrophysical Black Hole

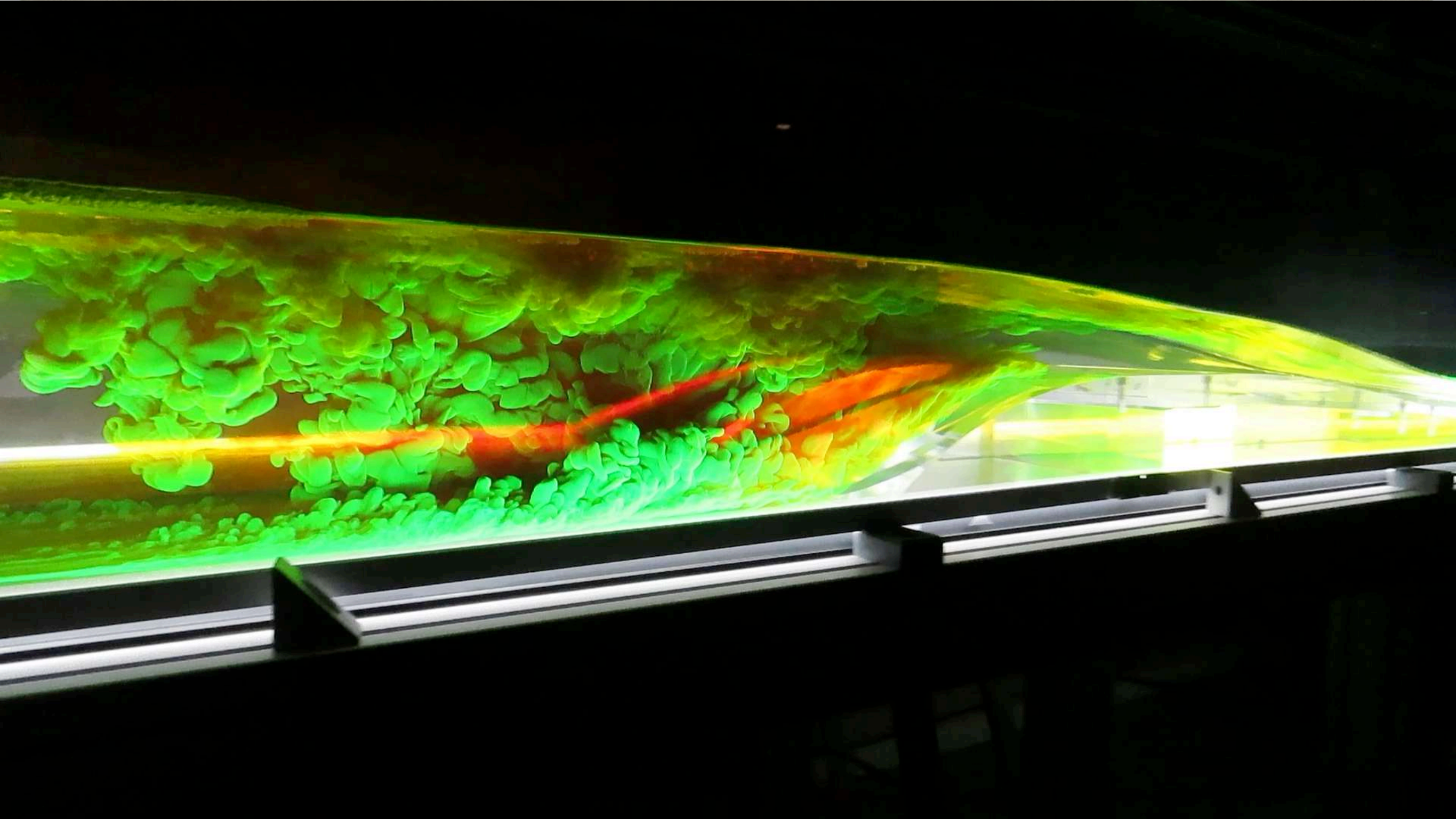


Introducing the Actual Black Holes



Courtesy Event Horizon Telescope (EHT) 2024

An Aquatic Black Hole Eating Matter

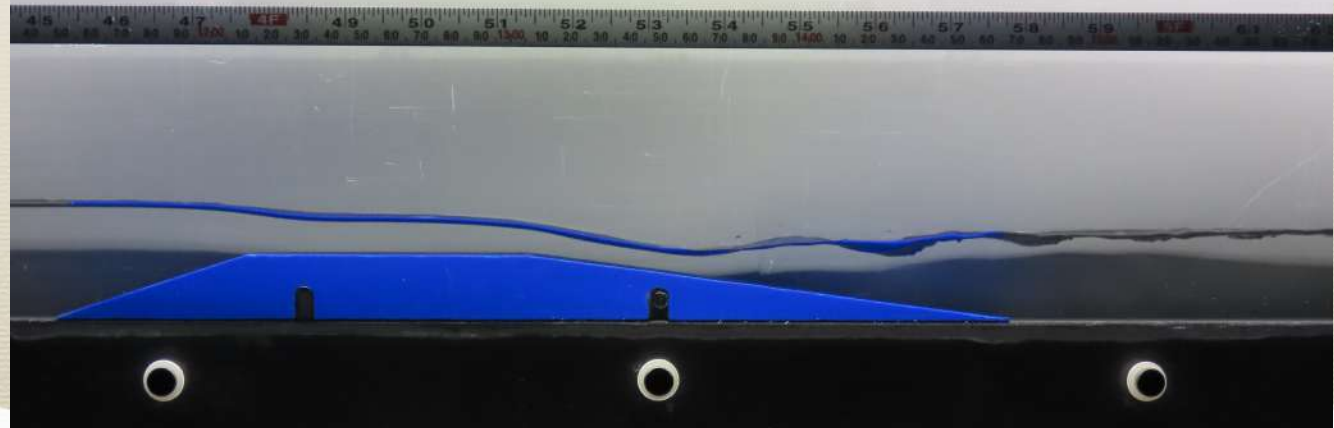
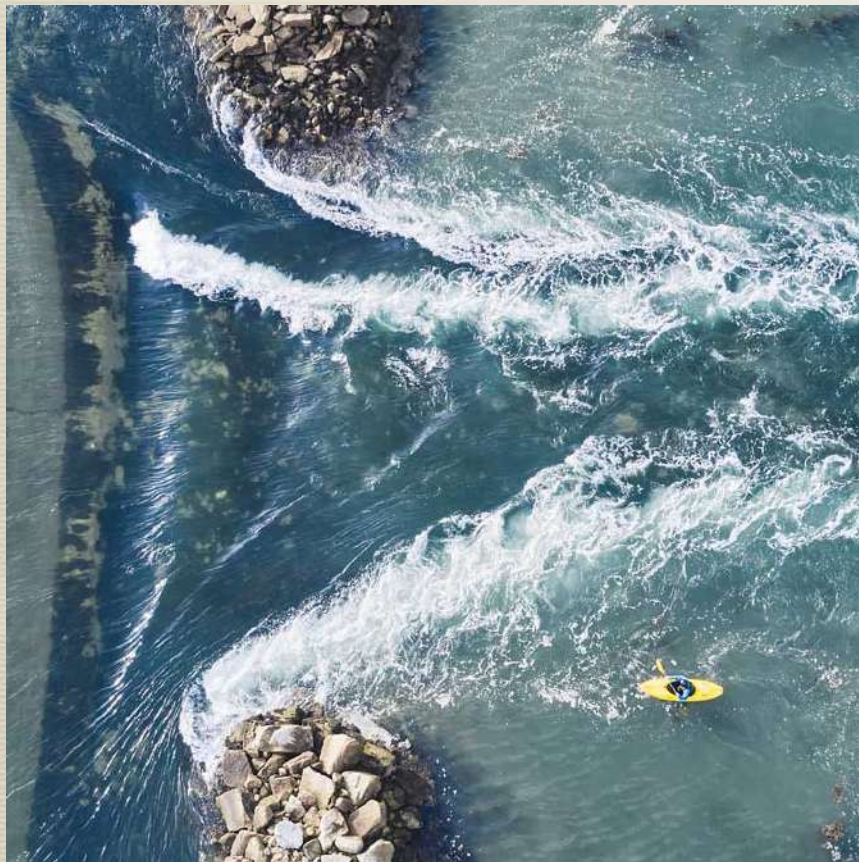


2D Hydraulic Black Hole and White Fountain

small tidal range
small upstream
water depth
flood tide

Transcritical accelerating
non-dispersive waterfall followed
by a transcritical decelerating

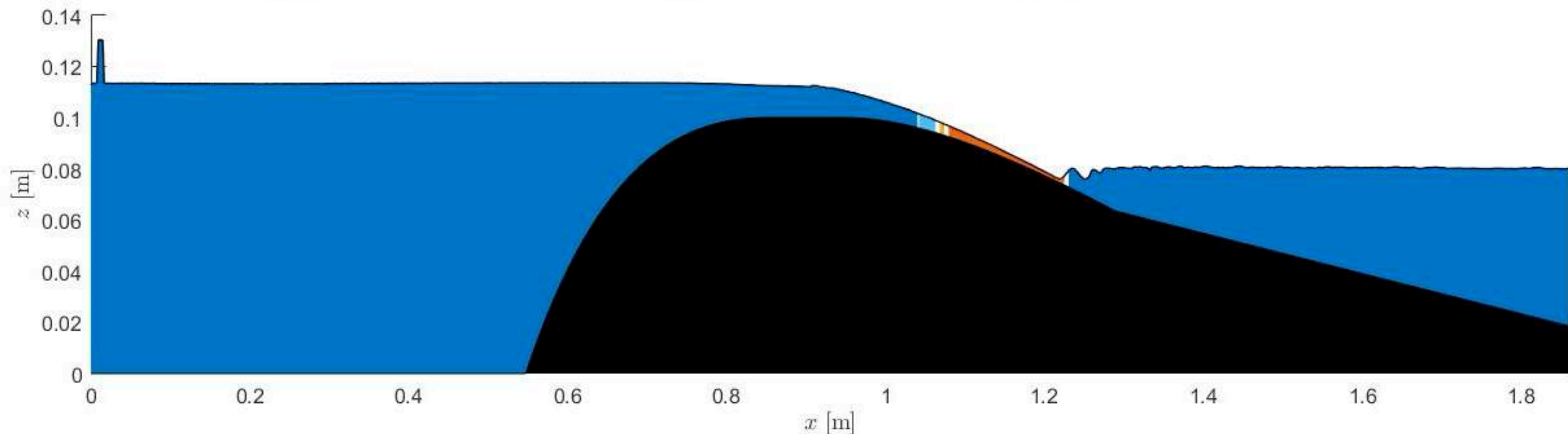
dispersive undulation



Filling up the Water Channel (or the Bathtub...)

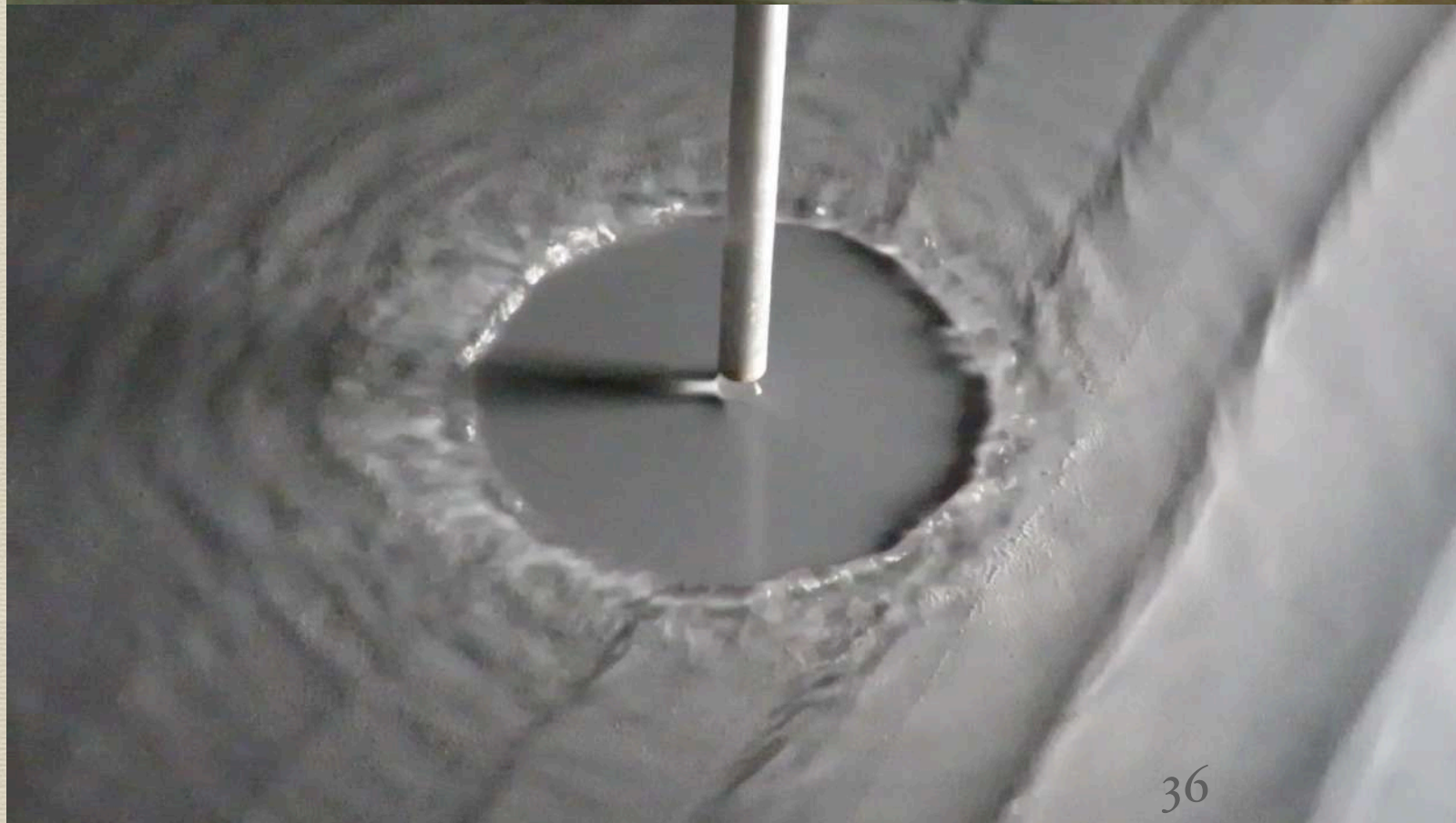
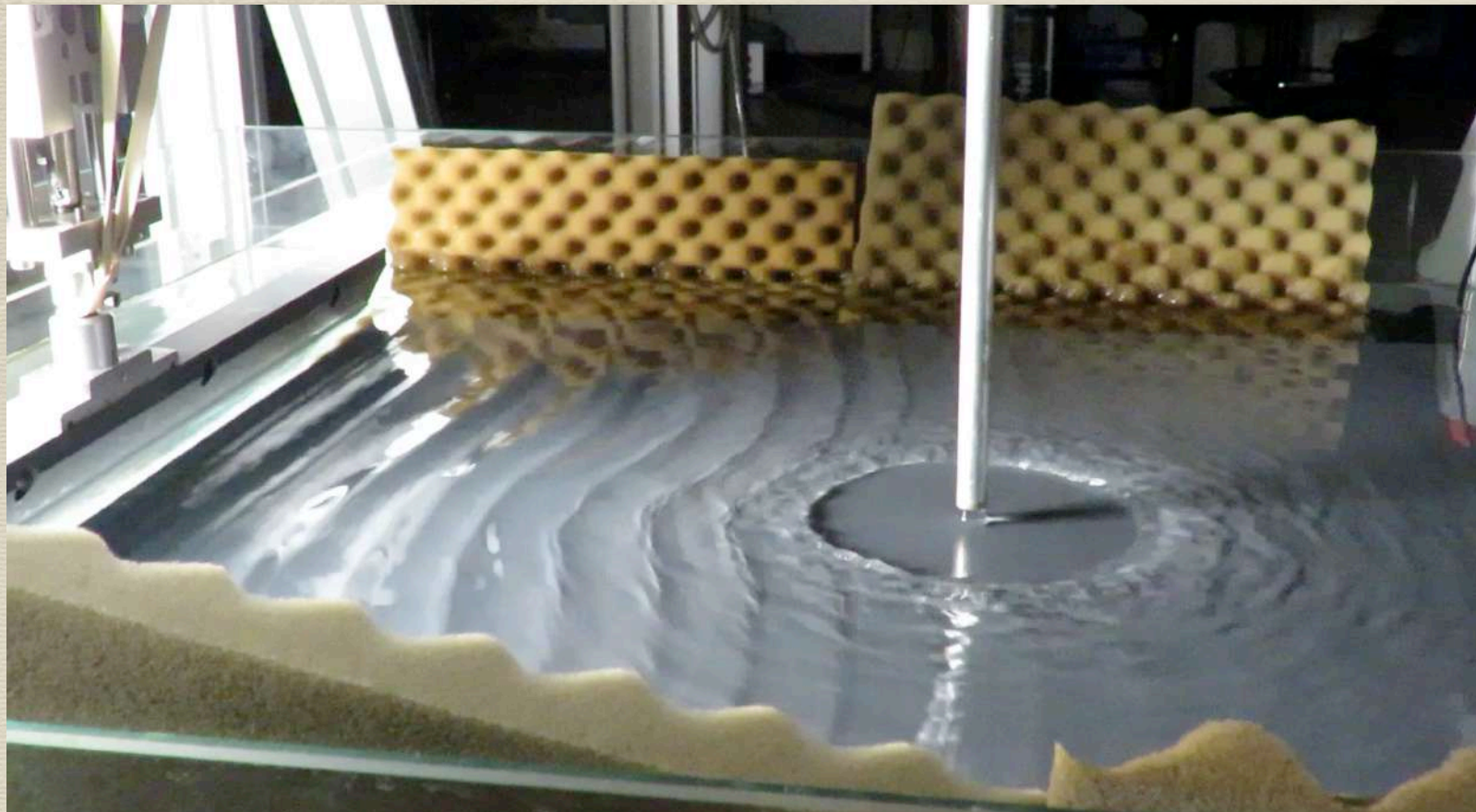
$Q = 23.13 \text{ L/min}$, $q = 1.00e - 03 \text{ m}^2/\text{s}$, $W_{\text{eff}} = 3.85e - 01 \text{ m}$, $t_{\text{acqui}} = 8.0 \text{ s}$, $f_{\text{acqui}} = 25.00\text{Hz}$, $dx = 4.660e - 04 \text{ m}$
 2 cameras (DantecSpeedSense1040), Canal HydroSedimentaire, 1 obstacle: Vancouver (0.100 m), Downstream gate : none

Fr < 1 & $U < c_{gmin} < U_L$	Fr < 1 & $c_{gmin} < U_L < U$	Fr > 1 & $c_{gmin} < U < U_L$
Fr < 1 & $c_{gmin} < U < U_L$	Fr > 1 & $U < c_{gmin} < U_L$	Fr > 1 & $c_{gmin} < U_L < U$

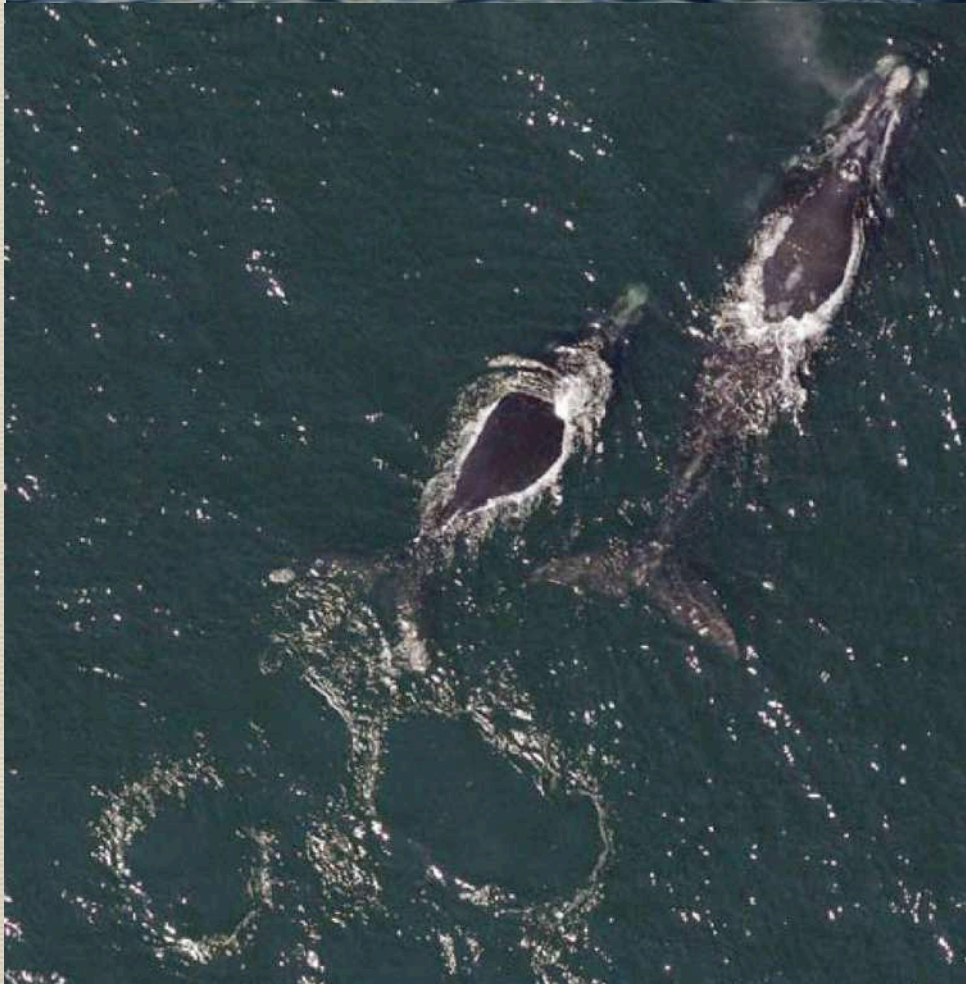


Alexis Bossard,
 PhD Thesis to be defended in
 december 2025
 Medium Water Channel of the
 Pprime Institute (France)

The Circular Jump as a 2D Hydraulic White Fountain



The Whale Flukeprint as a Dispersive White Hole Horizon (no metric)



© Robert Perry - CondorExpressPhotos.com

1D Black Hole, White Fountain and Wormhole in General Relativity

General Relativity

“The author... makes use of the mathematical analogies of the two problems to assist the imagination in the study of both.”

J. Clerk Maxwell (1861).

Point
of no return

Wormhole pinching

Blocking point

“The undiscovered country
from whose bourn
No traveller returns,
puzzles the will.”

Hamlet,
W. Shakespeare (1603).

LEGEND :



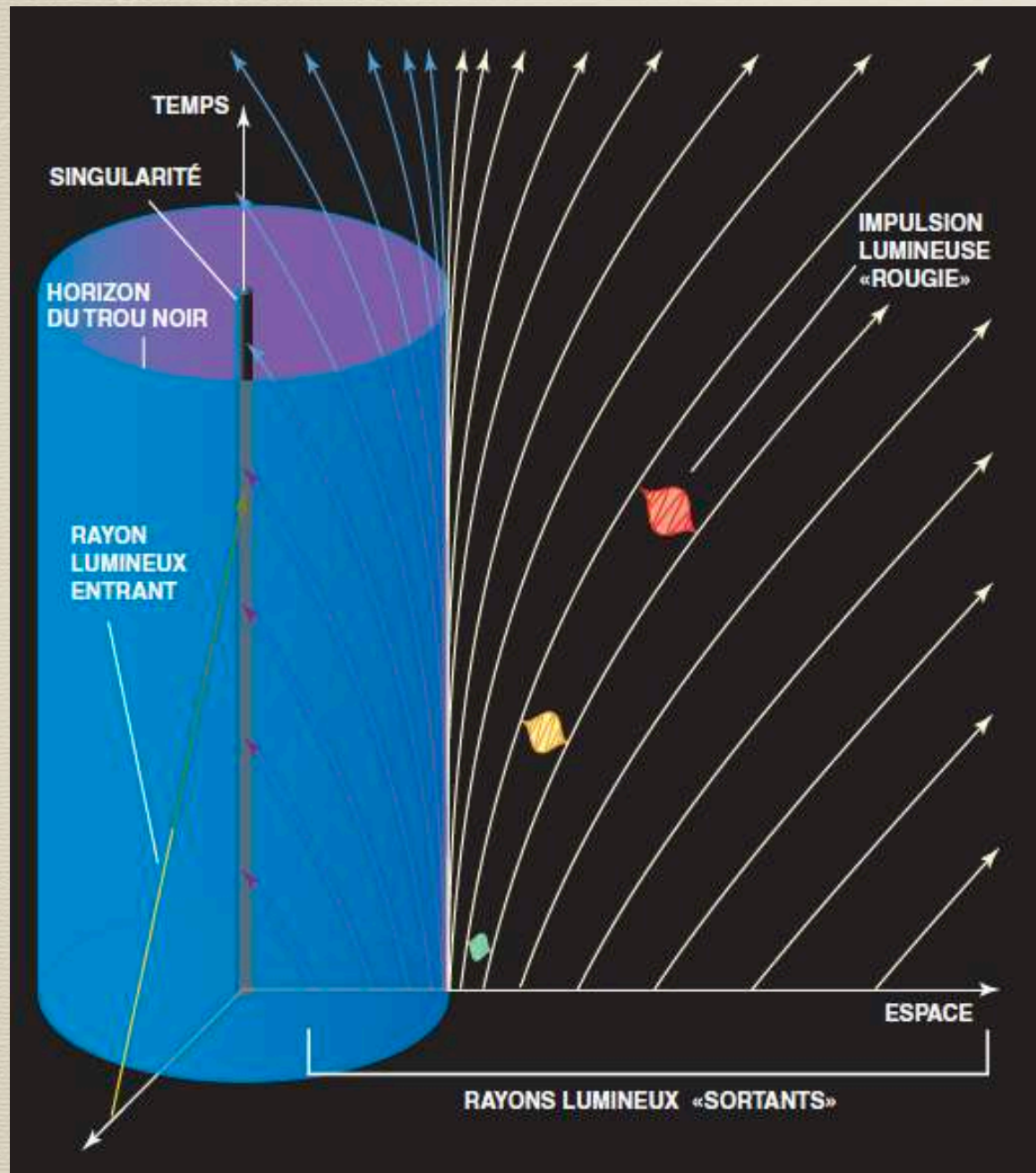
Long gravity wave



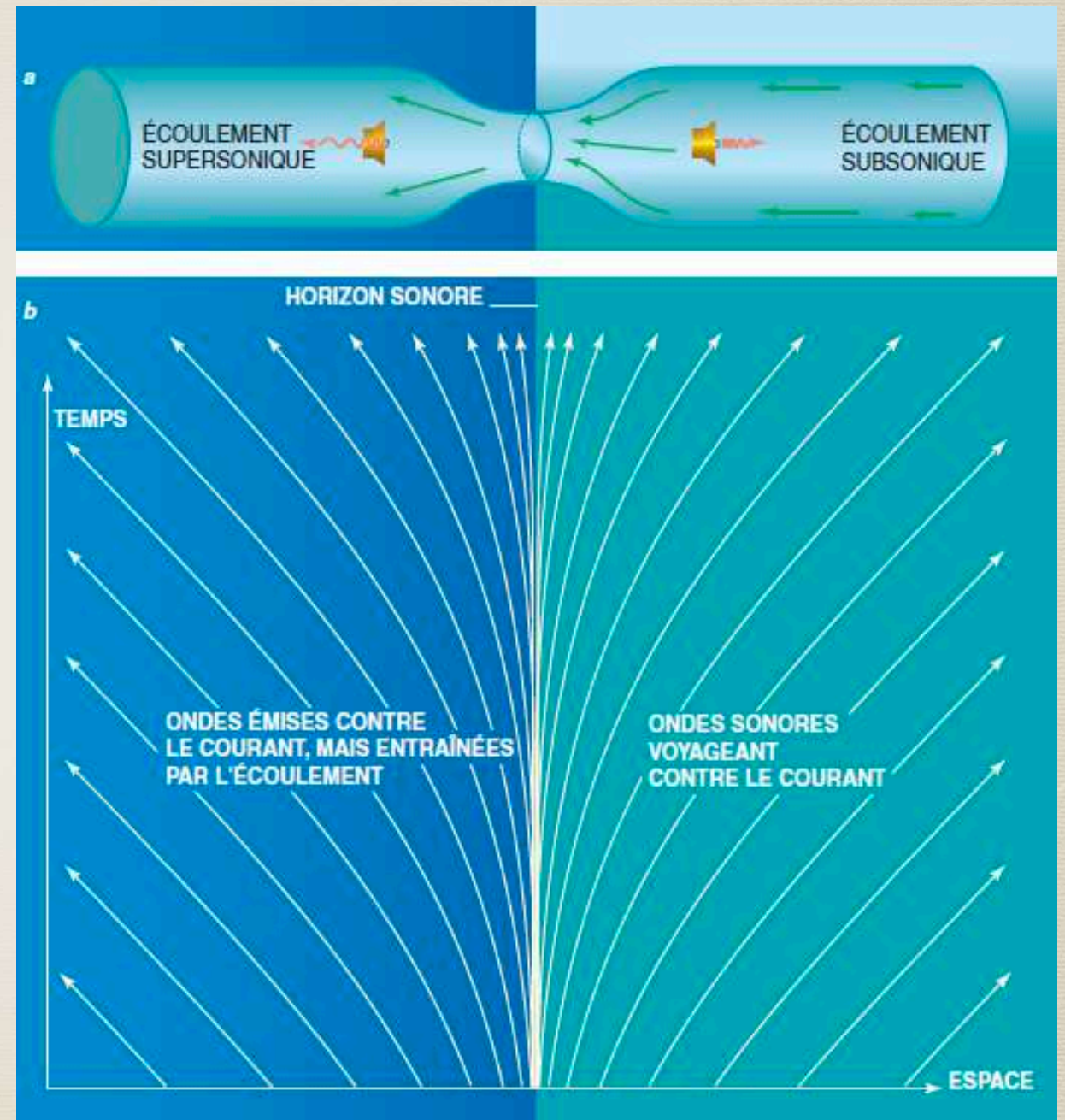
Long gravity wave

Spacetime Diagrams (a Kinematic Analogy)

A Black Hole in a P-G Spacetime



A Dumb Hole in a De Laval Nozzle



$$\frac{dx}{dt} = v(x) + c(x) \simeq \kappa x \quad \frac{dx}{\kappa x} \simeq dt \quad t \simeq t_0 + \frac{1}{\kappa} \ln\left(\frac{x}{x_0}\right) \quad \begin{array}{l} t \rightarrow \pm\infty \\ x \rightarrow 0^\pm \end{array}$$

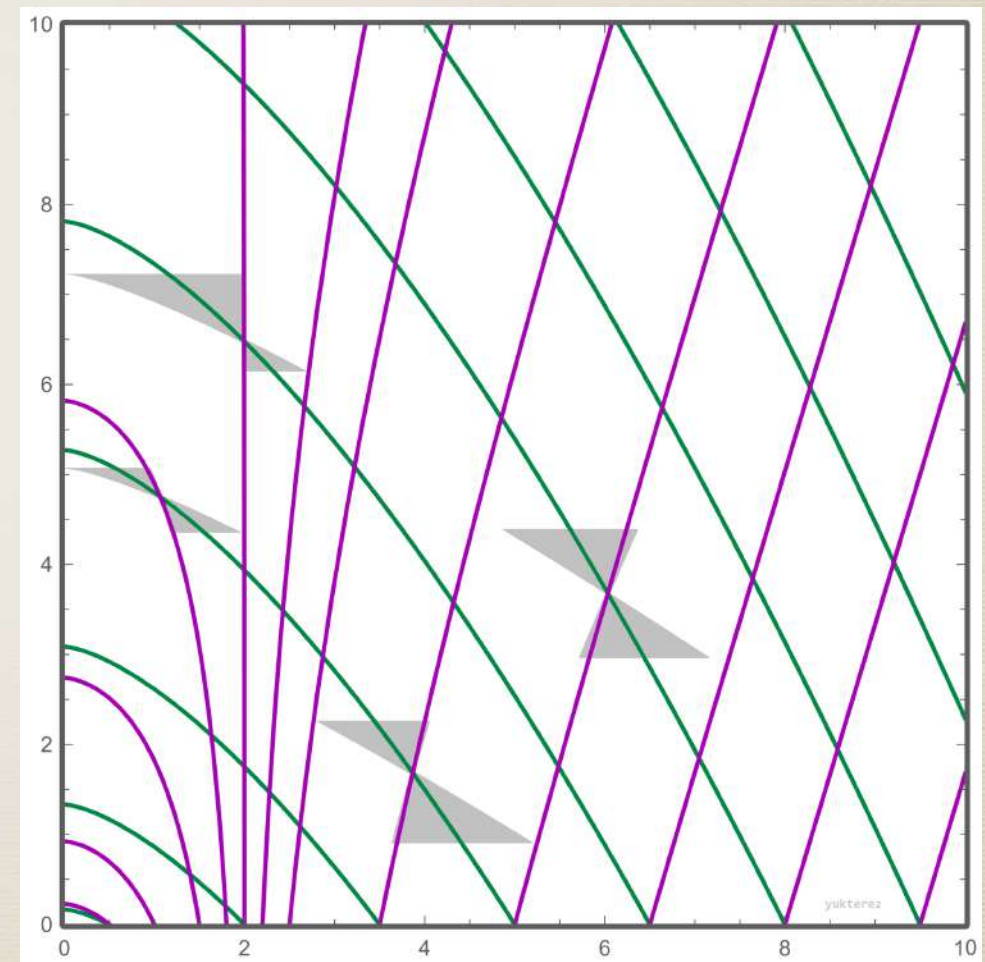
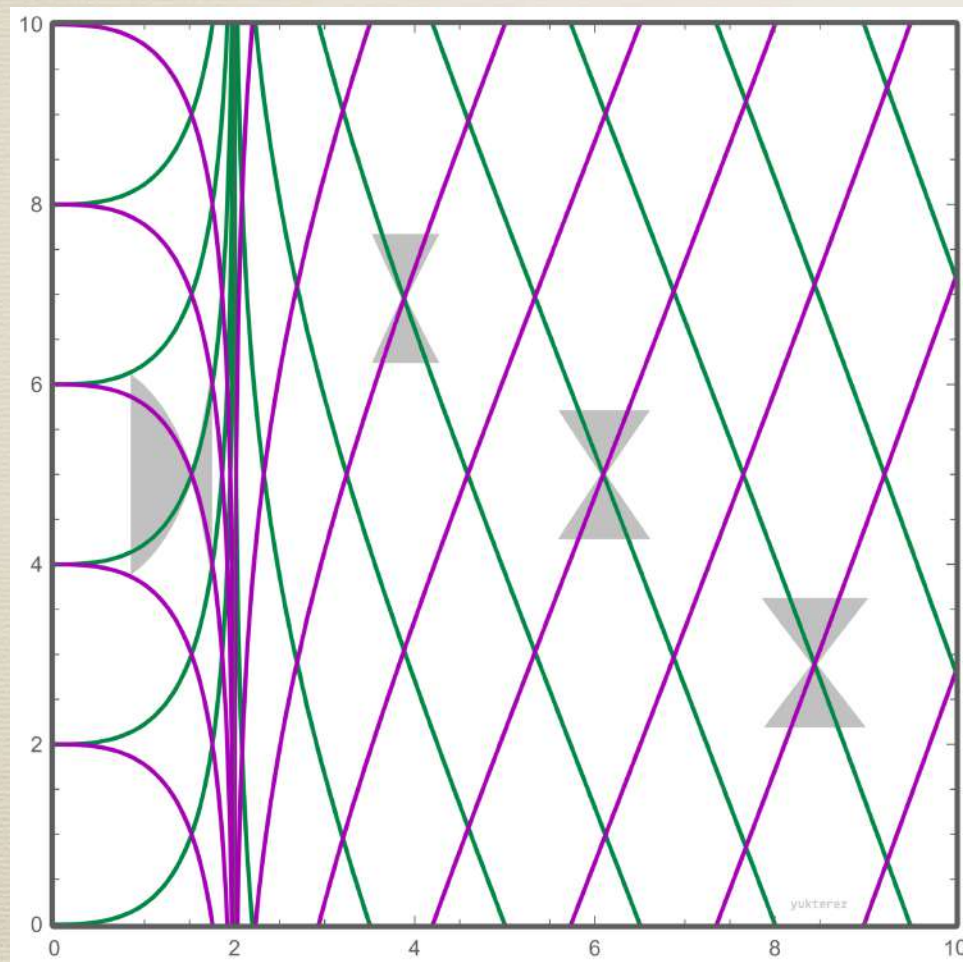
Courtesy Renaud Parentani 39 (Pour La Science, 2022).

Schwarzschild-Droste versus Painlevé-Gullstrand Space-times

« Painlevé wrote to Einstein to introduce his solution and invited Einstein to Paris for a debate. In Einstein's reply letter (December 7), he apologized for not being in a position to visit soon and explained why he was not pleased with Painlevé's arguments, emphasising that the coordinates themselves have no meaning. Finally, Einstein came to Paris in early April. On the 5th of April 1922, in a debate at the "Collège de France" with Painlevé, Becquerel, Brillouin, Cartan, De Donder, Hadamard, Langevin and Nordmann on "the infinite potentials", Einstein, baffled by the non quadratic cross term in the line element, rejected the Painlevé solution. » J. Fric

« The "trick" of the Painlevé proposal was that he no longer stuck to a full quadratic (static) form but instead, allowed a cross time-space product making the metric form no longer static but stationary and no longer direction symmetric but preferentially oriented. » J. Fric

P-G metric=acoustic metric in AG



Paul Painlevé, La mécanique classique et la théorie de la relativité, C. R. Acad. Sci. (Paris) 173, p. 677-680 (1921).

Allvar Gullstrand, Allgemeine Lösung des statischen Einkörperproblems in der Einsteinschen Gravitationstheorie, Arkiv för Matematik, Astronomi och Fysik, 16 (8), p. 1-15 (1922).

<https://astromontgeron.fr/Painleve-article-english.pdf>

First Identification of Wave Physics in an Effective Metric and of a Condensed Matter Hawking Temperature

The basic idea: consider fluid flow – Unruh (1981, 1995), Visser (1998)

Continuity:
$$\frac{\partial \rho}{\partial t} + \nabla(\rho \mathbf{v}) = 0$$

Euler's equation:
$$\frac{\partial \mathbf{v}}{\partial t} + \mathbf{v} \cdot \nabla \mathbf{v} = -\frac{\nabla p}{\rho} + \mathbf{F}$$

Assume fluid is irrotational ($\mathbf{v} = \nabla \phi$), inviscid and barotropic ($p = p(\rho)$) and linearize:

$$\rho \rightarrow \rho_0 + \rho_1 \quad \phi \rightarrow \phi_0 + \phi_1 \quad p \rightarrow p_0 + p_1$$

$$0 = -\partial_t \left(\frac{\partial \rho}{\partial p} \rho_0 (\partial_t \psi_1 + \mathbf{v}_0 \cdot \nabla \psi_1) \right) + \nabla \cdot \left(\rho_0 \nabla \psi_1 - \frac{\partial \rho}{\partial p} \rho_0 \mathbf{v}_0 (\partial_t \psi_1 + \mathbf{v}_0 \cdot \nabla \psi_1) \right)$$

First Identification of Wave Physics in an Effective Metric and of a Condensed Matter Hawking Temperature

The basic idea: consider fluid flow – Unruh (1981, 1995), Visser (1998)

1978



Continuity:

$$\frac{\partial \rho}{\partial t} + \nabla(\rho \mathbf{v}) = 0$$

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Assume fluid is irrotational ($\mathbf{v} = \nabla \phi$), inviscid and barotropic ($p = p(\rho)$) and linearize:

$$\rho \rightarrow \rho_0 + \rho_1 \quad \phi \rightarrow \phi_0 + \phi_1 \quad p \rightarrow p_0 + p_1$$

Relativistic wave equation:

$$\frac{1}{\sqrt{-g}} \frac{\partial}{\partial x^\mu} \left(\sqrt{-g} g^{\mu\nu} \frac{\partial}{\partial x^\nu} \phi_1 \right) = 0$$

with acoustic metric for massless scalar field:

where

$$g_{\mu\nu} = \frac{\rho_0}{c} \left(\begin{array}{c|c} -(c^2 - v^2) & -\mathbf{v}^T \\ \hline -\mathbf{v} & \mathbf{I} \end{array} \right) \quad g = [\det(g^{\mu\nu})]^{-1}$$



The Metric in the Dispersion Relation

Doppler shifted phonon spectrum in moving superfluid and BEC

review:

Barcelo, Liberati & Visser,
Analogue Gravity
Living Rev. Rel. 8 (2005) 12

$$E = cp + \mathbf{p} \cdot \mathbf{v}_s$$

c speed of sound
 $\mathbf{v}_s$ superfluid velocity

move $\mathbf{p} \cdot \mathbf{v}_s$ to the left

$$E - \mathbf{p} \cdot \mathbf{v}_s = cp$$

take square

$$(E - \mathbf{p} \cdot \mathbf{v}_s)^2 - c^2 p^2 = 0$$

$$g^{\mu\nu} p_\mu p_\nu = 0$$

$$p_\nu = (-E, \mathbf{p})$$

$$g^{00} = -1 \quad g^{0i} = -v_s^i \quad g^{ij} = c^2 \delta^{ij} - v_s^i v_s^j$$

effective metric

inverse metric $g_{\mu\nu}$ determines effective spacetime
in which phonons move along geodesic curves

$$ds^2 = -c^2 dt^2 + (d\mathbf{r} - \mathbf{v}_s dt)^2$$

reference frame for phonon is dragged
by moving liquid

$$ds^2 = g_{\mu\nu} dx^\mu dx^\nu$$

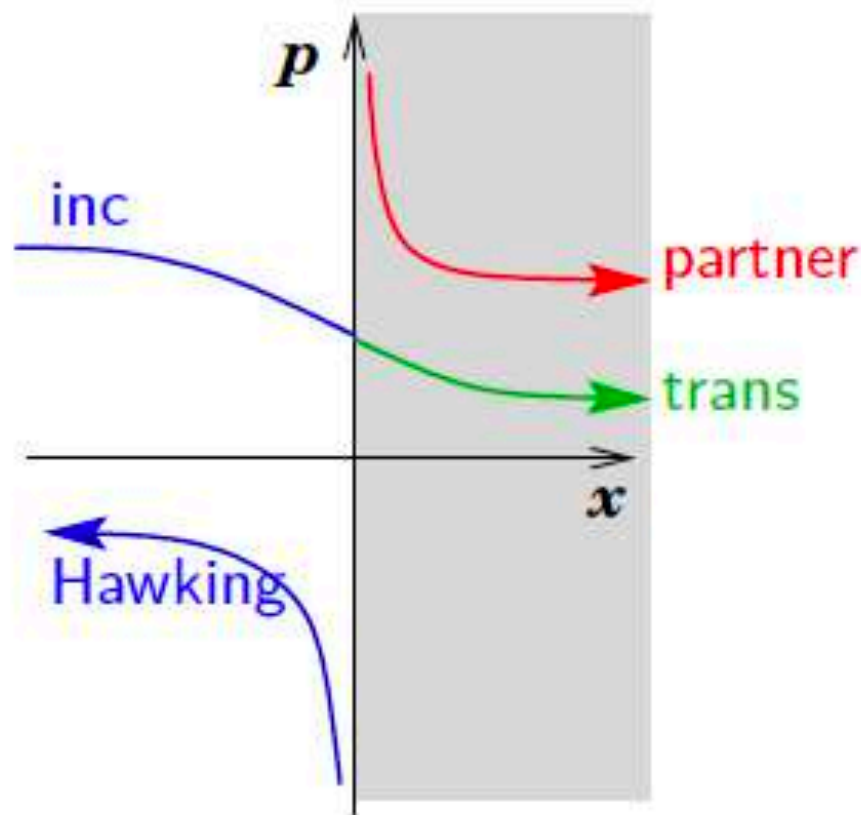
geometry is emergent

G. E. Volovik, Simulation of Panlevé-Gullstrand black hole in thin $^3\text{He-A}$ film,
Pis'ma ZhETF, 69, p. 662-668 or JETP Lett., 69, p. 705-713 (1999).

Tunneling Derivation of Hawking Temperature

long wave length limit : $E - v(x) \cdot p = \pm c(x) \cdot p \rightarrow p = \frac{E}{v(x) \pm c(x)}$

phase space :



Tunnel probability

$$P \propto e^{-2S/\hbar} \quad \text{where} \quad S = \left| \text{Im} \int p(x) dx \right|$$

near the horizon

$$\left| \frac{E}{v(x) - c(x)} \simeq \frac{E}{(x \pm i\epsilon) \frac{d}{dx}(v - c)|_0} \rightarrow \pm i\delta(x) \frac{E}{\frac{d}{dx}(v - c)|_0} \right.$$

$$\left. S \simeq \frac{\pi E}{\frac{d}{dx}(v - c)|_0} \right|$$

Hawking temperature

$$P \propto e^{-E/(k_B T_H)} \quad \text{with} \quad k_B T_H = \frac{\hbar}{2\pi} \left| \frac{d}{dx}(v - c) \right|_0$$

Surface Gravity: $\kappa = \left(\frac{d(v \pm c)}{dx} \right)_{x_H}$

Courtesy Nicolas Pavloff
44

$$k_B T_H = \frac{\hbar}{2\pi} \kappa$$

Gravity wave analogues of black holes

Ralf Schützhold* and William G. Unruh†

Department of Physics and Astronomy, University of British Columbia, Vancouver, British Columbia, Canada V6T 1Z1

(Received 22 May 2002; published 28 August 2002)

It is demonstrated that gravity waves of a flowing fluid in a shallow basin can be used to simulate phenomena around black holes in the laboratory. Since the speed of the gravity waves as well as their high-wave-number dispersion (subluminal vs superluminal) can be adjusted easily by varying the height of the fluid (and its surface tension) this scenario has certain advantages over the sonic and dielectric black hole analogs, for example, although its use in testing quantum effects is dubious. It can be used to investigate the various classical instabilities associated with black (and white) holes experimentally, including positive and negative norm mode mixing at horizons.

« For the Einsteinians, the ds^2 has a mystical and universal significance, constraining all phenomena to fit themselves in the mold of a sort of space-time form, like water in a vase. »

Paul Painlevé.

$$c = \sqrt{gh}$$

$$U(x) \Leftrightarrow V_{Schw}(r) = \sqrt{\frac{2GM}{r}}$$

$$\left(\frac{\partial}{\partial t} + \mathbf{v}_B^{\parallel} \cdot \nabla_{\parallel} \right)^2 \delta\Phi_{(0)} - gh_B \nabla_{\parallel}^2 \delta\Phi_{(0)} = 0$$

$$\square \delta\Phi_{(0)} = \frac{1}{\sqrt{-g}} \partial_{\mu} (\sqrt{-g} g^{\mu\nu} \partial_{\nu} \delta\Phi_{(0)}) = 0$$

$$g_{\text{eff}}^{\mu\nu} = \begin{pmatrix} 1 & \mathbf{v}_B^{\parallel} \\ \mathbf{v}_B^{\parallel} & \mathbf{v}_B^{\parallel} \otimes \mathbf{v}_B^{\parallel} - gh_B \mathbf{1} \end{pmatrix}$$

$$ds_{P.-G.}^2 = c^2 dt^2 - (d\vec{x} - \vec{U} dt)^2$$

PHYSICAL REVIEW D **66**, 044019 (2002)

$$kh \ll 1$$



Analogue Black Hole Flows

By courtesy of
Eric Lamballais
for the lending of
the pedagogical
open channel

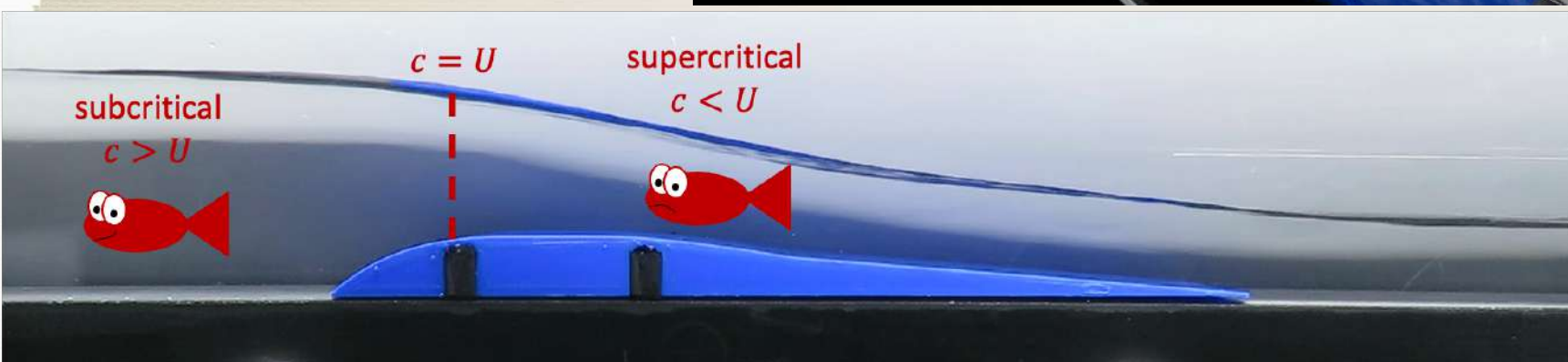


There is NO initial
water depth !

$$Fr = \frac{U}{c} = \frac{U}{\sqrt{gh}}$$

when

$$kh \ll 1$$



Small open channel features:

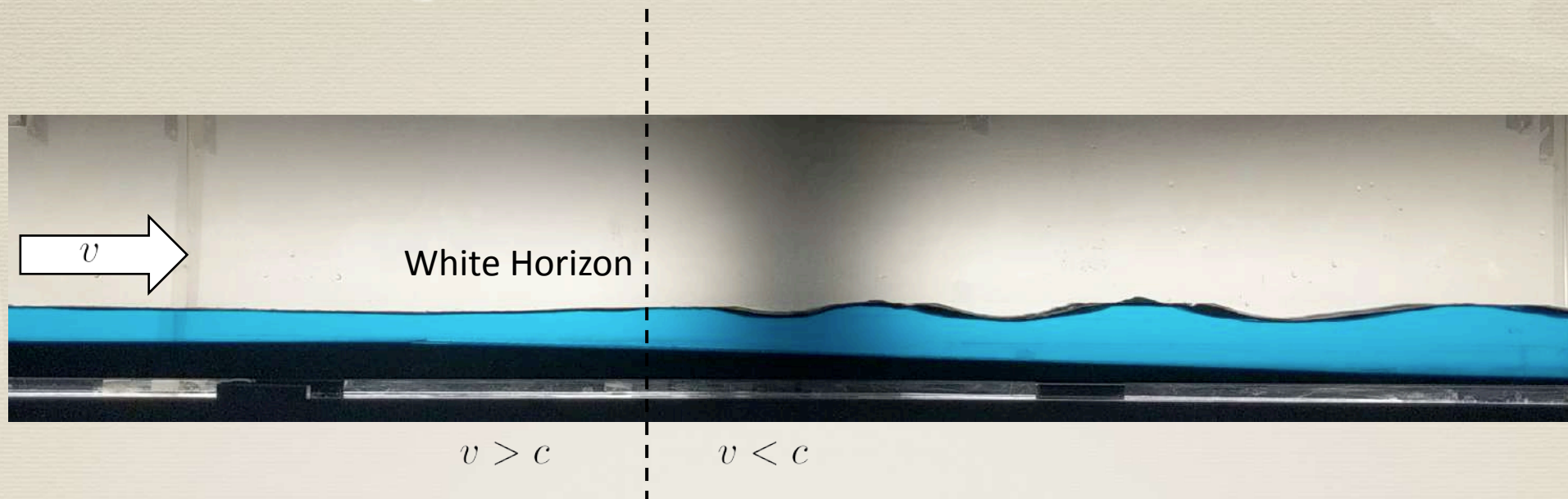
- Length : $L=2,5$ m
- Width : $w=5,4$ cm
- Flow rate range in L/min: 2-38
- Flow rate range in m^2/s : 0.0006-0.0115

Alexis Bossard,
PhD Thesis to be defended in
december 2025
Medium Water Channel of the
Pprime Institute (France)

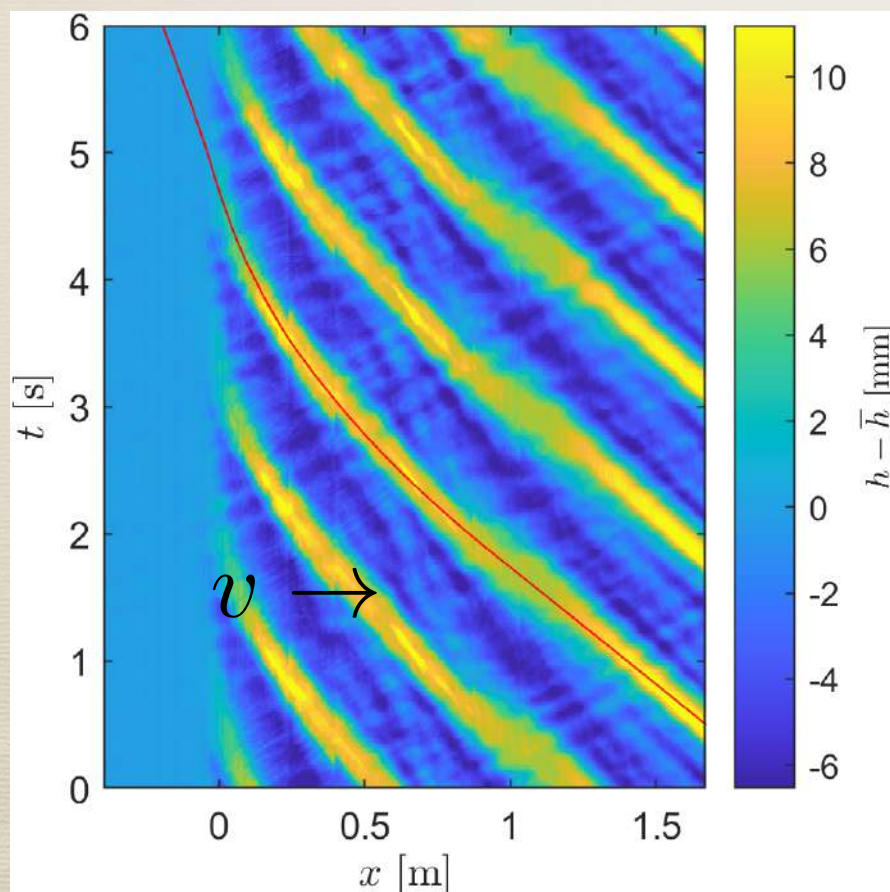


There is an initial
water depth !

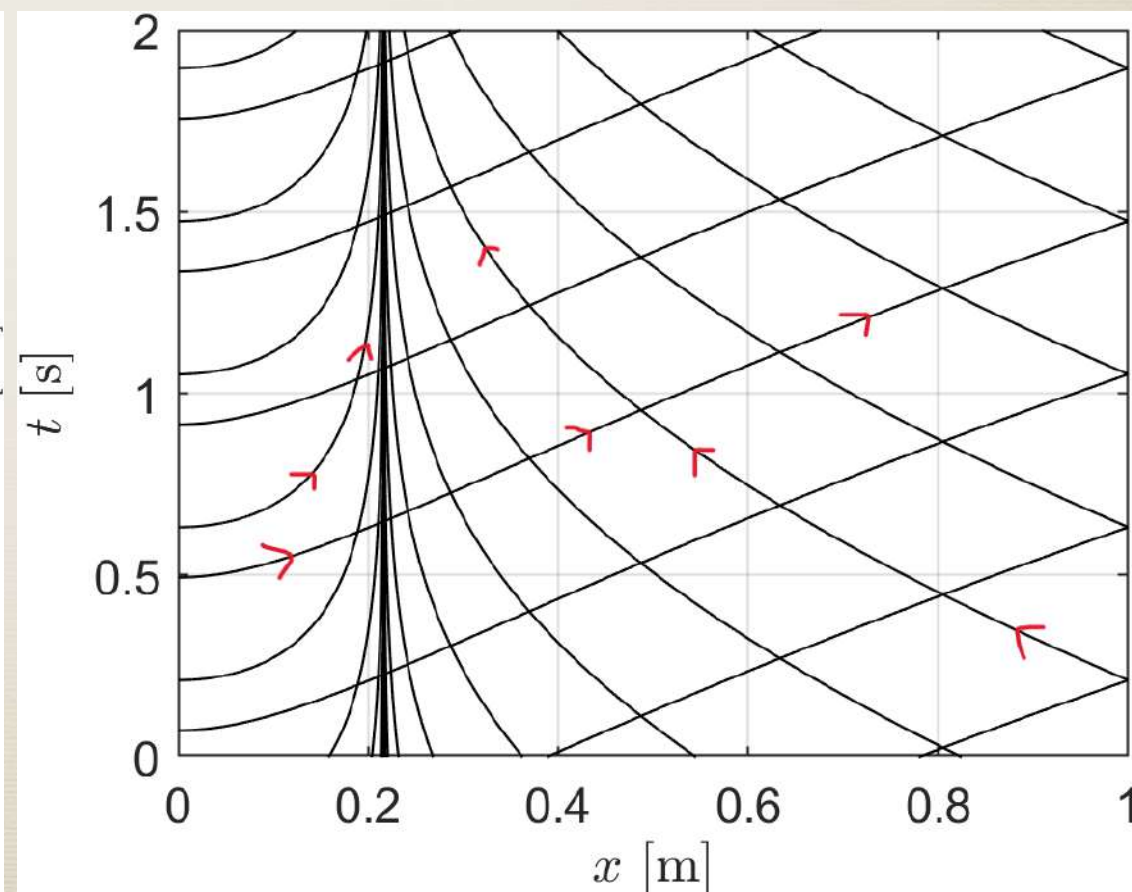
Space-Time Diagrams of a White Fountain=Cataract



Water interface recording

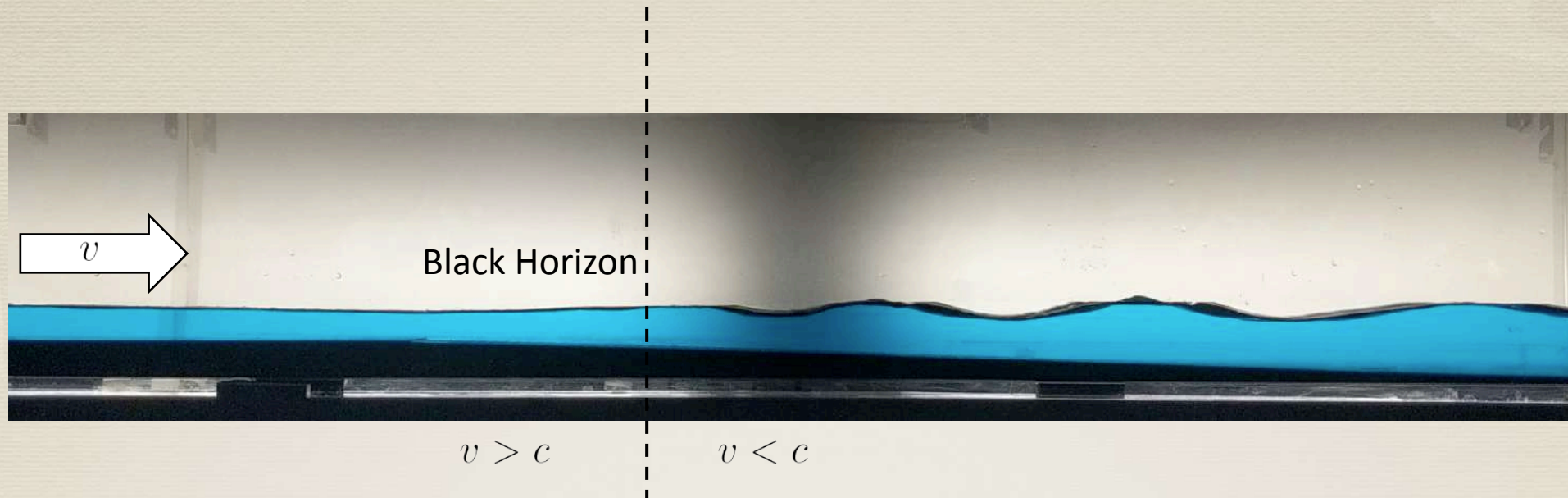


Spacetime diagram



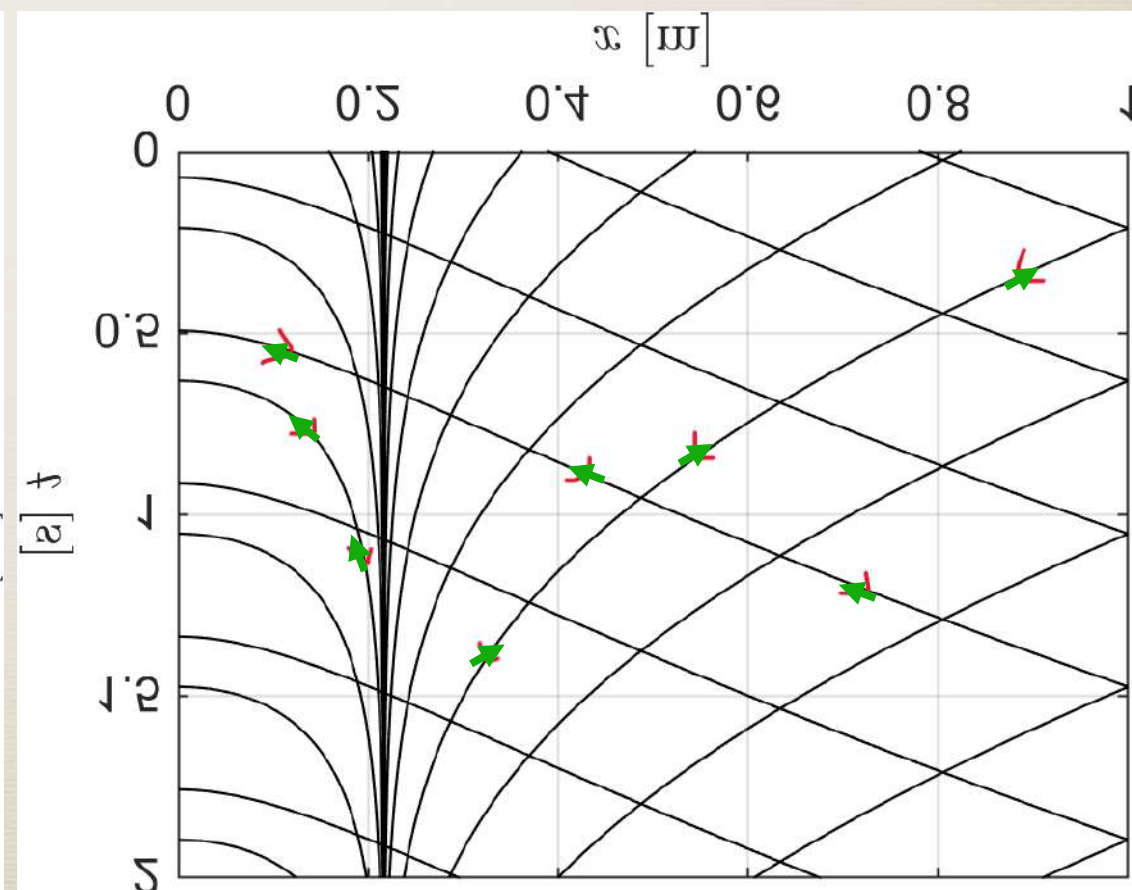
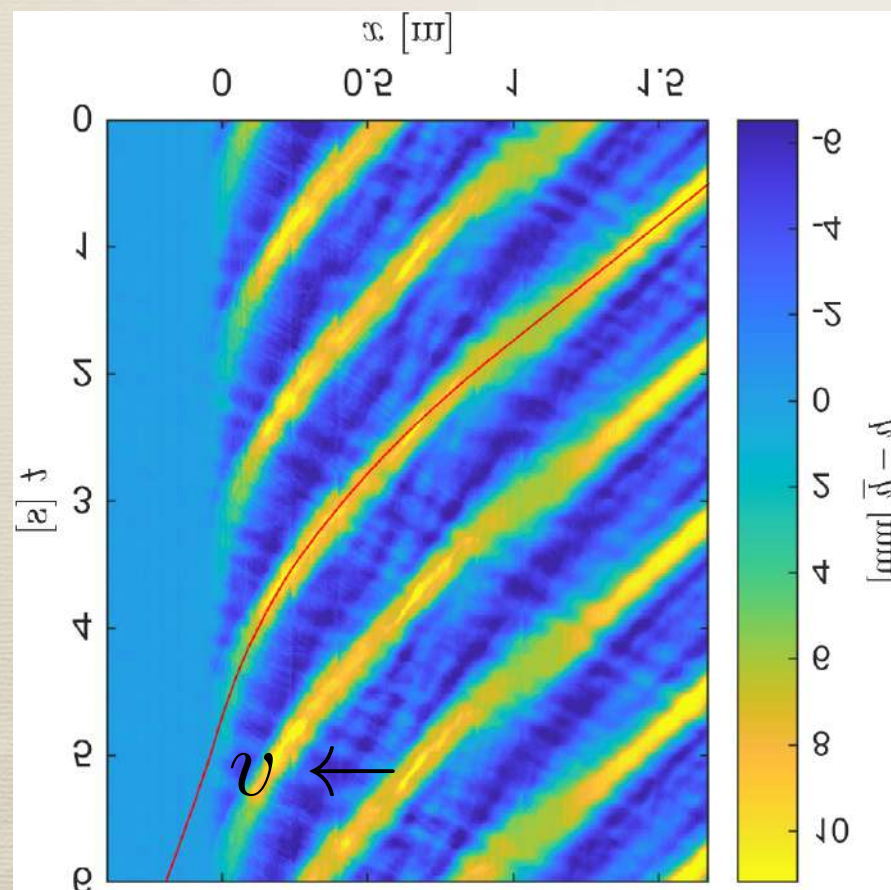
Courtesy Jaro Fransen from his Master Thesis at Eindhoven University of Technology (2024).

Space-Time Diagrams of a Black Hole=Waterfall



Water interface recording

Spacetime diagram

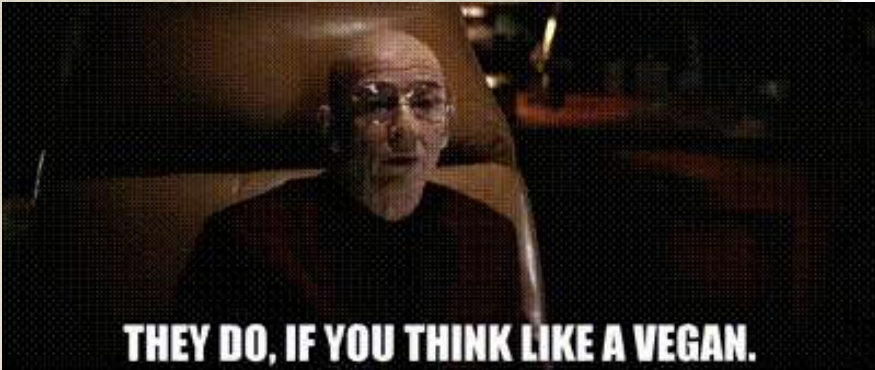
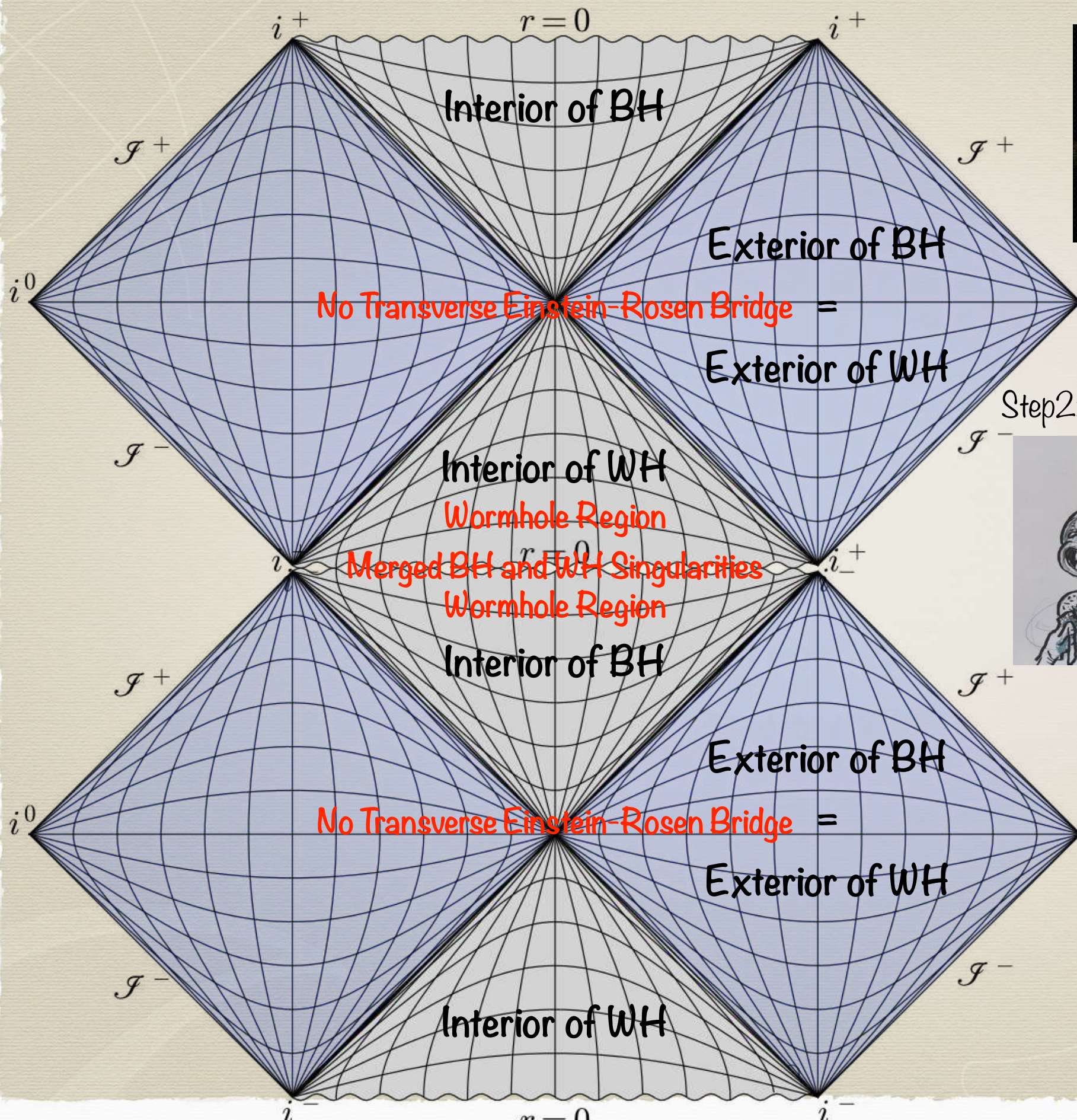


After Jaro Fransen from his Master Thesis at Eindhoven University of Technology (2024).

An Aquatic Wormhole with its Waterfall and related Cataract in Luxembourg
Pulling Pond (Schiessentümpel) of the Black Ernz River (Schwaarz lernz)

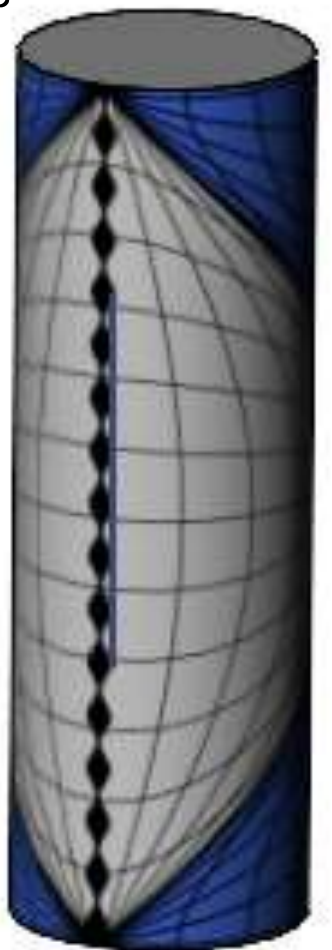


Folding the Penrose Diagram to Connect Past and Future Singularities of a NR-BH/WH pair

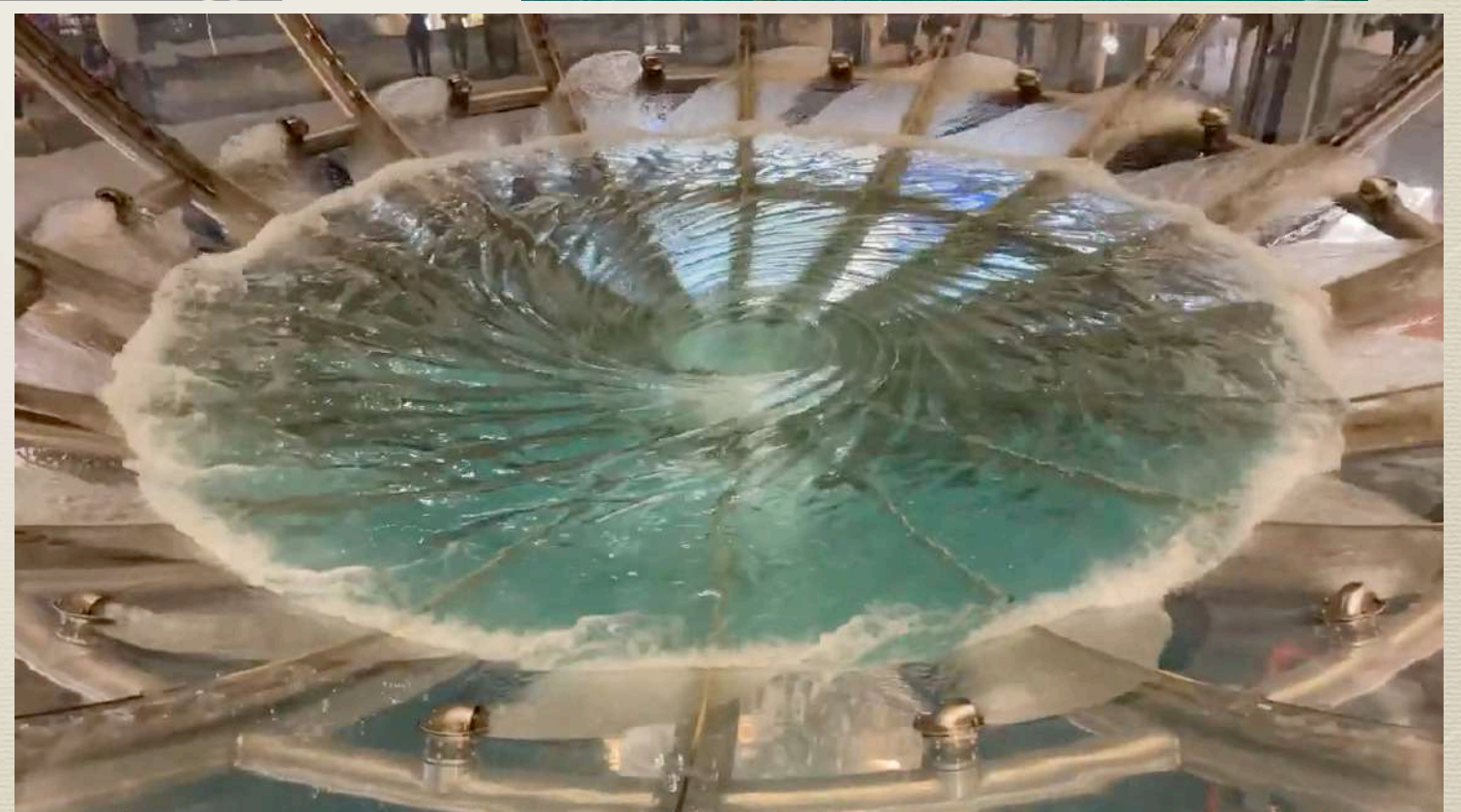
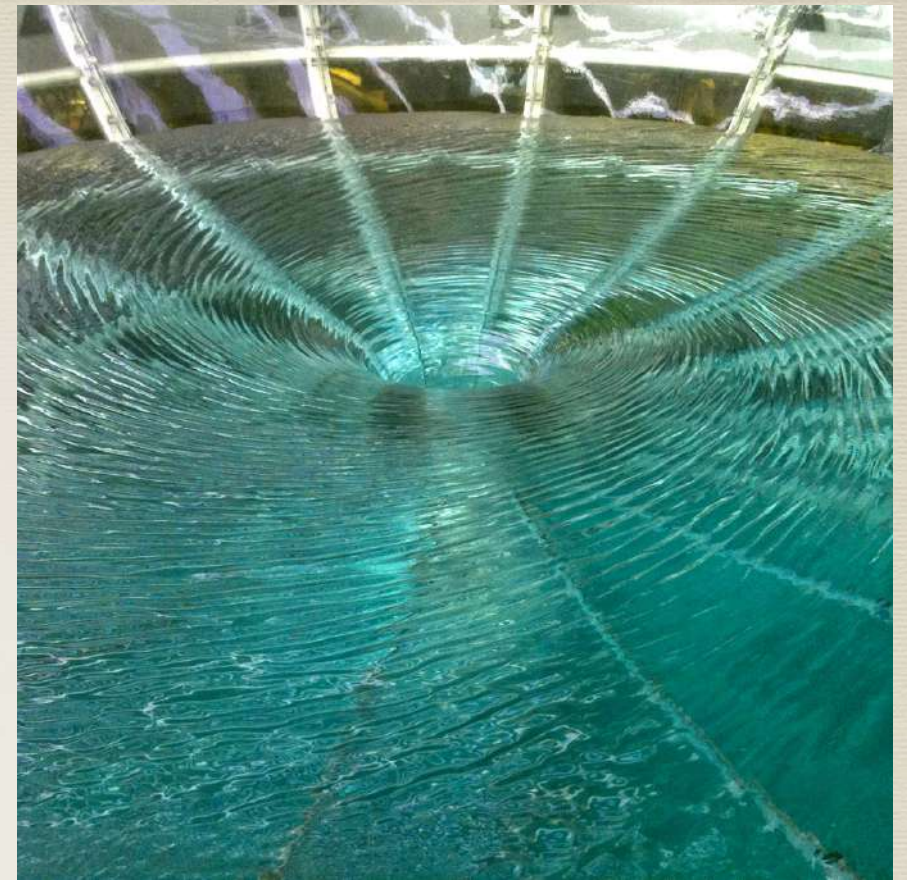
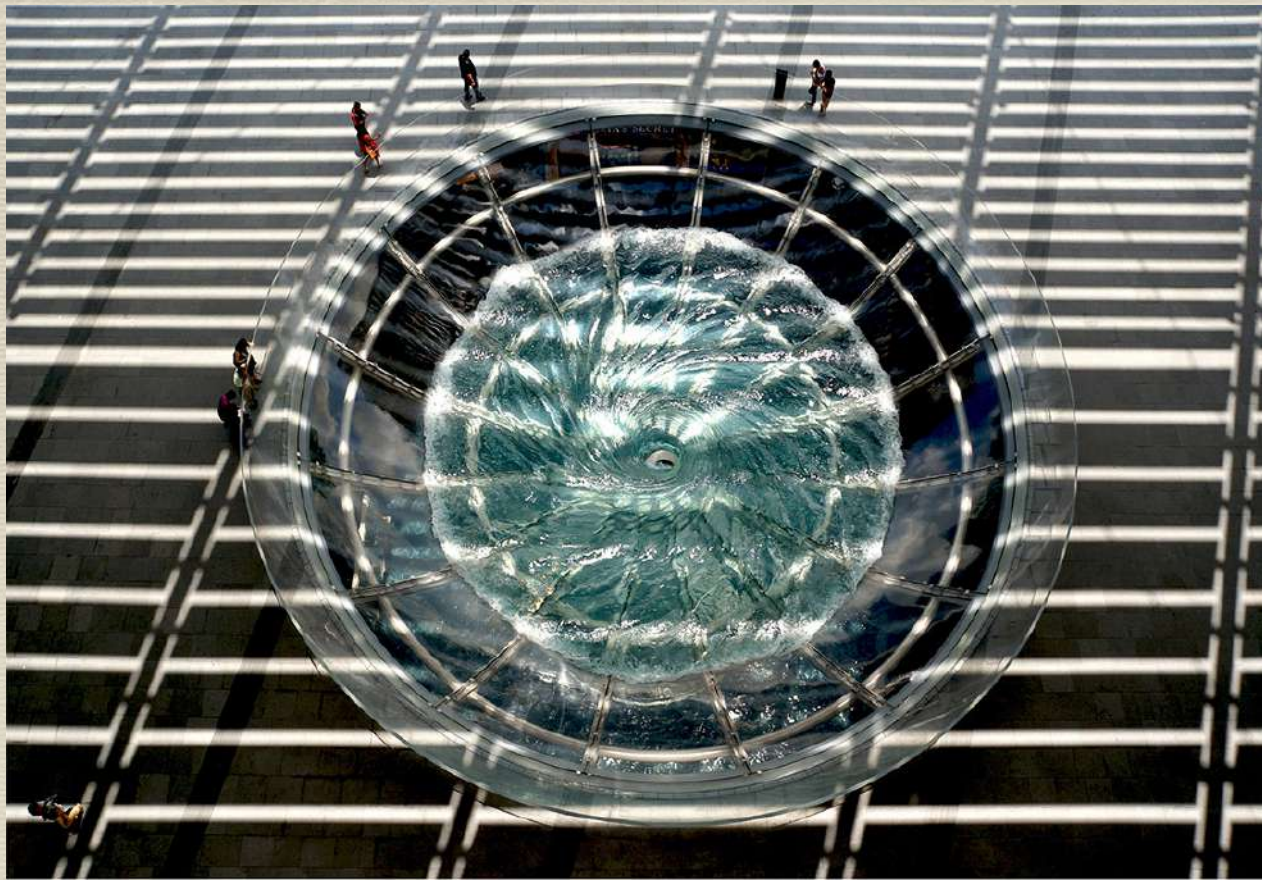


Step 1: singularities connexion

Step 2: exteriors connexion

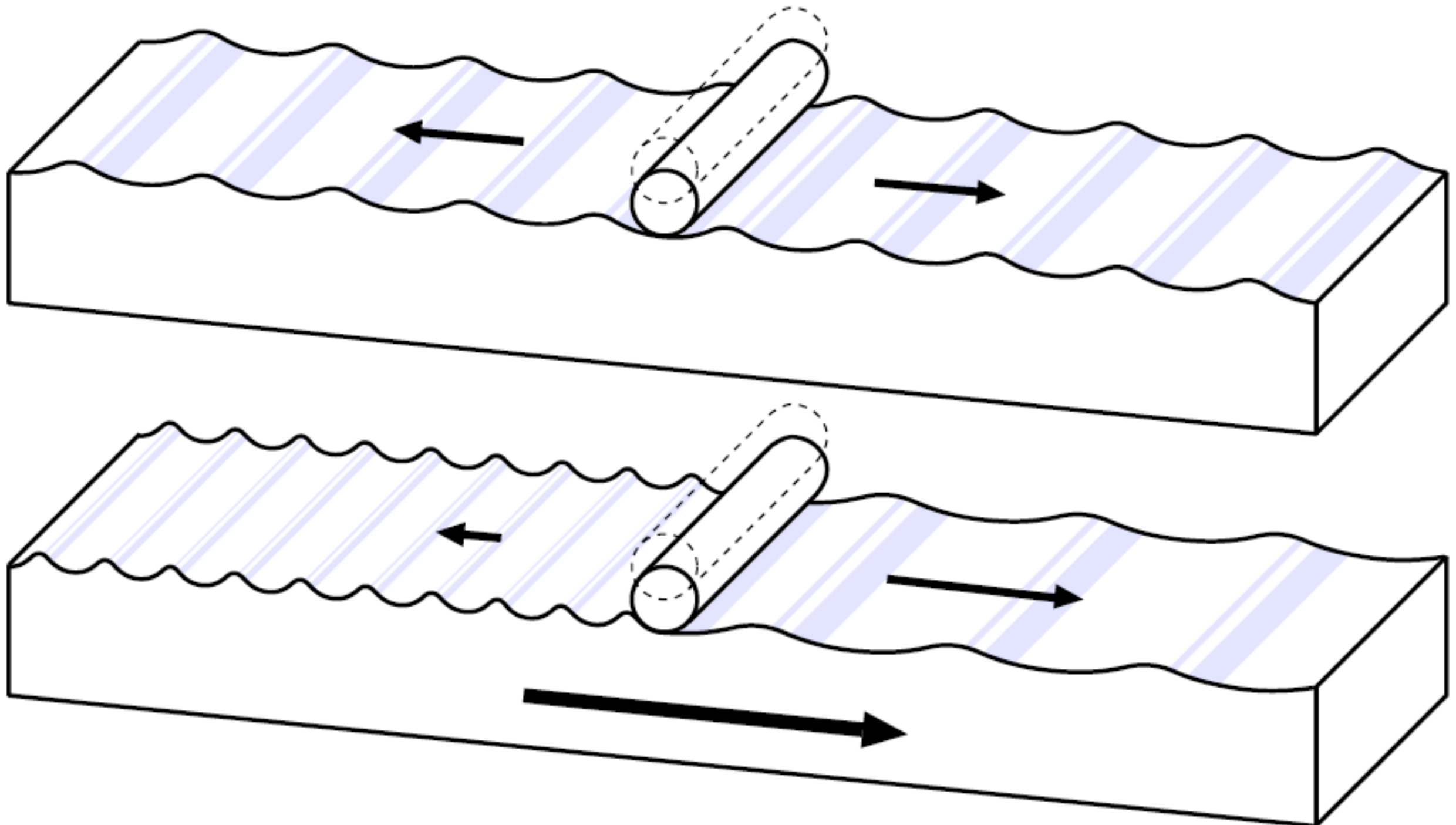


Experimentally Connected Past and Future Singularities of a R-BH



The Doppler Effect :

Blue-Shifting versus Red-Shifting



Intriguing Consequences of the Doppler Effect

$$kh \ll 1 \quad (\omega - \mathbf{U} \cdot \mathbf{k})^2 \simeq c^2 k^2 \quad c = \sqrt{gh}$$

+ For a dispersion relation with an “acoustic” branch,
the wavelength goes to zero at the blocking line

$$\omega = \Omega(k) = \mathbf{U} \cdot \mathbf{k} \pm c|k|$$

$$k_I = \frac{\omega}{U + c} \quad U < 0$$
$$U \rightarrow -c \quad \lambda_I \rightarrow 0$$

=> UV Problem at Blocking

+ Negative relative frequencies could appear

$$\omega' = \omega - \mathbf{k} \cdot \mathbf{U} < 0$$

=> Theoretical Curiosity ?

Invariants of the Massless Scalar Field Equation

$$ds^2 = [c^2 - U(x)^2]dt^2 + 2U(x)dt dx - dx^2$$

$$S = \frac{1}{2} \int dt dx \left[\frac{1}{c^2} |\partial_t \phi + U \partial_x \phi|^2 - |\partial_x \phi|^2 \right]$$

$$\partial_t(\partial_t \phi + U \partial_x \phi) + \partial_x(U \partial_t \phi + U^2 \partial_x \phi) - c^2 \partial_x^2 \phi = 0$$

$$N = \frac{i}{2c^2} \int_{-\infty}^{\infty} dx [\phi^* (\partial_t \phi + U \partial_x \phi) - \phi (\partial_t \phi^* + U \partial_x \phi^*)]$$

$\phi \rightarrow e^{i\alpha} \phi$
 $t \rightarrow t + a$

$$E = \frac{1}{2} \int_{-\infty}^{\infty} dx \left[\frac{1}{c^2} |\partial_t \phi|^2 + (1 - U^2/c^2) |\partial_x \phi|^2 \right]$$

$$U = cste$$

$$N = \frac{1}{c^2} \int_{-\infty}^{\infty} dk (\omega - Uk) |\tilde{\phi}(k)|^2$$

$$E = \frac{1}{c^2} \int_{-\infty}^{\infty} dk \omega (\omega - Uk) |\tilde{\phi}(k)|^2$$

The Wave Energy for Gravity Waves

In absence of a current :

$$E_0 = \frac{\rho g}{2} A^2 = \frac{\rho \omega^2}{2k \tanh(kh)} A^2$$

In presence of a current :

$$E_U = \frac{\rho \omega (\omega - \mathbf{U} \cdot \mathbf{k})}{2k \tanh(kh)} A^2$$

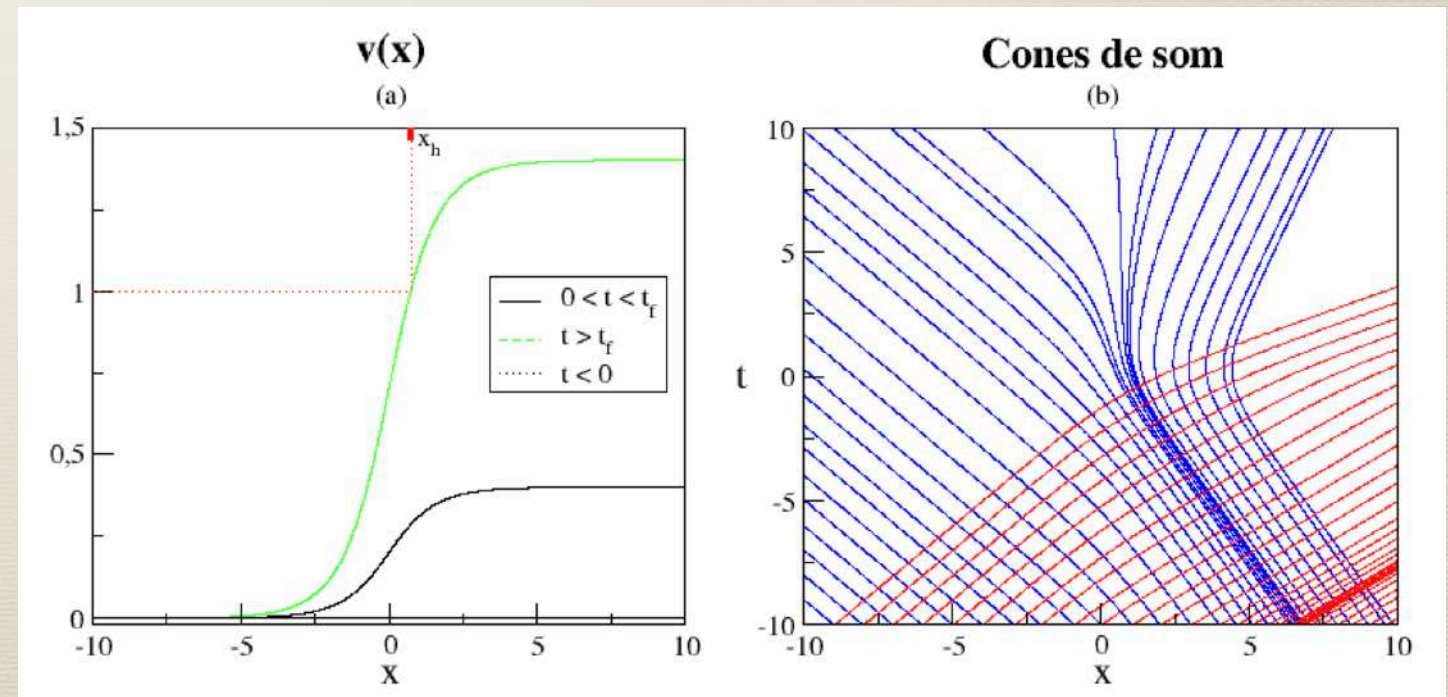
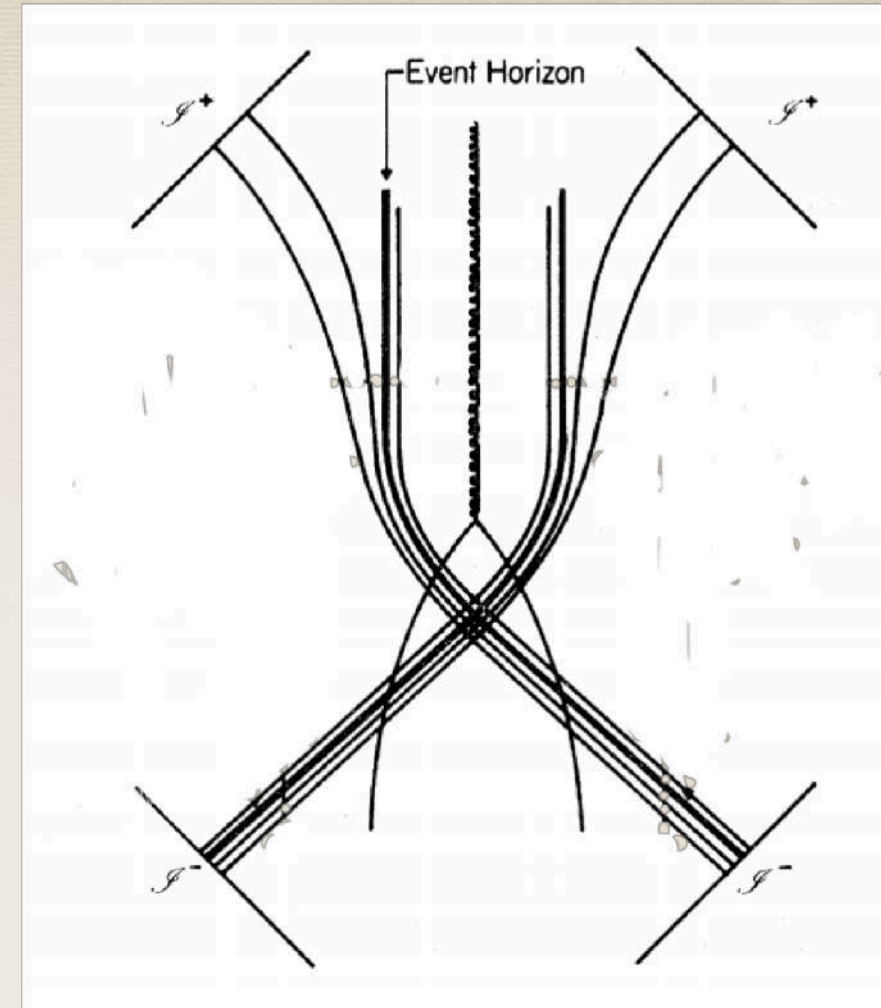
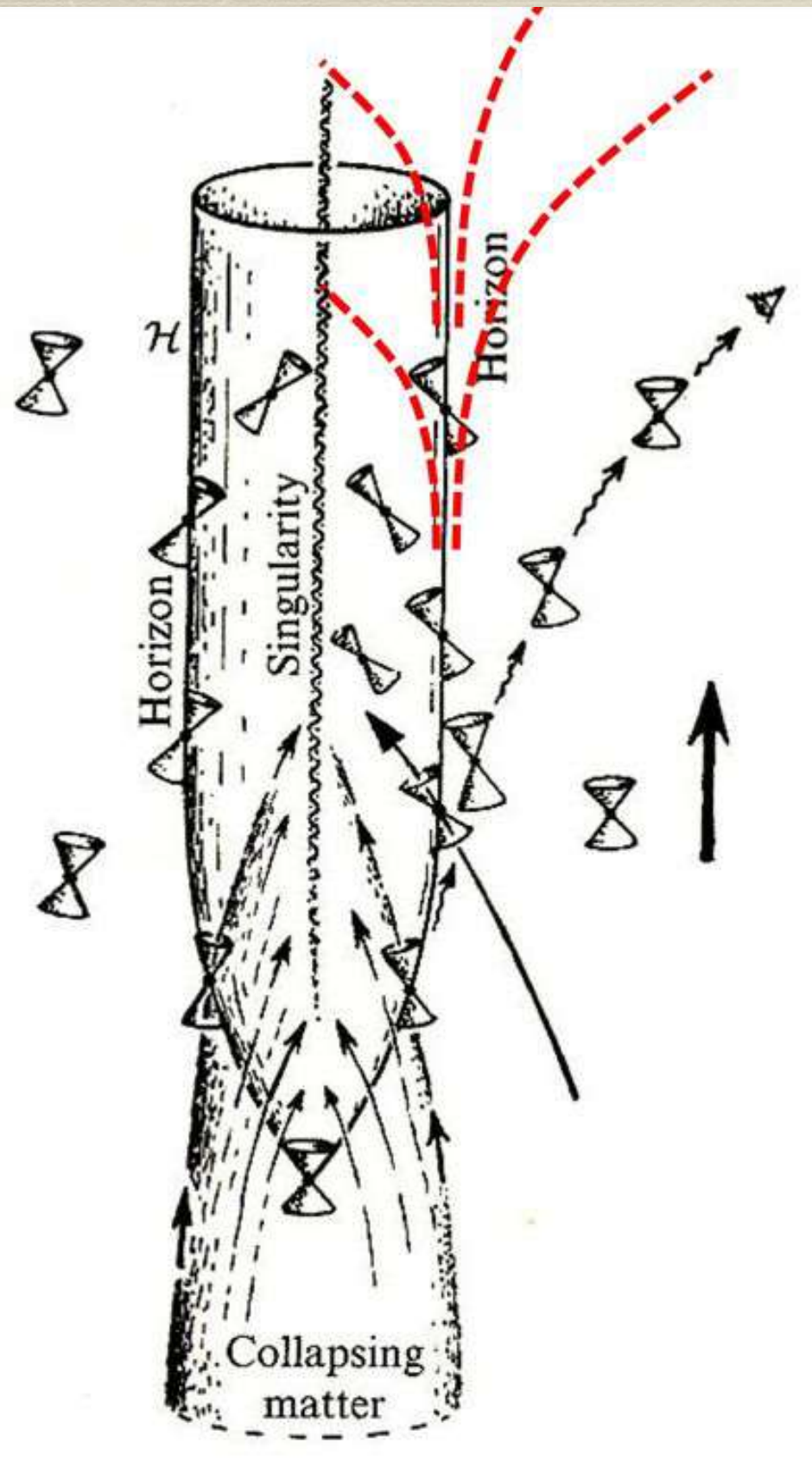
Invariance of Wave Action (number of particles in QM!) :

$$\frac{E_U}{\omega - \mathbf{U} \cdot \mathbf{k}} = \frac{E_0}{\omega}$$

$$\omega - \mathbf{U} \cdot \mathbf{k} = \pm \sqrt{gk \tanh(kh)}$$

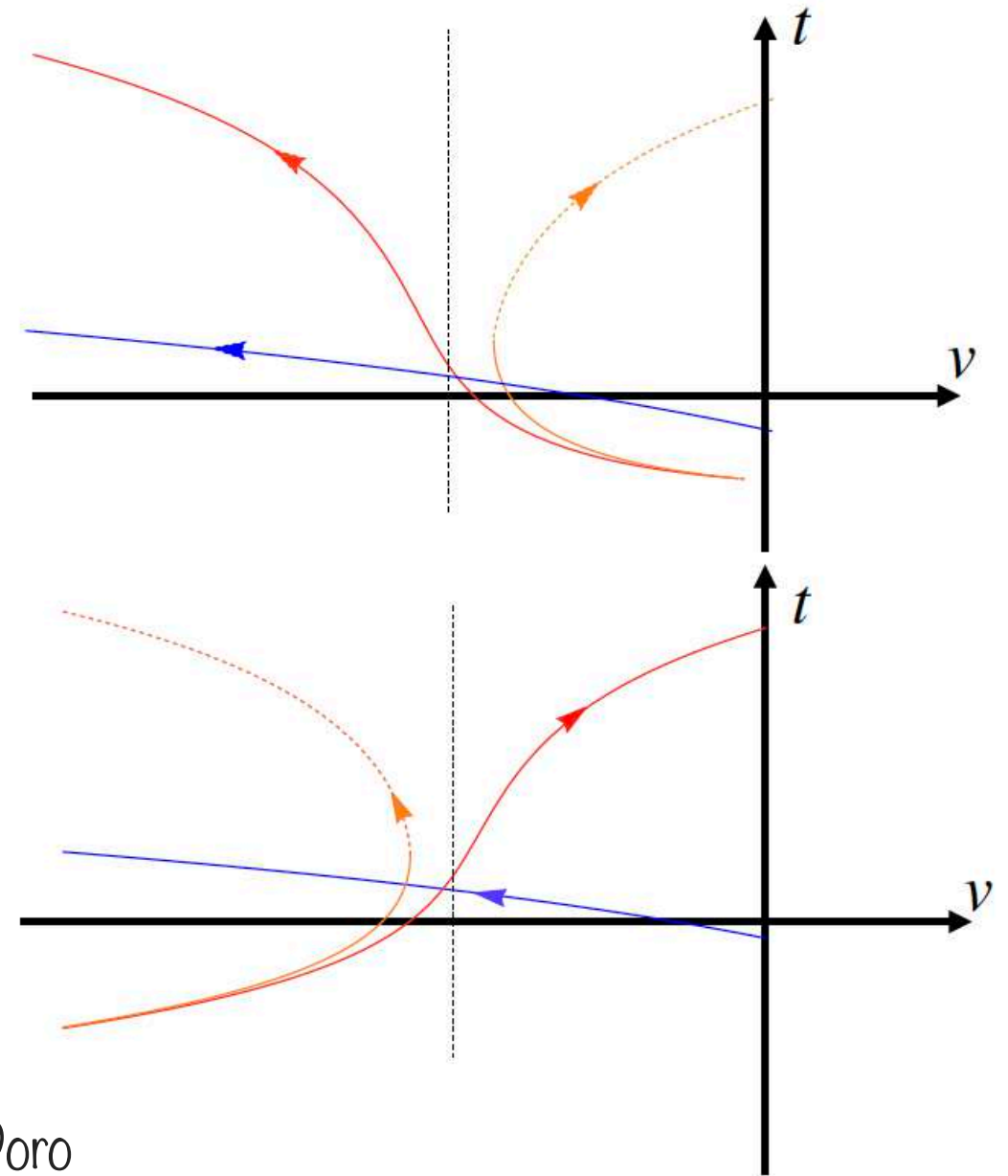
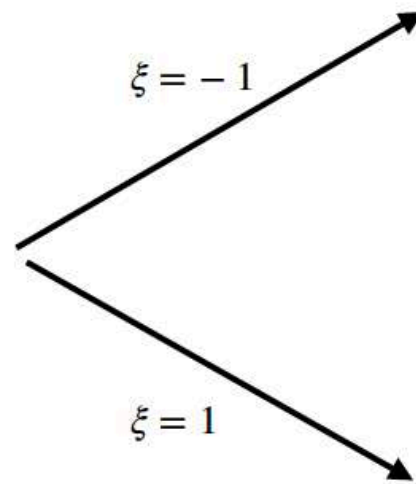
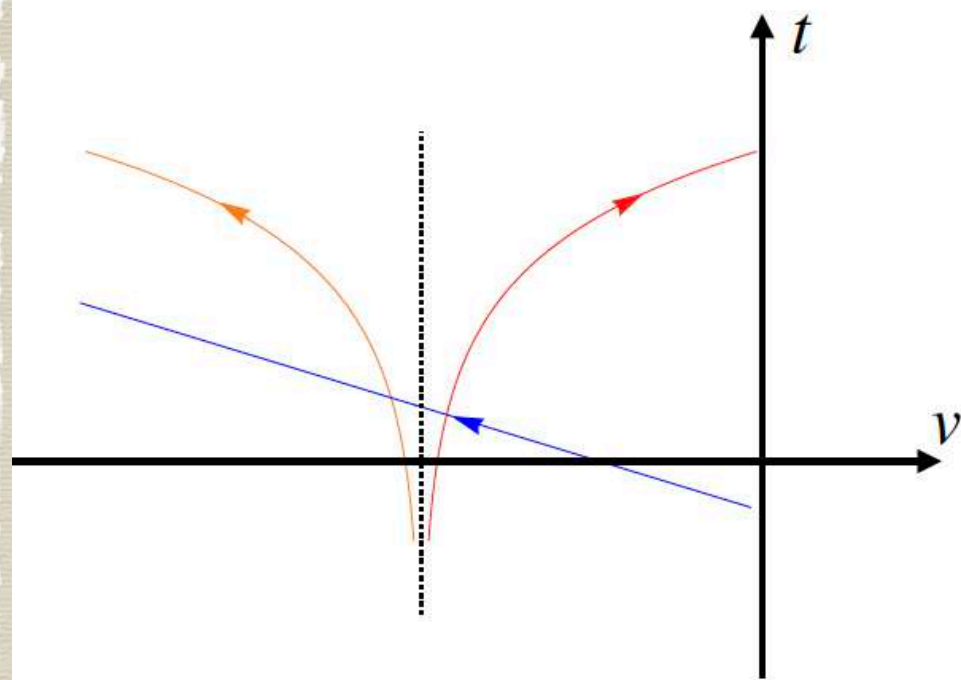
Negative energy waves are waves
with negative relative frequencies !

Ancestors of Hawking Radiation were created before the Collapse



Hawking Radiation With Modified Dispersion Relation

$$\omega^2 = c_s^2 k^2 \longrightarrow \omega^2 = c_s^2 k^2 \left(1 + \xi \frac{c_s^2 k^2}{\Lambda^2} \right)$$



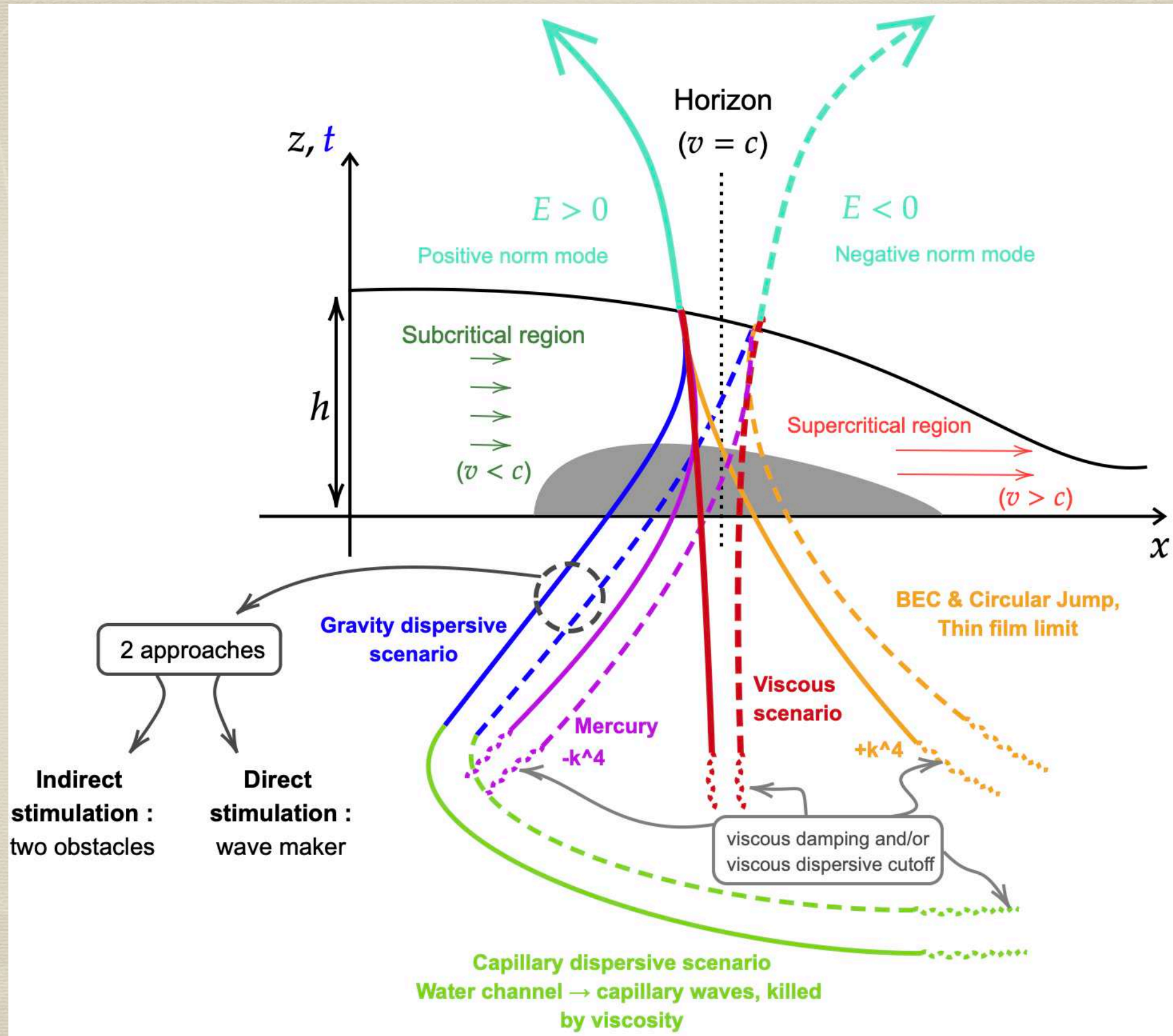
Hawking radiation is supposed to be robust against
the modification of the dispersion relations

Courtesy Francesco del Poro

Francesco del Poro, Beyond Lorentz invariance: a journey from Analogue to Hořava Gravity,
PhD Thesis (2024) under the supervision of Stefano Liberati
<https://iris.sissa.it/handle/20.500.11767/141091>

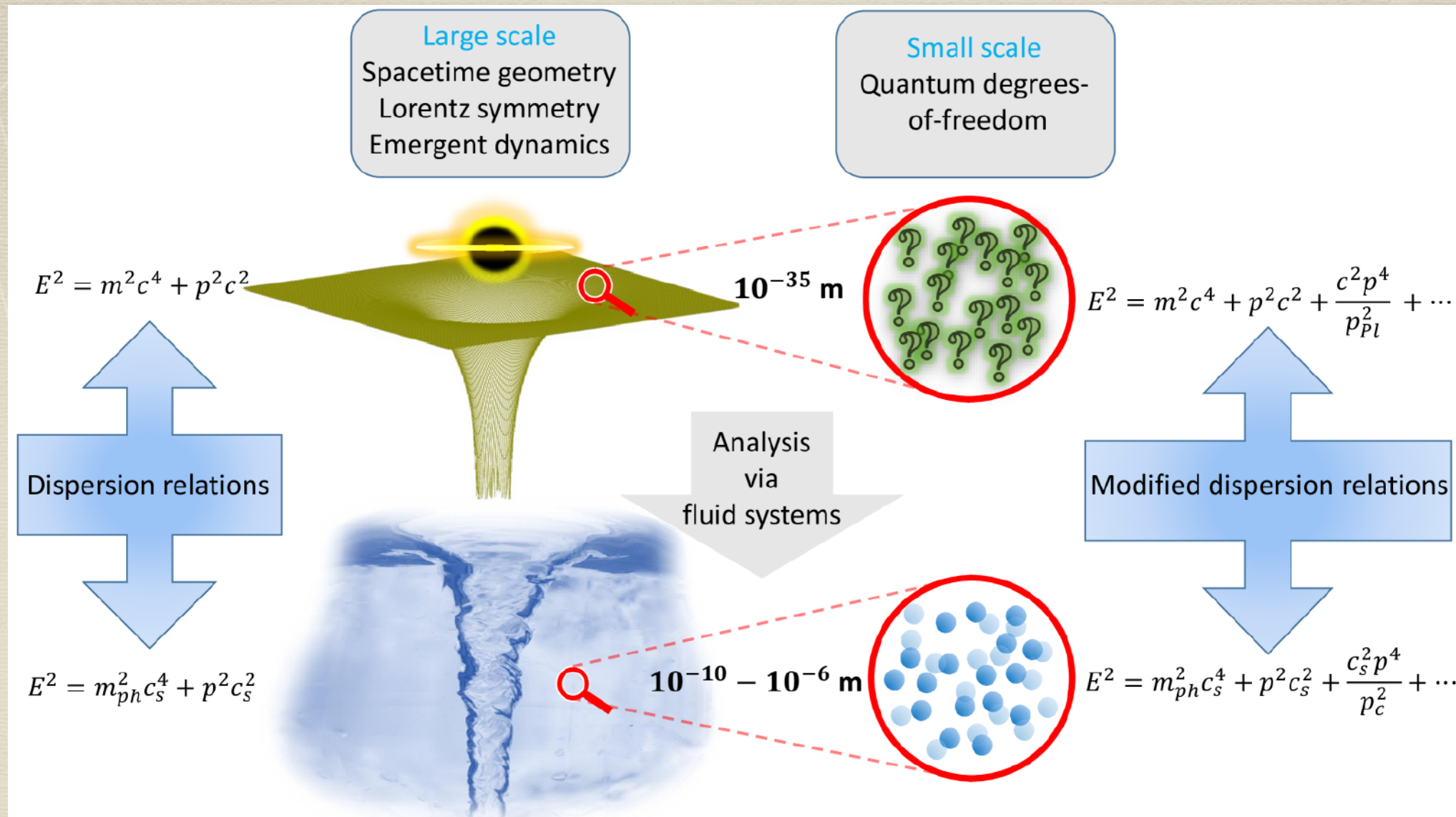
Francesco Del Porro, Stefano Liberati and Marc Schneider, Tunneling method for Hawking quanta in analogue gravity,
Comptes Rendus Physique 25, S2, p. 1-27 (2025)

Many Dispersive and Dissipative Scenarii for Wave Scattering at a BHH



Courtesy Ludivine Goncalves (current PhD Thesis)

Analogue Simulations of Quantum Gravity with Fluids



In fluids, the effective geometry breaks down at scales where the continuous description in terms of macroscopic variables no longer holds, in close analogy to the breakdown of spacetime geometry due to quantum gravitational effects. This scale is manifest by a breakdown of Lorentz symmetry as evidenced by the above modified phonon dispersion relations (see figure, right). There, m_{ph} is the mass of phonon excitations which is the analogue of the rest mass m , with energy and momentum being denoted by E and p for either system. For $m_{ph} = 0$ and truncating the expansion at the fourth-order in p the above expression reduces to the well-known Bogoliubov dispersion relation in a Bose gas [100, 101] or in a circular hydraulic jump [104, 105]. The critical momentum $p_c = M c_s$, that here plays the role of the Planck momentum p_{Pl} , depends on the mass M of the particles forming the condensate and is inversely proportional to the coherence length $\xi = \hbar/(M c_s)$. When $p \ll p_c$, the excitations follow the standard relativistic energy-momentum relation.

